

Low Energy Nuclear Experiments

Ben Kay, Argonne National Laboratory

National Nuclear Physics Summer School, 8-19 July 2019

Overview, part 1 (*general properties of nuclei, mostly macroscopic*)

What can experimentalists determine about a nuclear system in the lab?

- *History ... the isotopes, the facilities we use*
- *What can we measure/is observable?*
- *Questions to ask about the nucleus*
 - *How much do they weigh?*
 - *What size are they?*
 - *What shape are they?*

Attempt to use many accessible examples from recent literature, leaning towards the study of exotic nuclei where possible

Overview, part 2 *(mostly direct reactions, not so exotic, microscopic)*

The connection between direct reactions and nuclear structure

- *History*
- *Reactions, reaction types, direct reactions*
- *Observables*
- *Energies, momentum*
- *Spectroscopic factors, occupancies (in context of 'modern' [but stable-beam] examples)*
- *Other reactions (pairing, cluster, charge exchange)*

Attempt to steer clear of reactions for reaction's sake, rather using them as a meaningful tool to gain insights into topical nuclear structure properties

Overview, part 3 *(mostly direct reactions, quite exotic, microscopic)*

The connection between direction reactions and nuclear structure

- *History*
- *Exotic beams*
- *Kinematics*
- *Spectrometers (with a focus on solenoidal spectrometers)*
- *A few examples from the last few years (2014, 2017, 2017, 2019) (what drove them, reaction choices, results, commentary)*

Reading

- Slides from past schools (NNPSS [Heather Crawford's are exemplary], EBSS) are impressive (next slide for references)
- Books are good, but often dense and not always transparent (on direct reactions, my personal favorites are N. K. Glendenning's *Direct Nuclear Reactions*, and C. A. Bertulani and P. Danielewicz's *Introduction to Nuclear Reactions*.)
- Great papers (some of the older ones can be wonderfully pedagogical, others far less so). I will attempt to highlight some as I go through.

Past schools ... slides on reactions

2002 (ORNL 1st), 2003 (NSCL 2nd), 2004 (ANL 3rd), 2005 (LBNL 4th) 2006 (ORNL 5th), 2007 (NSCL 6th), 2008 (ANL 7th), 2009 (LBNL 8th) 2010 (ORNL 9th)

<https://people.nslc.msu.edu/~zegers/ebss2011/cizewski.pdf> (J. Cizewski of Rutgers, NSCL 2011) ... **10th in EBSS series**

<http://www.phy.anl.gov/atlas/EBSS2012/NuclearReactions-ExperimentI.pptx> (L. Trache of Texas A&M, ANL 2012) ... **11th**

<http://www.phy.anl.gov/atlas/EBSS2012/NuclearReactions-ExperimentII.pptx>

<http://fribusers.org/documents/2013/ebssLectures/reactions1.pdf> (Grigory Rogachev of FSU, LBNL 2013) ... **12th**

<http://fribusers.org/documents/2013/ebssLectures/reactions2.pdf>

<http://fribusers.org/documents/2013/ebssLectures/reactions3.pdf>

http://fribusers.org/documents/2014/ebssLectures/hoffman_1.pdf (Calem Hoffman of Argonne, ORNL 2014) ... **13th**

http://fribusers.org/documents/2014/ebssLectures/hoffman_2.pdf

http://aruna.physics.fsu.edu/ebss_lectures/F_Lecture2.pdf (Ben Kay of Argonne, FSU 2015) ... **14th**

https://people.nslc.msu.edu/~iwasaki/EBSS2016/reaction_1.pdf (Alan Wuosmaa of UConn, NSCL 2016) ... **15th**

https://people.nslc.msu.edu/~iwasaki/EBSS2016/reaction_2.pdf

https://people.nslc.msu.edu/~iwasaki/EBSS2016/reaction_3.pdf

<http://www.phy.anl.gov/ebss2017/ebss-2017-zegers.pdf> (Remco Zegers of NSCL, ANL 2017) ... **16th**

