

Physics 601, Homework I

Problemw 1 and 2; To solve these problems I simply used Eq. (15) of the notes and computed $B(E)$ using that $\phi_{\text{in}}(r) = A\sqrt{2(E + V_0)}$ for both positive and negative E . Then $B(E) = k_{\text{in}}(E) \cot[k_{\text{in}}(E)r_0]$, where $k_{\text{in}}(E) = \sqrt{2(E + V_0)}$. This is then substituted into Eq.(15) to get $K = \tan \delta(E)$. The attached Mathematica notebook shows a plot of the phase shift for $E > 0$. But we can also use the same formulas for $E < 0$ to find where $K = i$. Plots are also shown where the one bound state is located. A check is made to see that there is no bound state close to $E = 0$. Finally, since $\ell = 0$ solutions are simply sin functions the solution was also obtained by matching the functions directly using $\phi(r) = A \sin(k_{\text{in}}r + \delta)$ and $\phi_{\text{in}} = A_{\infty} \sin(kr)$ as a check. Notice that when using the arctan function it is necessary to eliminate jumps by π in the phase shift to obtain a smooth plot.

Problem 3. In this case it is easier match solutions at boundaries directly. There are three regions call them $j = 0, 1, 2$. In each region the solution is $\phi_j = A_j \sin(k_j r + \delta_j)$. The functions and their derivatives are matched at the boundaries r_j , $j = 0, 1$ using that $\delta_0 = 0$. This gives the equations

$$k_{j+1} \cot(k_{j+1}r_j + \delta_j) = k_j \cot(k_j r_j + \delta_j), \quad j = 0, 1 \quad (1)$$

so that

$$\delta_2 = \delta = \text{arccot}((k_1/k_2) \cot(k_1 r_1 + \delta_1)) - k_2 r_1 \quad (2)$$

$$\delta_1 = \text{arccot}((k_0/k_1) \cot(k_0 r_0 + \delta_0)) - k_1 r_0 \quad (3)$$

$$\delta_0 = 0 \quad (4)$$

These equations can be used for all E , provide one takes $\sqrt{-C} = i\sqrt{C}$, $C > 0$. Mathematica does this automatically. The second attached notebook shows the solution in this case. Note that there are lot of jumps by π .

We could also use equation (15) or (16) of the notes. In this case $B(E) = k_1 \cot(k_1 r_1 + \delta_1)$ is the logarithmic derivative at the boundary to the outer region.

Use of B(E) matrix for a square well

■ Set up u, v, uin, derivatives, and B(E) (mainly inner region)

```

V0 = 2.
2.
u[x_] := Sin[x]
v[x_] := Cos[x]
kin[V0_, en_] := Sqrt[2 * (en + V0)]
B[en_, r_] := kin[V0, en] * Cot[kin[V0, en] * r]

```

■ Set up K=Tan(delta) (mainly outer region).

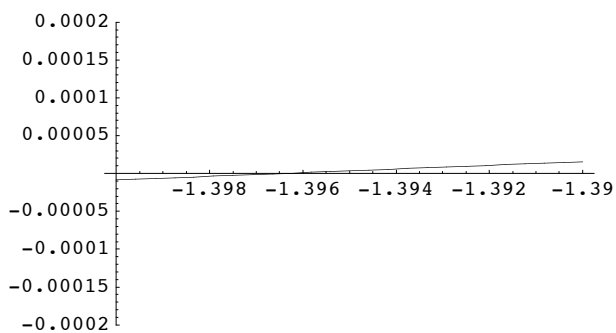
```

k[en_] := Sqrt[2 * en]
K[en_, r0_] :=
-u[k[en] * r0] / v[k[en] * r0] + (k[en] / v[k[en] * r0]^2) / (B[en, r0] + k[en] * Tan[k[en] * r0])

```

■ Find values of energy for which K=-i.

