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Review, Jan 27, 2012
We showed,
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$$\boldsymbol{\sigma} \cdot \mathbf{p} \psi_A = \frac{1}{\varsigma} (E - V(r) + mc^2) \psi_B.$$

$$\boldsymbol{\sigma} \cdot \mathbf{p} \psi_B = \frac{1}{c} (E - V(r) - mc^2) \psi_A.$$

and used

$$\begin{aligned} \boldsymbol{\sigma} \cdot \mathbf{p} &= \frac{1}{f} (\boldsymbol{\sigma} \cdot \hat{r}) \left(-i\hbar r \frac{\partial}{\partial r} + i\boldsymbol{\sigma} \cdot \mathbf{L} \right) \\ \psi_{A,B} &= \frac{1}{r} \phi_{A,B}(r) \mathcal{Y}_{j,m,\ell_{A,B}} \\ \boldsymbol{\sigma} \cdot \hat{r} \mathcal{Y}_{j,m,\ell=j\pm 1/2} &= -\mathcal{Y}_{j,m,\ell=j\mp 1/2} \\ \boldsymbol{\sigma} \cdot \mathbf{L} \mathcal{Y}_{j,m,\ell} &= \frac{2}{\hbar} \mathbf{S} \cdot \mathbf{L} \mathcal{Y}_{j,m,\ell} = \frac{2}{\hbar} (j(j+1) - \ell(\ell+1) - 3/4) \mathcal{Y}_{j,m,\ell} \end{aligned}$$

to obtain

$$\hbar c \left(\frac{dF}{dr} + \frac{\kappa}{r}F\right) = (E - V(r) + mc^2)G$$

$$\hbar c \left(\frac{dG}{dr} - \frac{\kappa}{r} G \right) = -(E - V(r) - mc^2)F$$

Our task is to solve these coupled first order differential equations.

Solution

Use atomic units and rationalize variables $\varepsilon = E/mc^2 \quad e^2/\hbar c = \alpha = \text{fine structure constant}, \quad c = \frac{1}{\alpha},$ $\lambda = \frac{1}{\alpha\sqrt{1-\varepsilon^2}}, \quad \rho = 2\lambda r,$

$$\frac{dG}{d\rho} - \frac{\kappa}{\rho}G = \left(\frac{1}{2}\sqrt{\frac{1-\varepsilon}{1+\varepsilon}} - \frac{Z\alpha}{\rho}\right)F$$
$$\frac{dF}{d\rho} + \frac{\kappa}{\rho}F = \left(\frac{1}{2}\sqrt{\frac{1+\varepsilon}{1-\varepsilon}} + \frac{Z\alpha}{\rho}\right)G$$

$$F = \sqrt{1+\varepsilon}e^{-
ho/2}(\phi_1+\phi_2), \ G = \sqrt{1-\varepsilon}e^{-
ho/2}(\phi_1-\phi_2)$$
 to get

$$\frac{d\phi_1}{d\rho} = \left(1 - \frac{\alpha\varepsilon}{\sqrt{1 - \varepsilon^2 \frac{Z}{\rho}}}\right)\phi_1 + \left(-\frac{\kappa}{\rho} - \frac{Z\alpha}{\sqrt{1 - \varepsilon^2}}\right)\phi_2 \quad (1)$$
$$\frac{d\phi_2}{d\rho} = \left(-\frac{\kappa}{\rho} + \frac{\alpha}{\sqrt{1 - \varepsilon^2}}\frac{Z}{\rho}\right)\phi_1 + \frac{\alpha\varepsilon}{\sqrt{1 - \varepsilon^2 \frac{Z}{\rho}}}\phi_2 \quad (2)$$

Solution continued

Solve Eq.(2) for ϕ_2 and substitute into Eq.(1). This gives a second order differential equation for ϕ_1 whose solution can be recognized as

$$\phi_{1} = \frac{\mathbf{n}'}{\sqrt{-\kappa + \alpha Z/\sqrt{1-\varepsilon^{2}}}} \rho^{\gamma} F(-\mathbf{n}'+1, 2\gamma+1, \rho)$$

$$\phi_{2} = \sqrt{-\kappa + \alpha Z/\sqrt{1-\varepsilon^{2}}} \rho^{\gamma} F(-\mathbf{n}', 2\gamma+1, \rho)$$
where

$$n' = \frac{Z\alpha\varepsilon}{\sqrt{1-\varepsilon^2}} - \gamma$$
$$\gamma = \sqrt{\kappa^2 - Z^2\alpha^2}$$

and F(a, b, z) is the confluent hypergeometric function of Eq.(A.59) in Appendix A.5. As for the Schrödinger equation, bound states occur when the series for F(a, b, z) terminates. This requires that n' is an integer greater than or equal to 1. Solve for ε in terms of n':

$$\begin{split} \varepsilon &= \frac{1}{\sqrt{1 + \frac{Z^2 \alpha^2}{(n'+\gamma)^2}}} = \frac{1}{\sqrt{1 + \frac{Z^2 \alpha^2}{(n-\delta_j)^2}}}\\ \text{With } \kappa^2 &= (j+1/2)^2, \ n' = n - |\kappa| \text{ and } \delta_j = |\kappa| - \gamma \text{ we see that the equation for the eigenenergy agrees with Eq. (2.35) of the text.} \end{split}$$