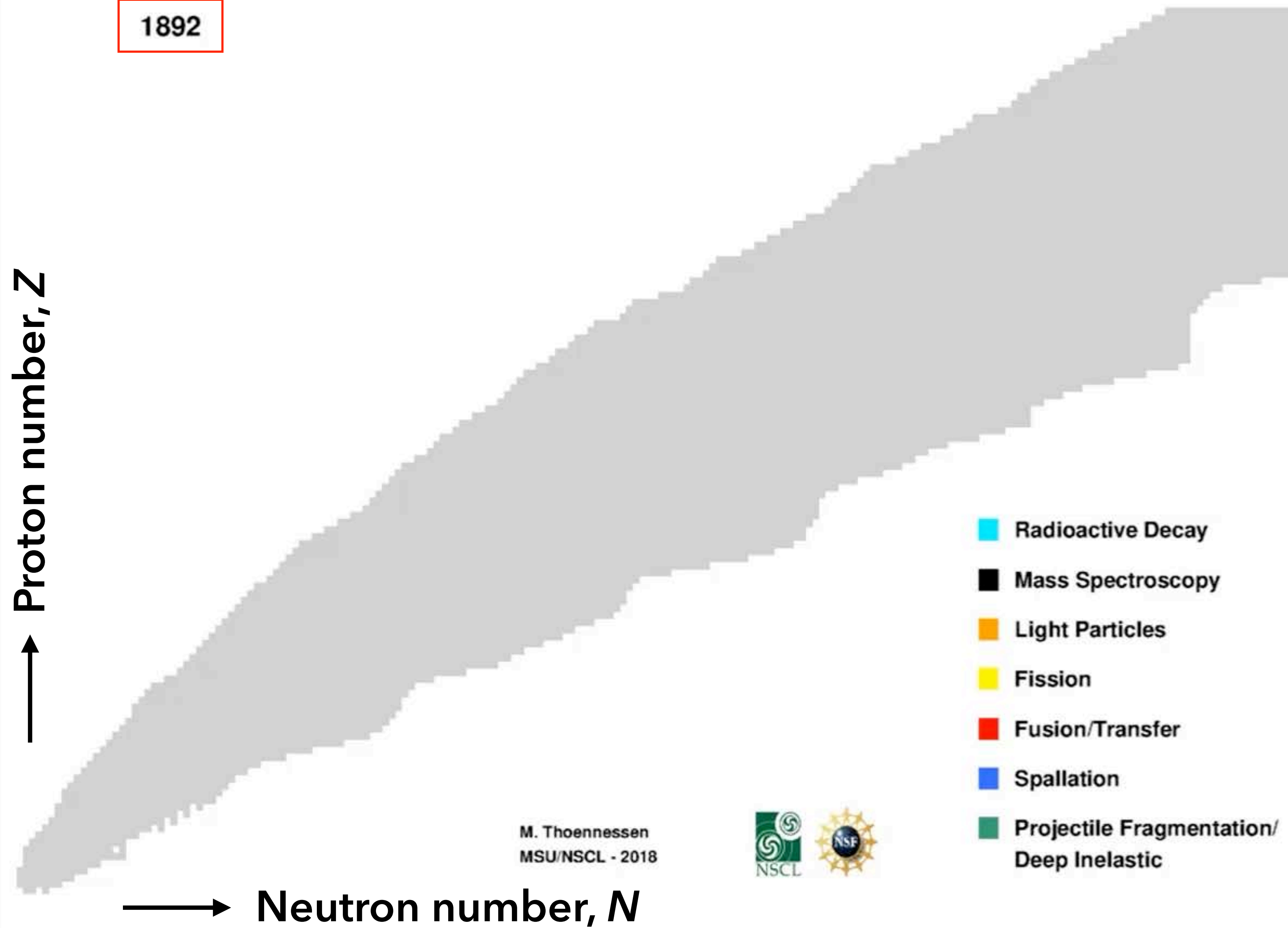
And soon, these slides ... **17th**

Part 1: General overview

Isotopes, masses, sizes, shapes

Isotopes

1892



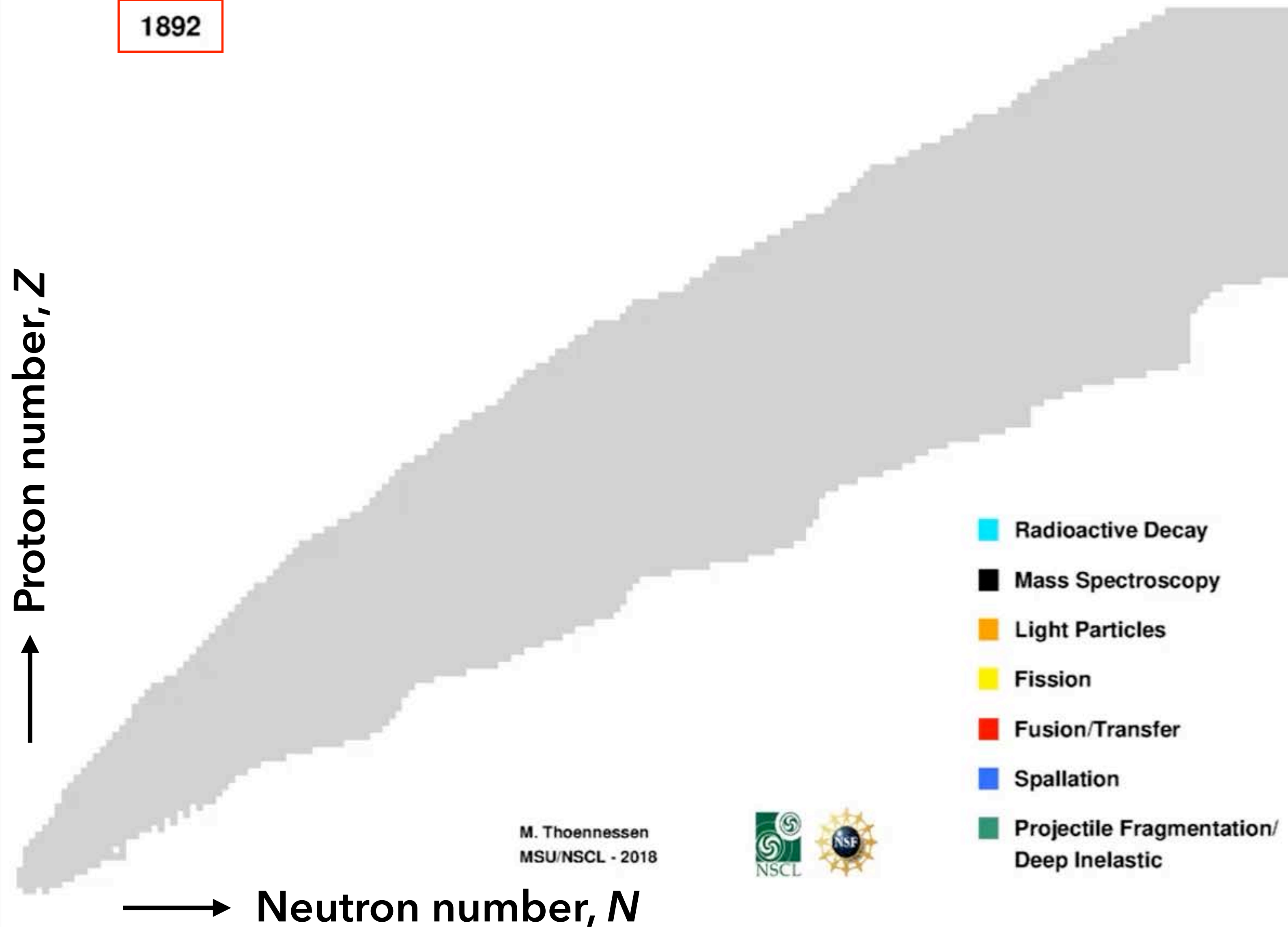
M. Thoennessen
MSU/NSCL - 2018



<https://people.nslc.msu.edu/~thoenness/isotopes/>

Isotopes

1892



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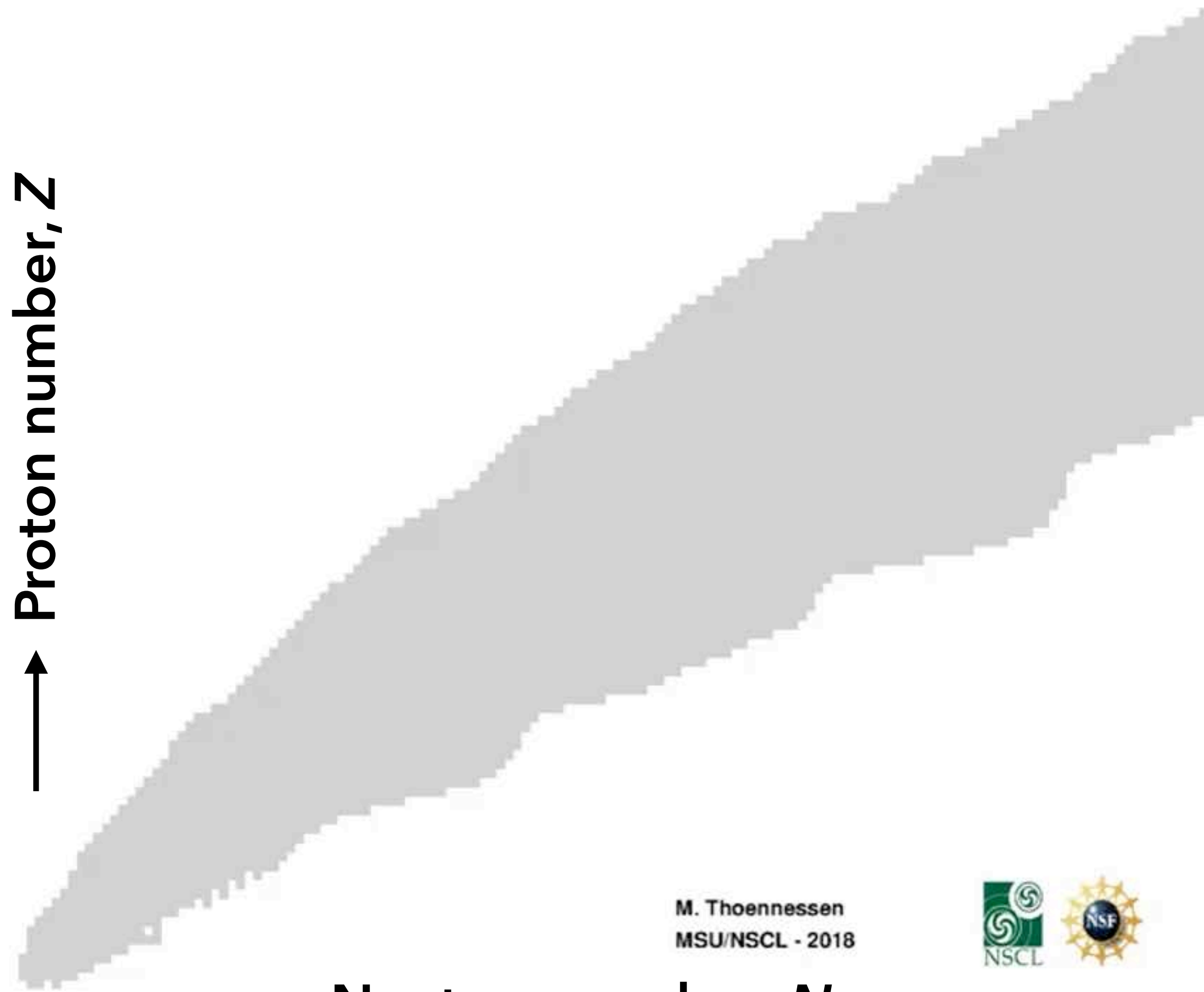


<https://people.nsl.mscl.msu.edu/~thoennes/isotopes/>

Isotopes

1892

Proton number, Z



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Neutron number, N

<https://people.nscl.msu.edu/~thoennes/isotopes/>

Physics Letters B 795 (2019) 266–270



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Contents lists available at ScienceDirect

Physics Letters B

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Discovery of ^{68}Br in secondary reactions of radioactive beams

K. Wimmer^{a,b,*}, P. Doornenbal^b, W. Korten^c, P. Aguilera^d, A. Algora^{e,f}, T. Ando^a, T. Arici^{g,h}, H. Baba^b, B. Blankⁱ, A. Boso^j, S. Chen^b, A. Corsi^c, P. Davies^k, G. de Angelis^l, G. de France^m, D.T. Doherty^c, J. Gerl^g, R. Gernhäuserⁿ, D.G. Jenkins^k, S. Koyama^a, T. Motobayashi^b, S. Nagamine^a, M. Niikura^a, A. Obertelli^{c,b}, D. Lubosⁿ, B. Rubio^e, E. Sahin^o, T.Y. Saito^a, H. Sakurai^{a,b}, L. Sinclair^k, D. Steppenbeck^b, R. Taniuchi^a, R. Wadsworth^k, M. Zielinska^c

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ARTICLE INFO

Article history:

Received 15 May 2019

Received in revised form 10 June 2019

Accepted 10 June 2019

Available online 12 June 2019

Editor: D.F. Geesaman

Keywords:

Radioactive beams

New isotope

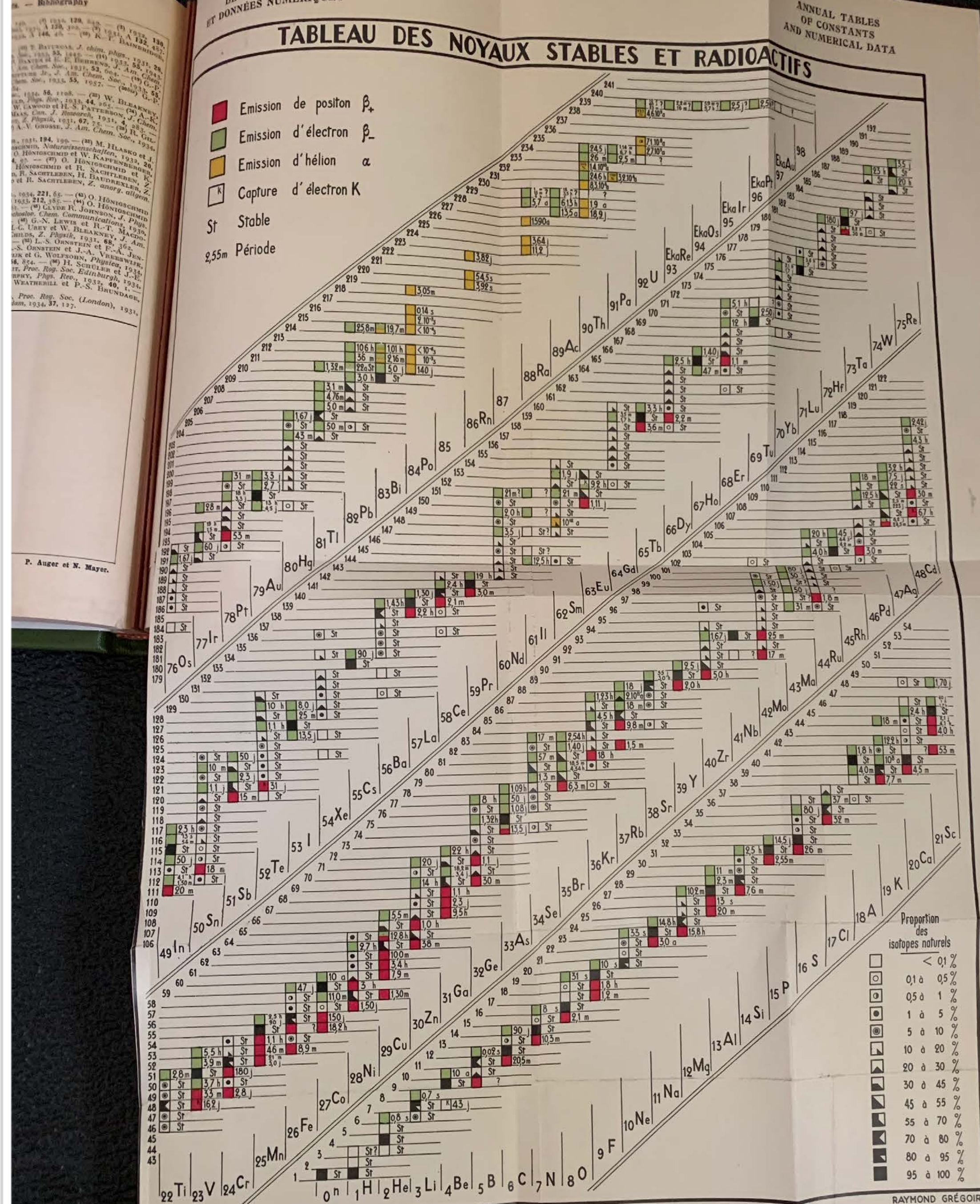
Direct reaction

ABSTRACT

The proton-rich isotope ^{68}Br was discovered in secondary fragmentation reactions of fast radioactive beams. Proton-rich secondary beams of $^{70,71,72}\text{Kr}$ and ^{70}Br , produced at the RIKEN Nishina Center and identified by the BigRIPS fragment separator, impinged on a secondary ^9Be target. Unambiguous particle identification behind the secondary target was achieved with the ZeroDegree spectrometer. Based on the expected direct production cross sections from neighboring isotopes, the lifetime of the ground or long-lived isomeric state of ^{68}Br was estimated. The results suggest that secondary fragmentation reactions, where relatively few nucleons are removed from the projectile, offer an alternative way to search for new isotopes, as these reactions populate preferentially low-lying states.

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1937

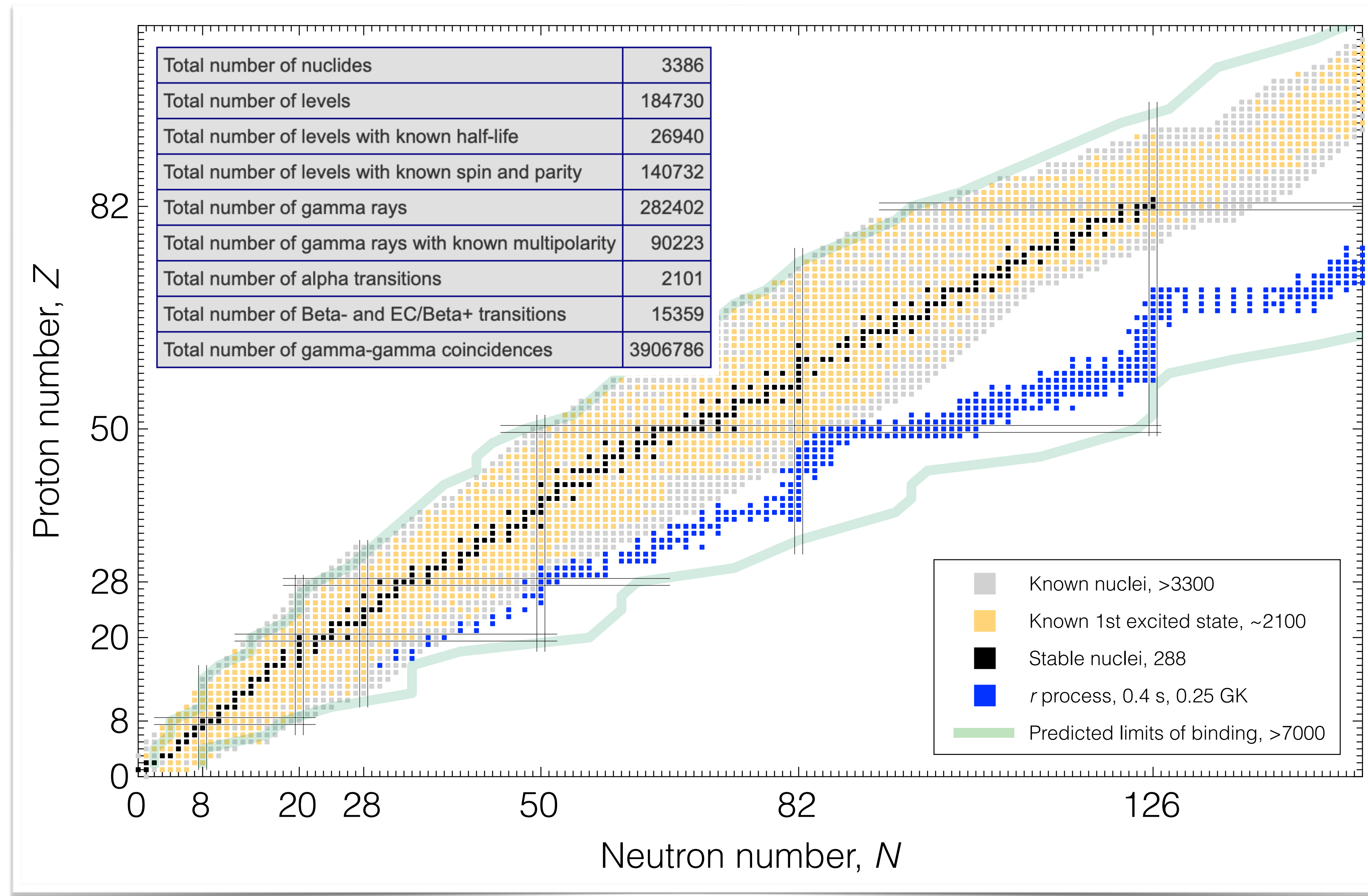


World pop: 2.3 B
Gas: 16 cents/gal.

From Table Annuelles Volume XI, Chapter 41-45, 1937 by Piere Curie

Today ...

World pop: 7.53 B
Gas: 250 cents/gal.



Limits from Eler et al., Nature **486**, 509 (2012), other information from <https://www.nndc.bnl.gov/nudat2/help/index.jsp>