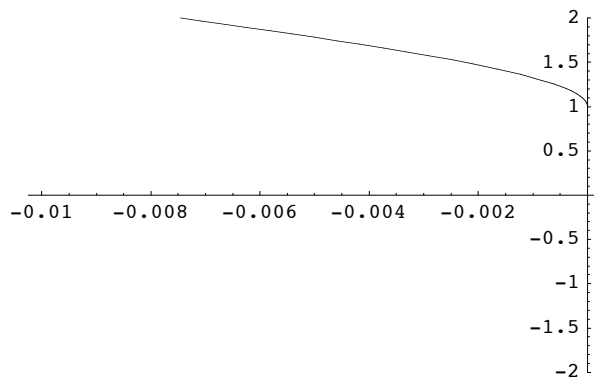
```
Plot[Im[K[en, 2.33] + I], {en, -1.4, -1.39}, PlotRange -> {- .0002, .0002}];
```



```
K[-1.398, 2.333]
```

```
0. - 1. i
```

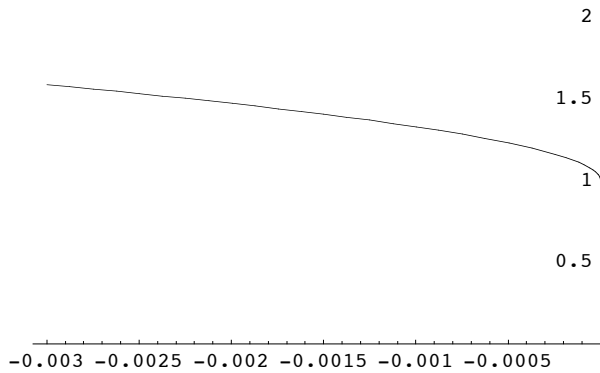
```
Plot[Im[K[en, 2.33] + I], {en, -.01, -.0}, PlotRange -> {-2, 2}];
```



```
K[-0.000001, 2.33]
```

```
0. + 0.0101899 i
```

```
Plot[Im[K[en, 2.33] + I], {en, -.003, -.0}, PlotRange -> {-.2, 2}];
```



```
K[-0.000001, 2.33] / k[-0.000001]
```

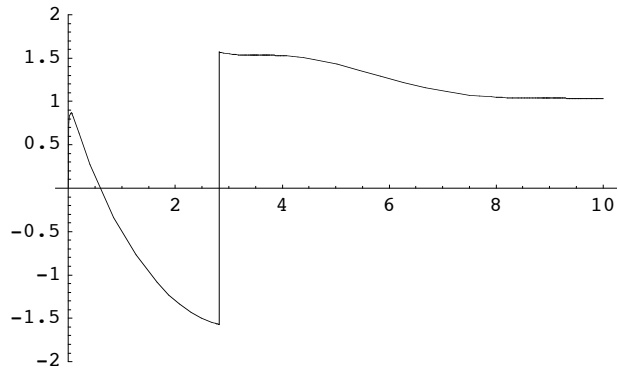
```
7.20538 + 0. i
```

We see that there is one bound state, at about $E=-1.4$. There is no bound state at $k=0.0$

■ Find the phase shift $\delta(E)$.

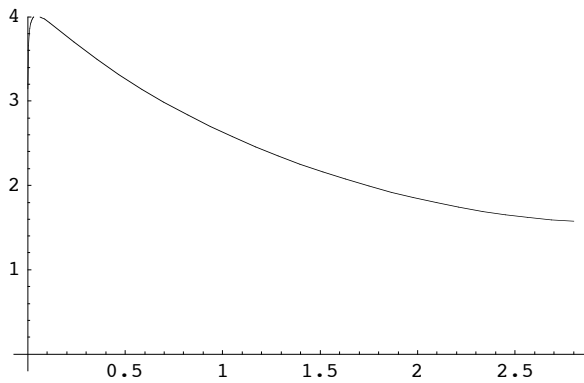
```
delta[en_, r0_] := ArcTan[K[en, r0]]
```

```
Plot[delta[en, 2.33], {en, 0, 10.}, PlotRange -> {-2, 2}];
```

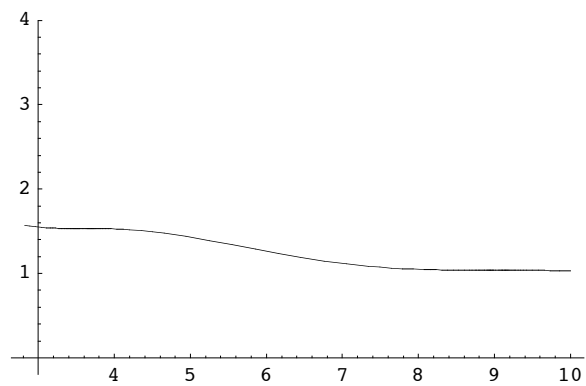


Eliminate the break at $E \sim 3$ au. The steps to do this are given below.

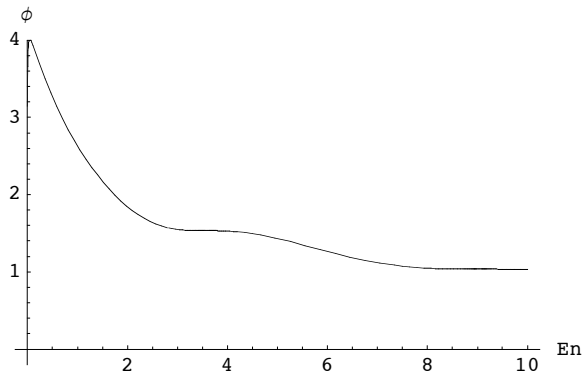
```
p1 = Plot[delta[en, 2.33] + Pi, {en, 0, 2.8}, PlotRange -> {-0.2, 4}];
```



```
p2 = Plot[delta[en, 2.33], {en, 2.82, 10}, PlotRange -> {-0.2, 4}];
```



```
pf = Show[p1, p2, AxesLabel -> {En,  $\phi$ };
```

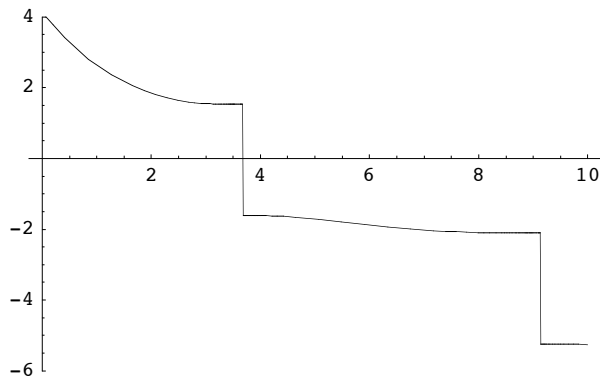


The breaks have now been completely eliminated.

■ Check: Use a more direct solution.

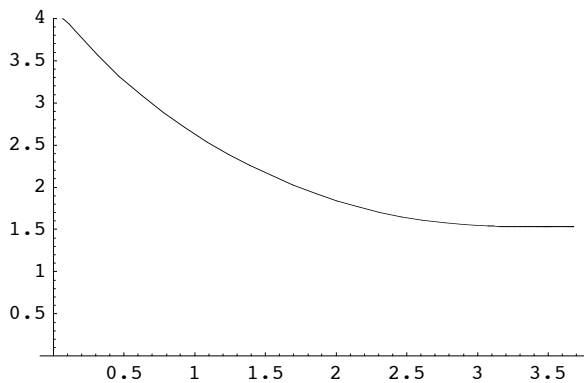
```
deltaa[en_, r0_] := ArcTan[(k[en] / kin[V0, en]) * Tan[kin[V0, en] * r0]] - k[en] * r0
```

```
pla = Plot[deltaa[en, 2.33] + 2 * Pi, {en, 0.001, 10}, PlotRange -> {-6, 4};
```