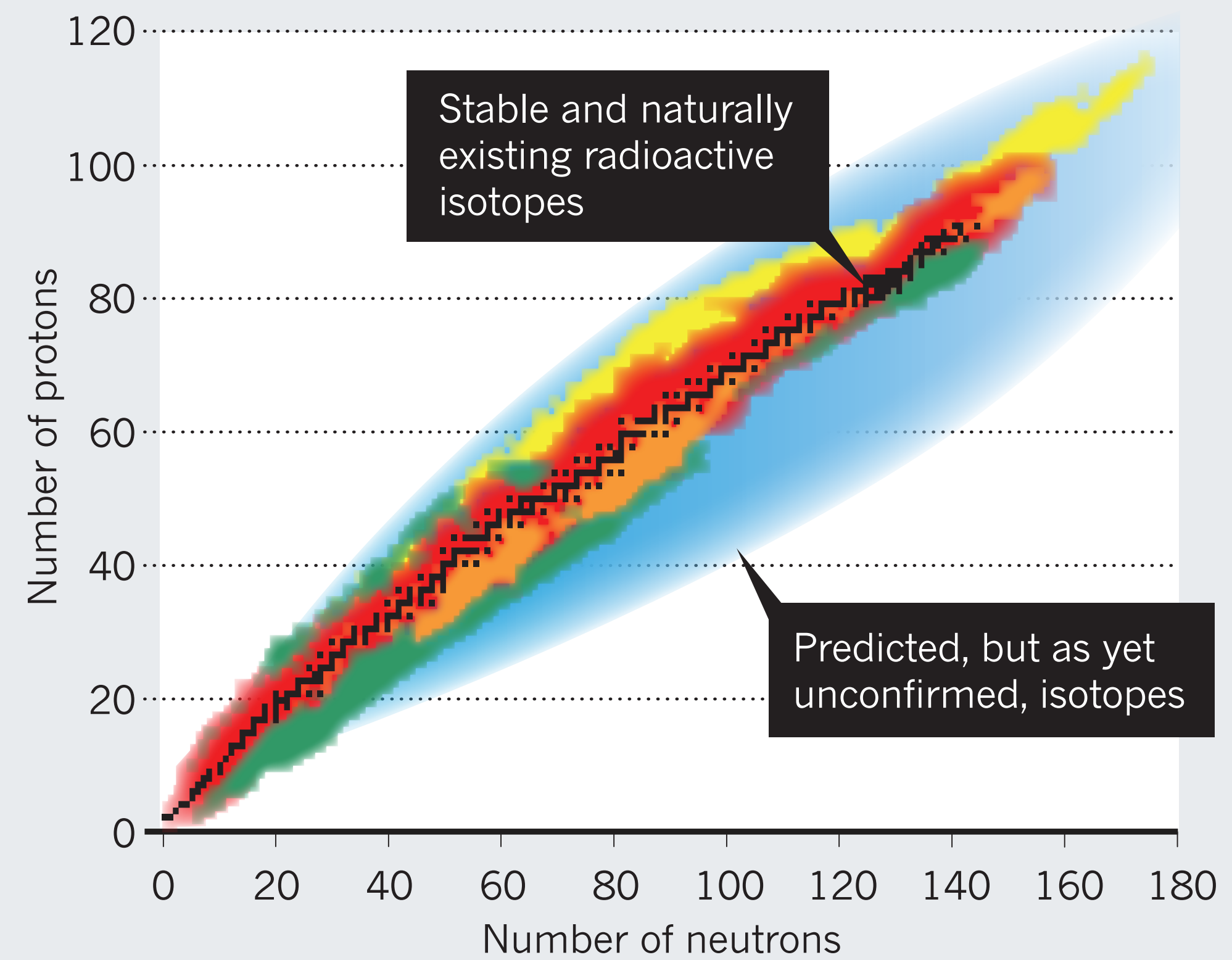
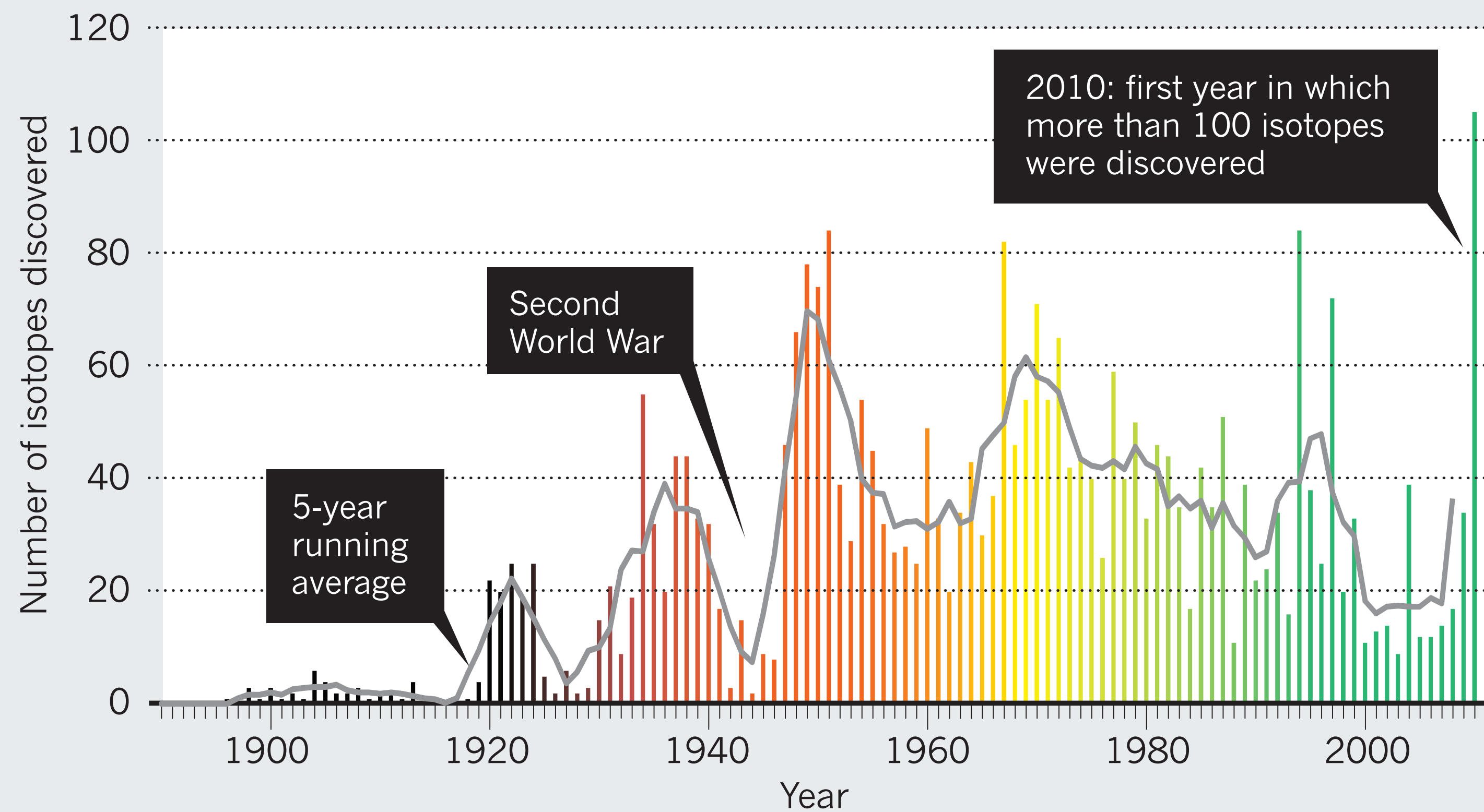
Isotopes

THE NUCLIDE TRAIL

Isotope discovery over the past 100 years (below) has jumped with each introduction of new technology. Some 2,700 radioactive isotopes have been discovered so far (below right), but about 3,000 more are predicted to exist.

Isotope-discovery technique

- Light particle reactions
- Neutron reactions
- Fusion
- Fragmentation/spallation