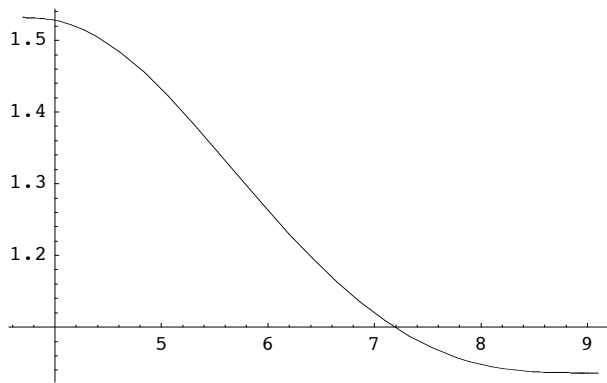


Eliminate the breaks. Note that this way shows 3 breaks.

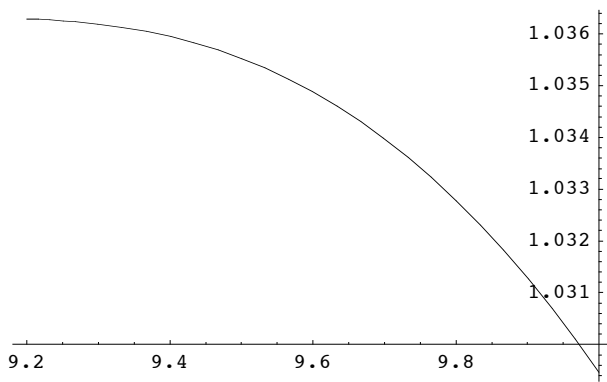
```
pla = Plot[deltaa[en, 2.33] + 2 * Pi, {en, 0.001, 3.68}, PlotRange -> {0, 4};
```



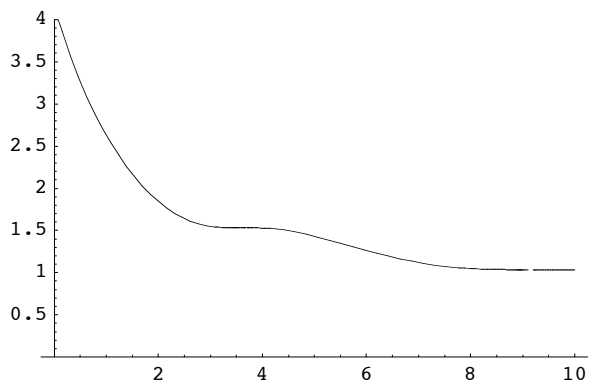
```
p2a = Plot[deltaa[en, 2.33] + 3 * Pi, {en, 3.7, 9.1}];
```



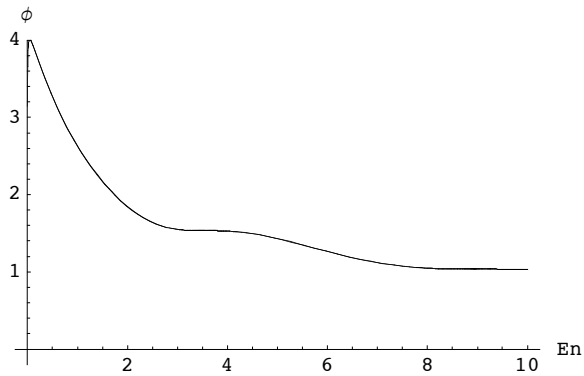
```
p3a = Plot[deltaa[en, 2.33] + 4 * Pi, {en, 9.2, 10}];
```



```
pfp = Show[p1a, p2a, p3a];
```



```
Show[pf, pf];
```

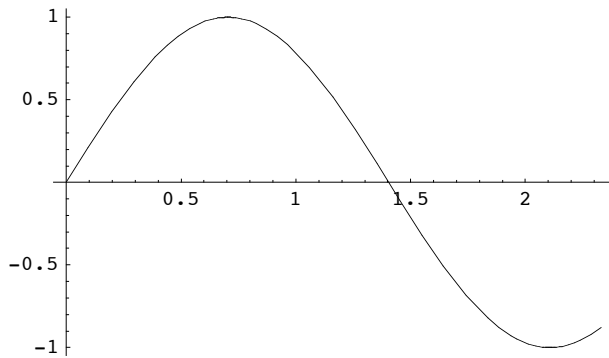


This plot shows that the same phase shift is obtained both ways.

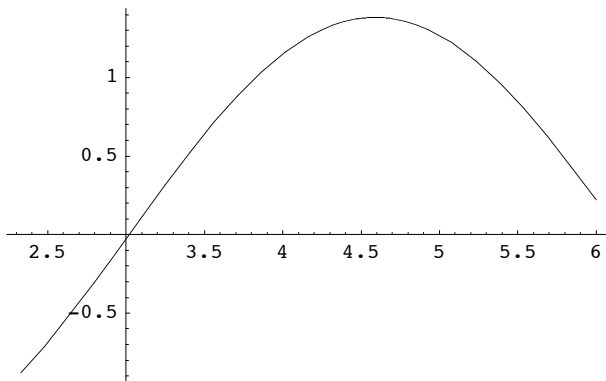
```
Phiin[en_, r_] := Sin[kin[2, en] * r]
delt[en_, r0_] := ArcTan[K[en, r0]]
Phi[en_, r_, r0_] :=
  Sin[k[en] * r + delt[en, r0]] * (Sin[kin[2, en] * r0] / Sin[k[en] * r0 + delt[en, r0]])

Phiin[en_, r_] := Sin[kin[2, en] * r]

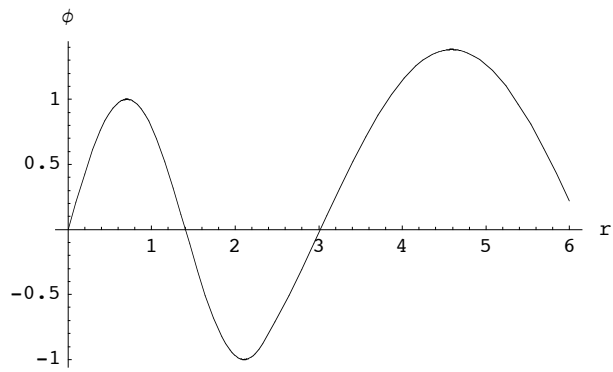
p1 = Plot[Phiin[.5, r], {r, 0, 2.33}];
```



```
p2 = Plot[Phi[.5, r, 2.33], {r, 2.33, 6}];
```



```
Show[p1, p2, AxesLabel -> {r,  $\phi$ }]
```



- Graphics -

Above is a plot of the wave function ϕ vs. r for $E=0.5$ au.