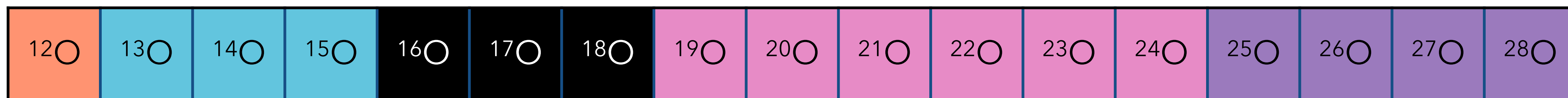
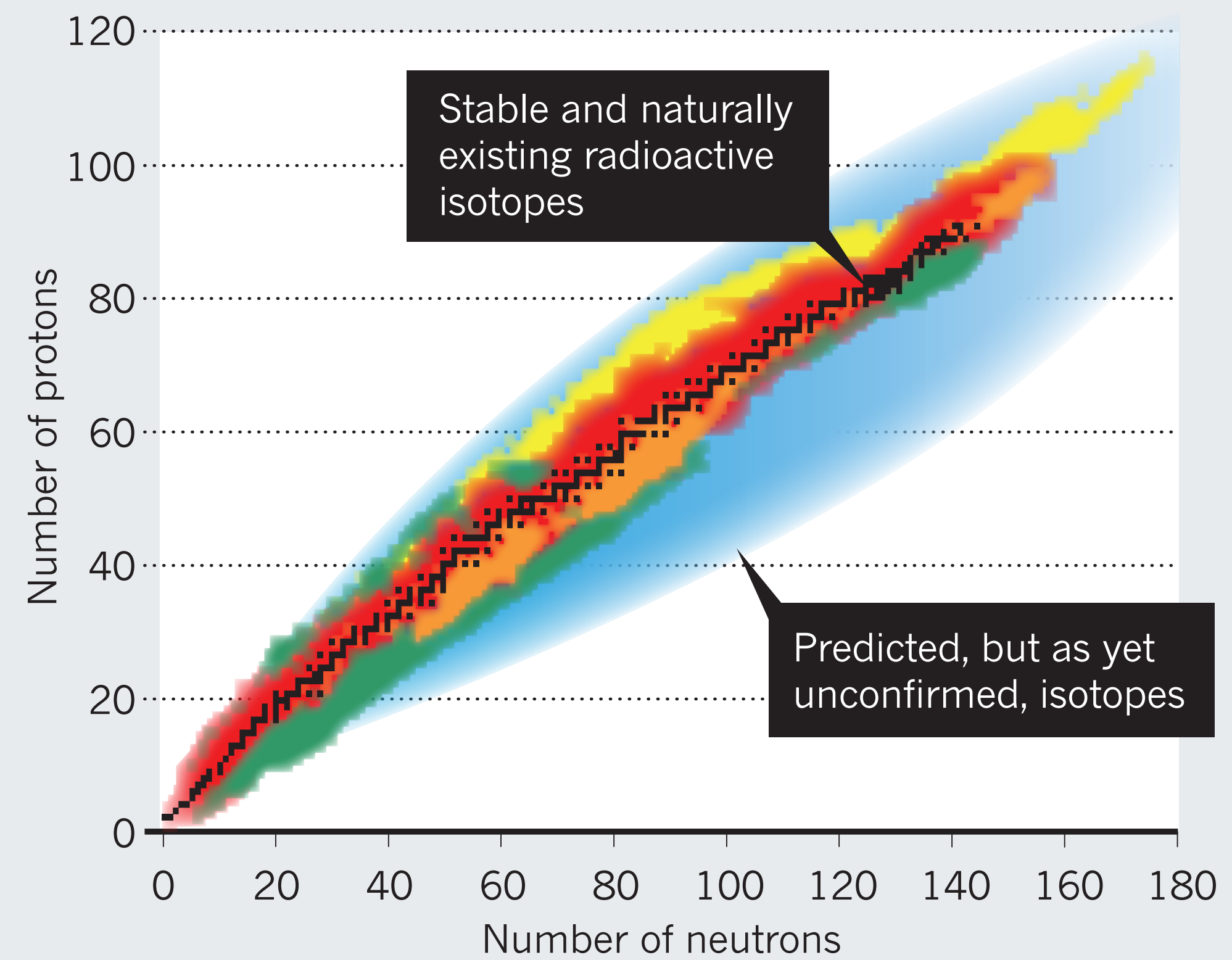
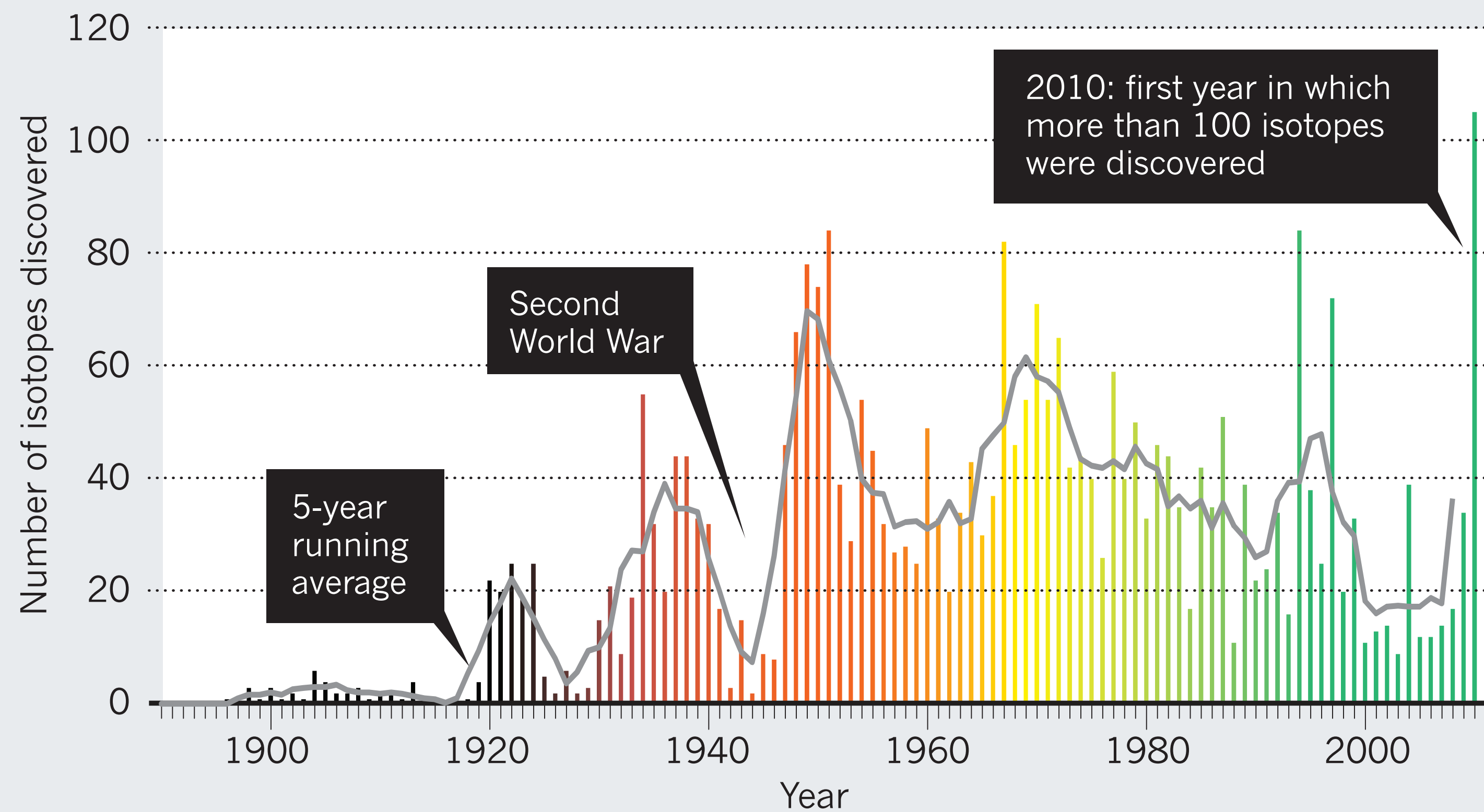
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M. Thoennessen and B. Sherrill, *Nature* **473**, 25 (2011)