■ Program for Eqs. (2) and (3) of the answer sheet

```

k0[en_] := Sqrt[2 * (en + 2.)]
k1[en_] := Sqrt[2 * (en - 2.)]
k2[en_] := Sqrt[2 * (en)]

r0 = 2.33
2.33

r1 = 2.6
2.6

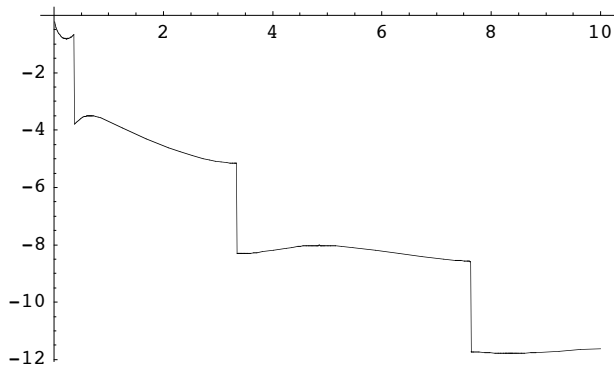
Clear[delta1]

delta1[en_] := ArcCot[(k0[en] / k1[en]) * Cot[k0[en] * r0]] - k1[en] * r0

delta2[en_] := ArcCot[(k1[en] / k2[en]) * Cot[k1[en] * r1 + delta1[en]]] - k2[en] * r1

Plot[delta2[en], {en, 0.01, 10}];

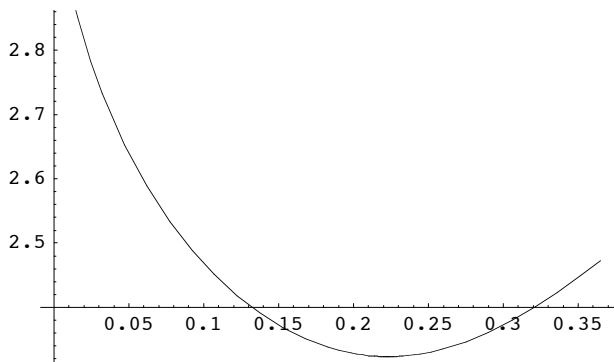
```



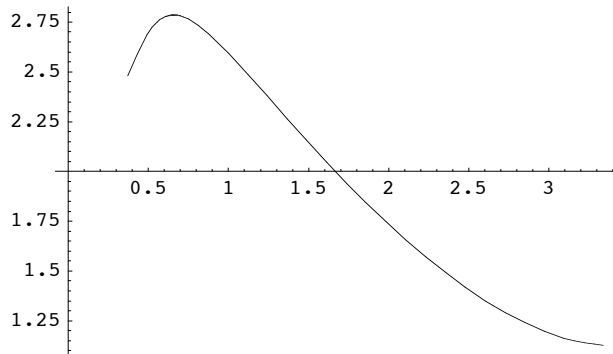
```

p1 = Plot[delta2[en] + Pi, {en, 0.001, 0.365}];

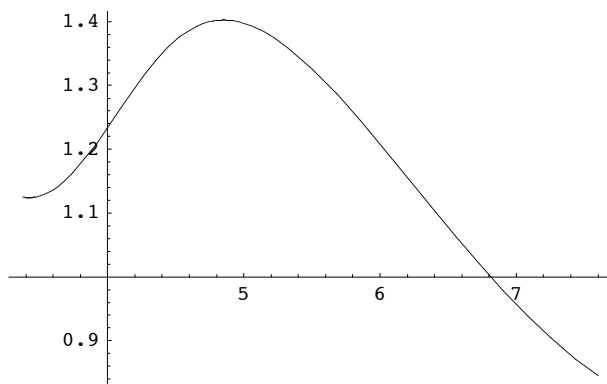
```



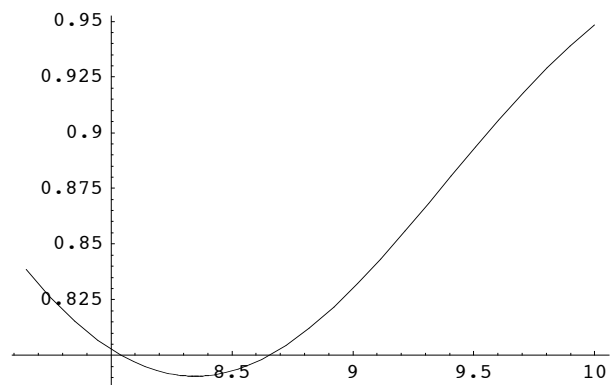
```
p2 = Plot[delta2[en] + 2 * Pi, {en, 0.37, 3.34}];
```



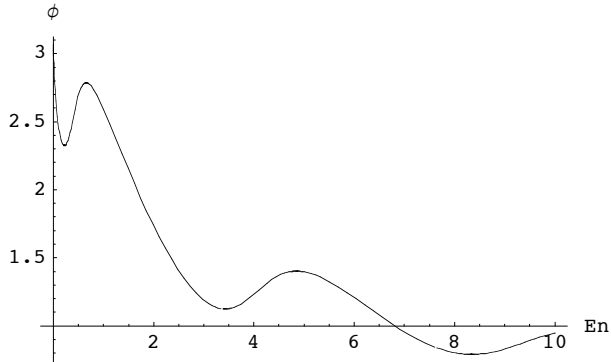
```
p3 = Plot[delta2[en] + 3 * Pi, {en, 3.38, 7.6}];
```



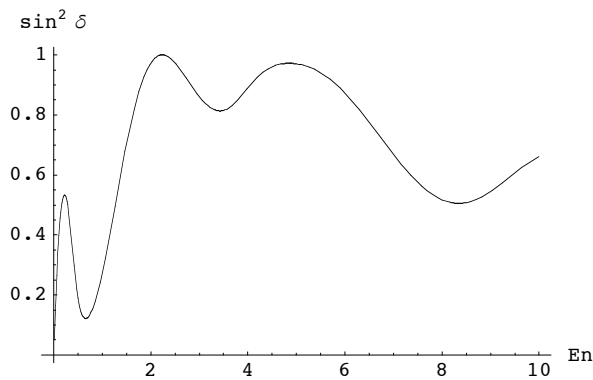
```
p4 = Plot[delta2[en] + 4 * Pi, {en, 7.65, 10}];
```



```
Show[p1, p2, p3, p4, AxesLabel -> {En,  $\phi$ }] ;
```



```
Plot[Sin[delta2[en]]^2, {en, 0.01, 10}, AxesLabel -> {En, sin^2( $\delta$ )}];
```

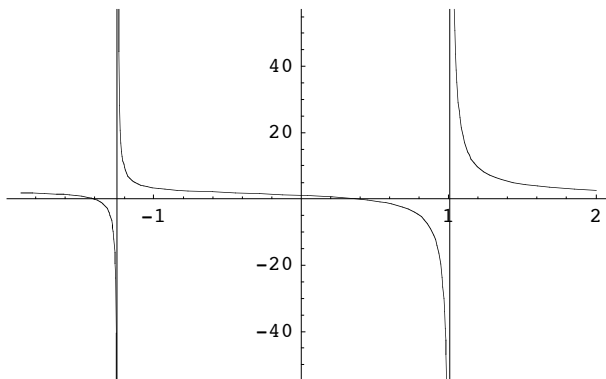


We have not discussed this yet, but the \sin^2 of the phase shift is the physically significant parameter. It is a smooth function of the phase shift and does not show the effects of jumps.

We can also calculate $B(E)$ from the logarithmic derivative of the inner function at $r=r_1$.

```
B[en_] := k1[en] * Cot[k1[en] * r1 + delta1[en]]
```

```
Plot[B[en], {en, -1.9, 2}];
```



Notice that $B(E)$ is real for all values of E . It is an analytic function of E except at the first order poles where it becomes infinite.

Use of B(E) matrix for a square well

■ Set up u, v, uin, derivatives, and B(E) (mainly inner region)

```

v0 = 2
2
u[x_] := Sin[x]
v[x_] := Cos[x]
kin[v0_, en_] := Sqrt[2 * (en + v0)]
B[en_, r_] := kin[v0, en] * Cot[kin[v0, en] * r]

```

■ Set up K=Tan(delta) (mainly outer region).