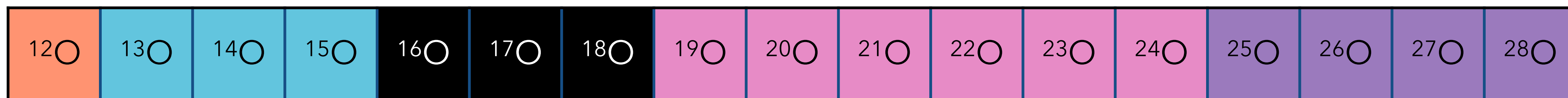
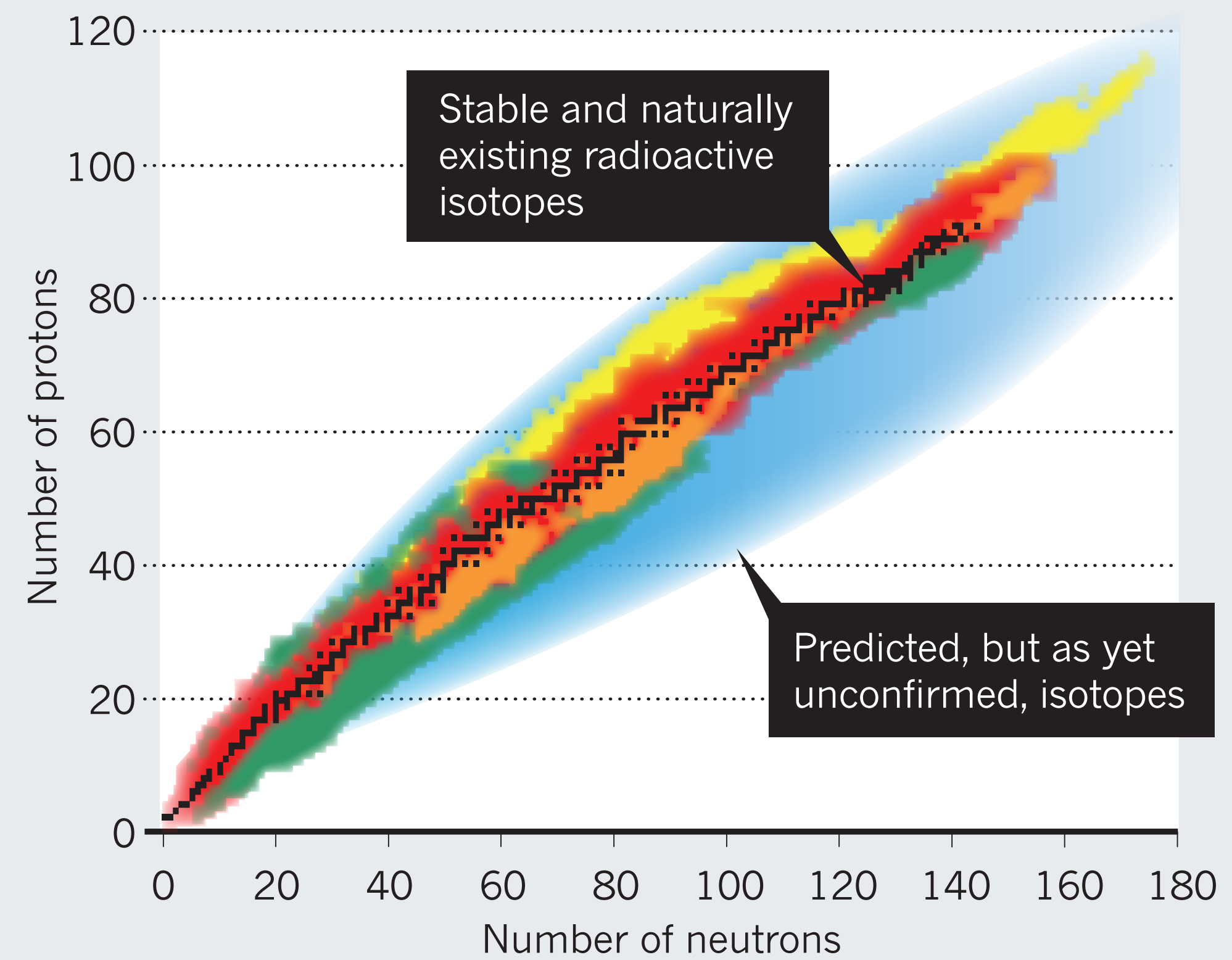
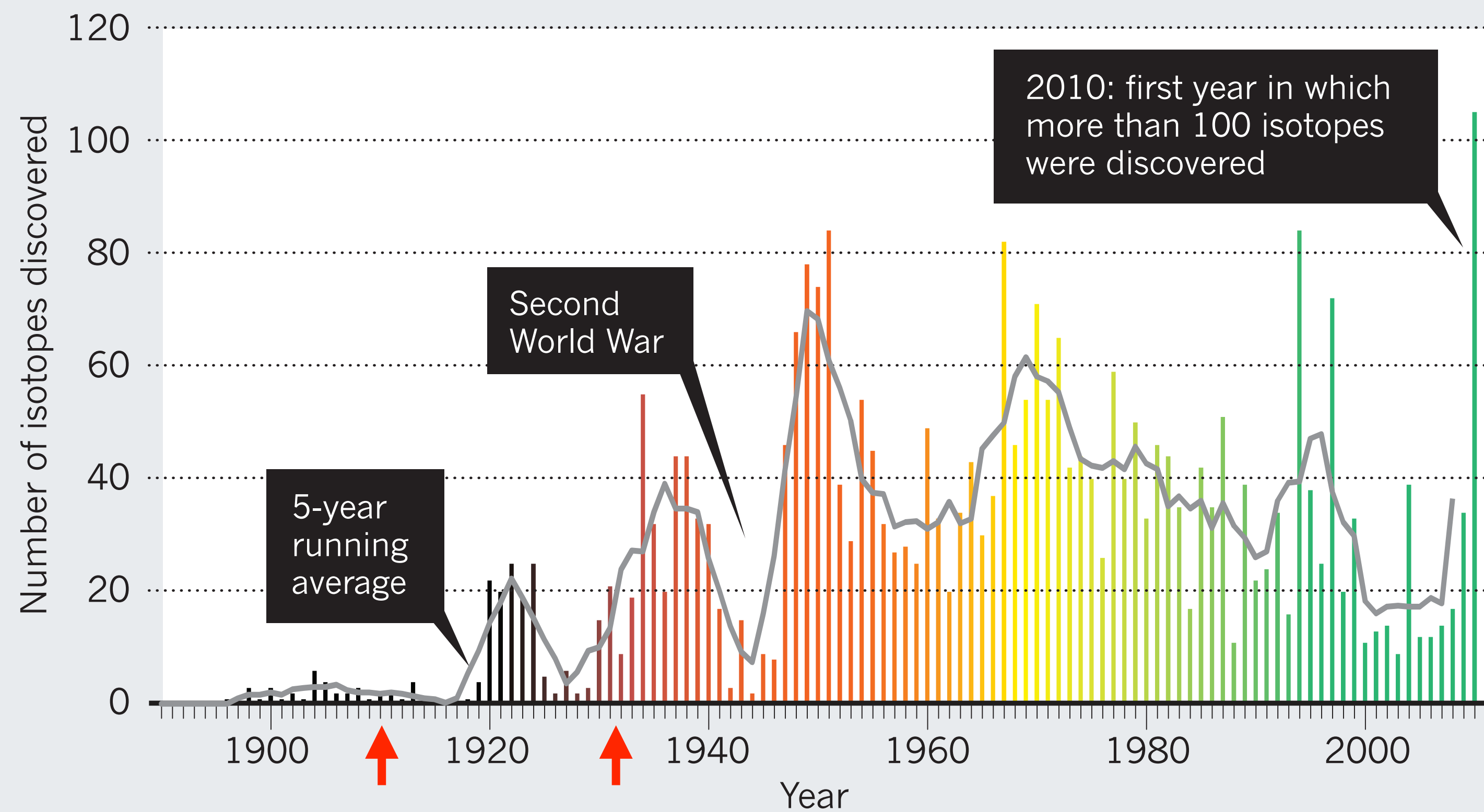
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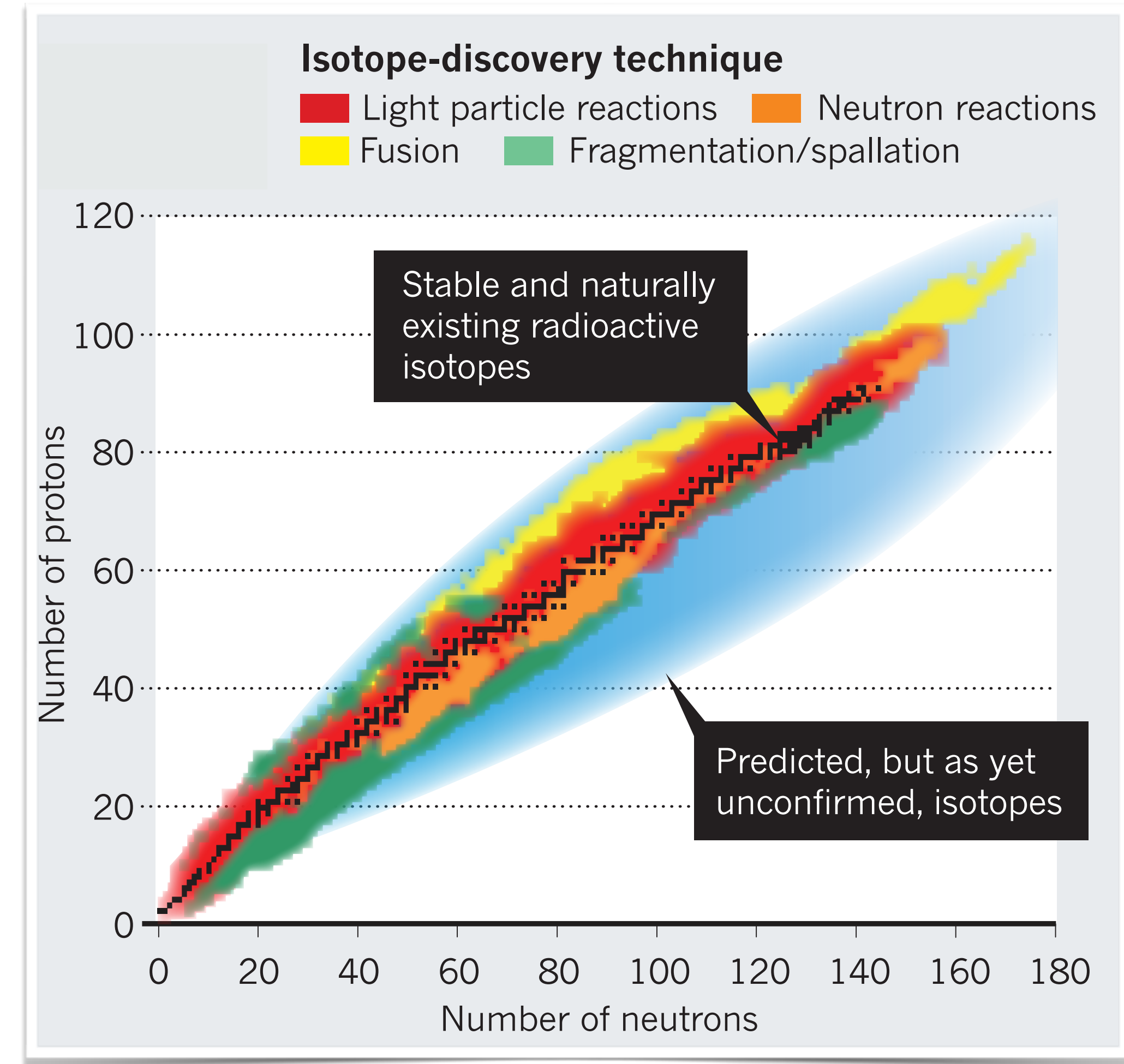
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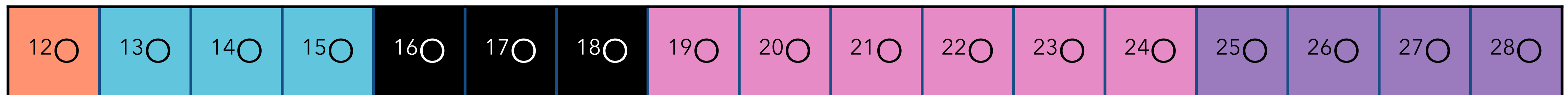
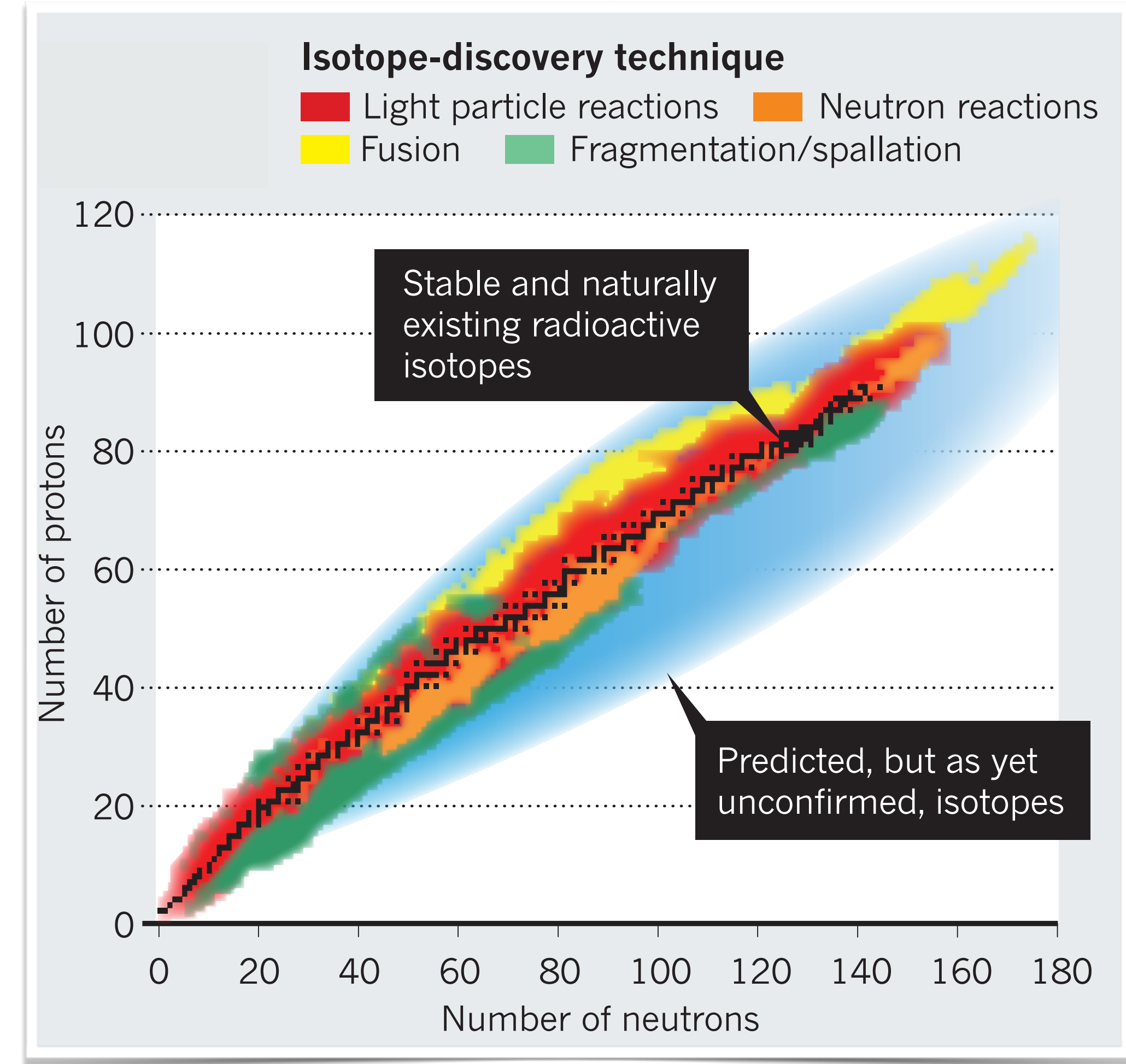
Isotopes