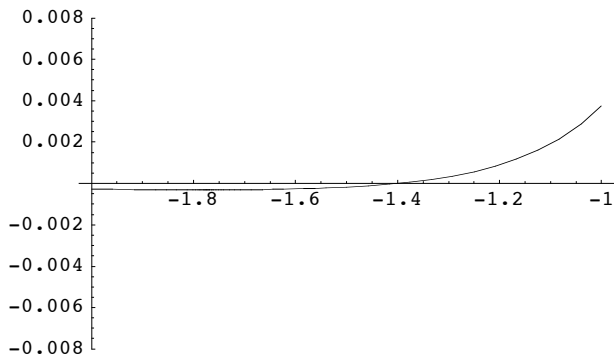
```

k[en_] := Sqrt[2 * en]
K[en_, r0_] :=
-u[k[en] * r0] / v[k[en] * r0] + (k[en] / v[k[en] * r0]^2) / (B[en, r0] + k[en] * Tan[k[en] * r0])

```

■ Find values of energy for which K=-i.

```
Plot[Im[K[en, 2.33] + I], {en, -2, -1}, PlotRange -> {- .008, .008}];
```



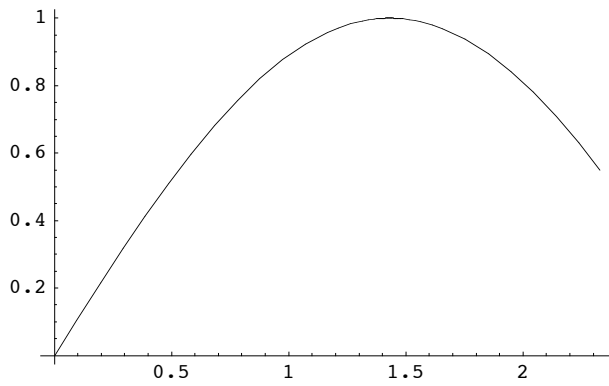
```

K[-1.396416, 2.33] + I
0. - 4.01306 × 10-10 i

```

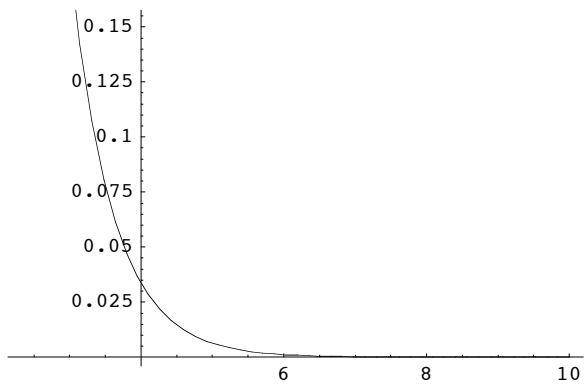
This verifies that the bound state energy is -1.396416 au \sim -1.4 au. The wave function is shown below.

```
p1 = Plot[Sin[kin[V0, -1.396416] * r], {r, 0, 2.33}];
```



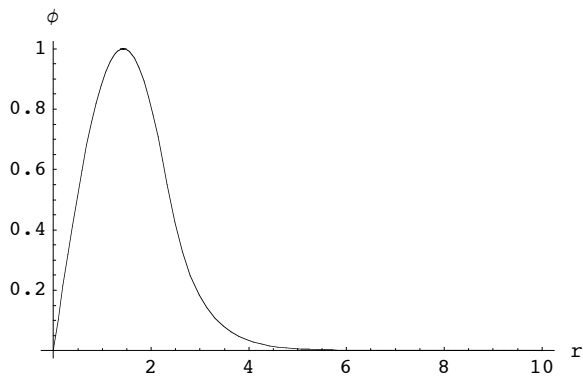
- Graphics -

```
p2 = Plot[Sin[kin[V0, -1.396416] * 2.33] * Exp[I * k[-1.396416] * (r - 2.33)], {r, 2.33, 10}];
```



- Graphics -

```
Show[p1, p2, AxesLabel -> {r,  $\phi$ }]
```



The above is a graph of the un-normalized wave function.