First challenge for an experimentalist is to make/probe the nucleus you want to study ...



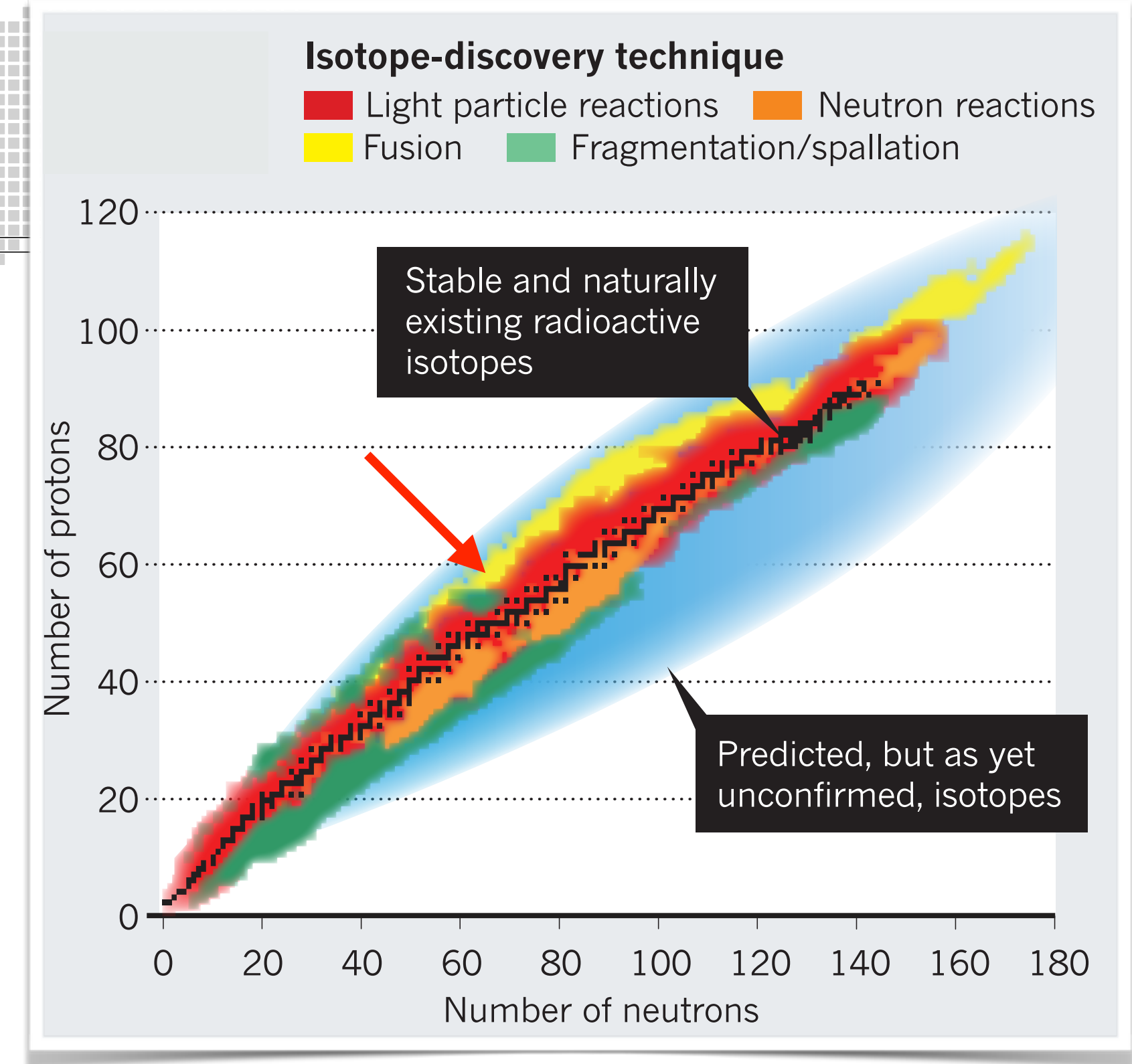
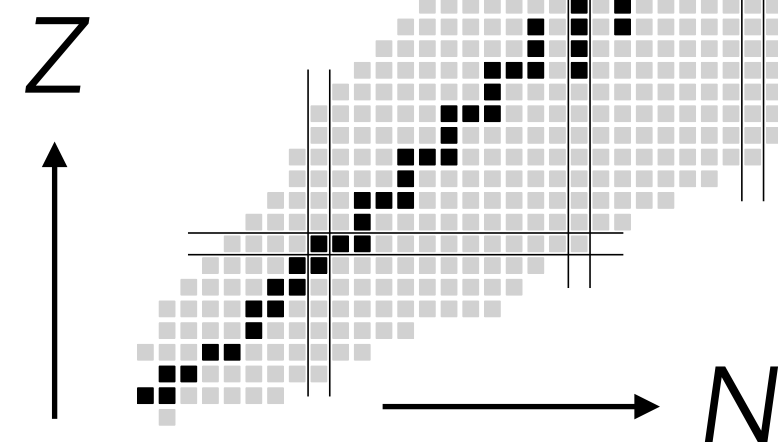
Isotopes

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E.g. fusion(-evaporation)

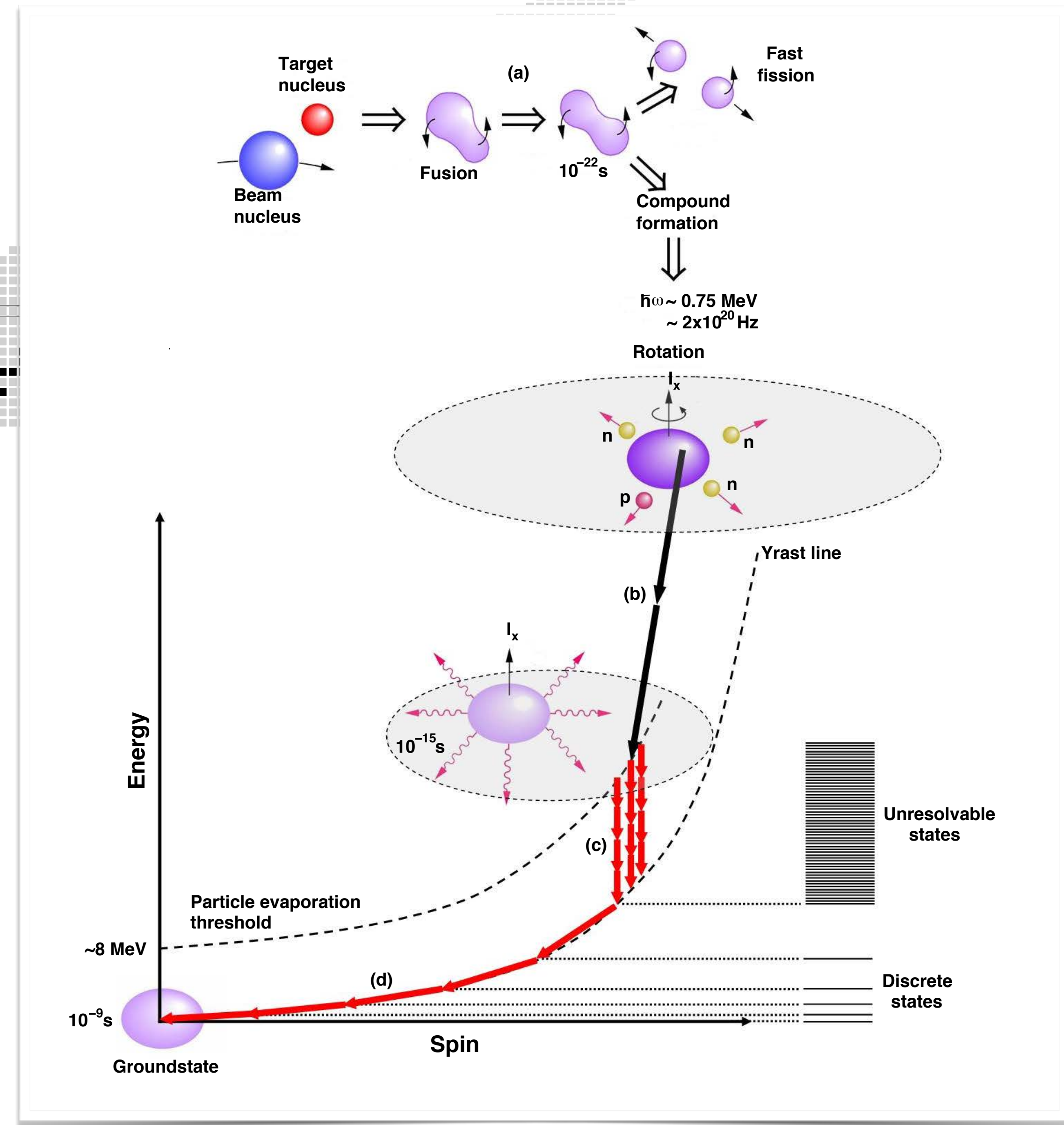
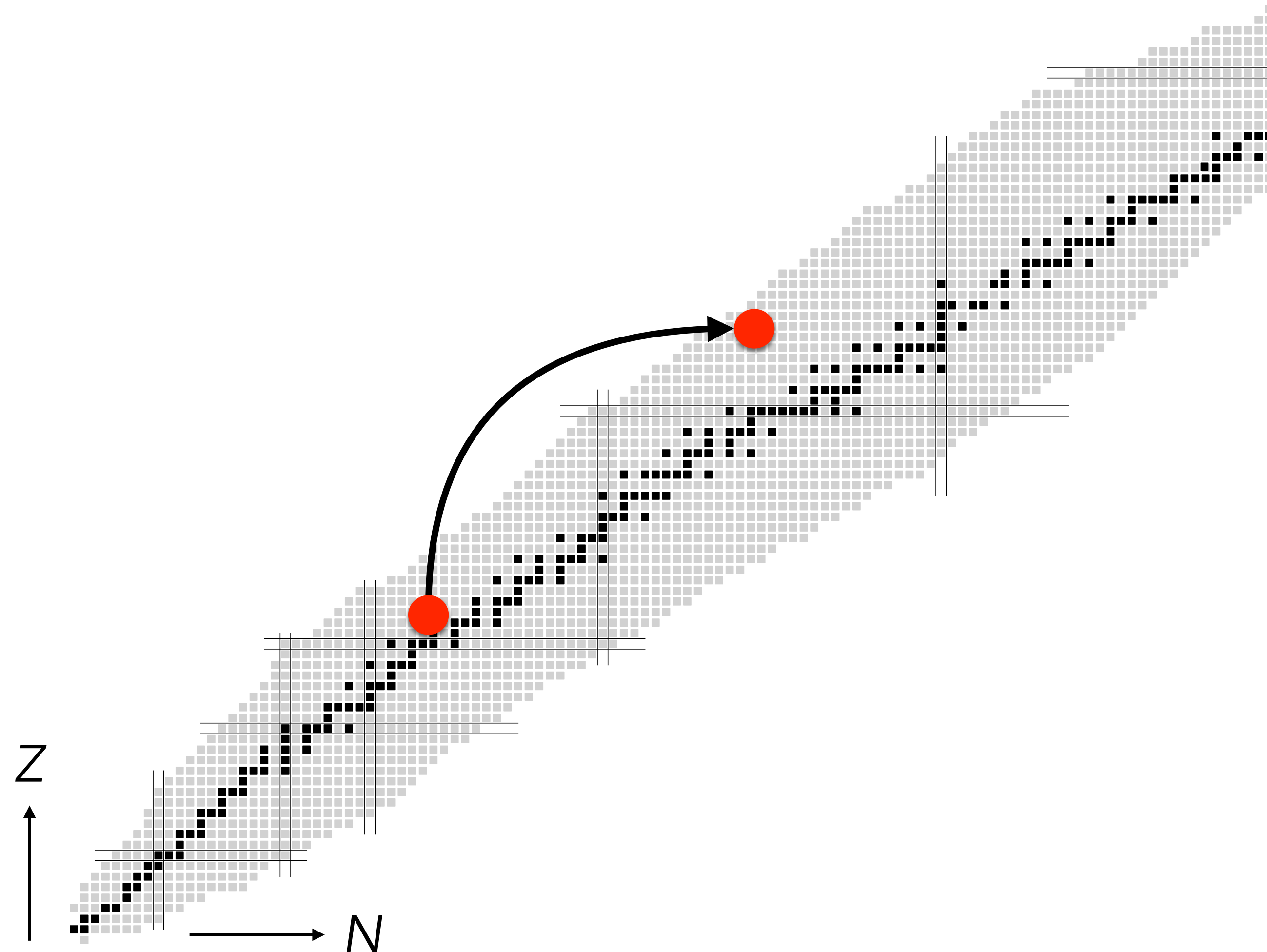
^{122}Ce discovered via $^{64}\text{Zn} + ^{64}\text{Zn} \rightarrow ^{122}\text{Ce} + \alpha + 2n$



J. F. Smith et al. *PLB* **625**, 203 (2005), Bark et al., *Nucl. Phys. A* **514**, 503 (1990) (and A. N. Deacon's thesis, Manchester 2006)

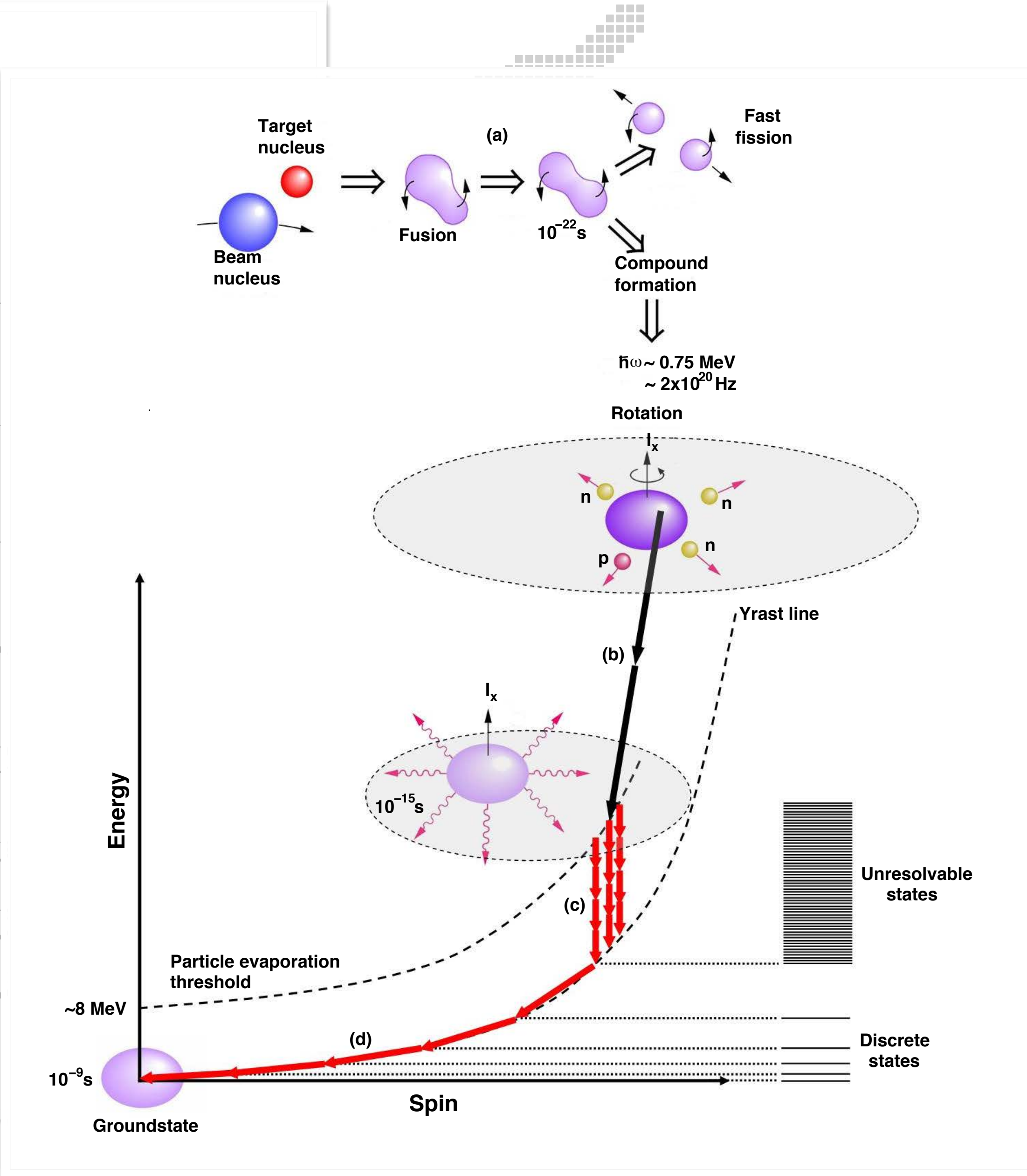
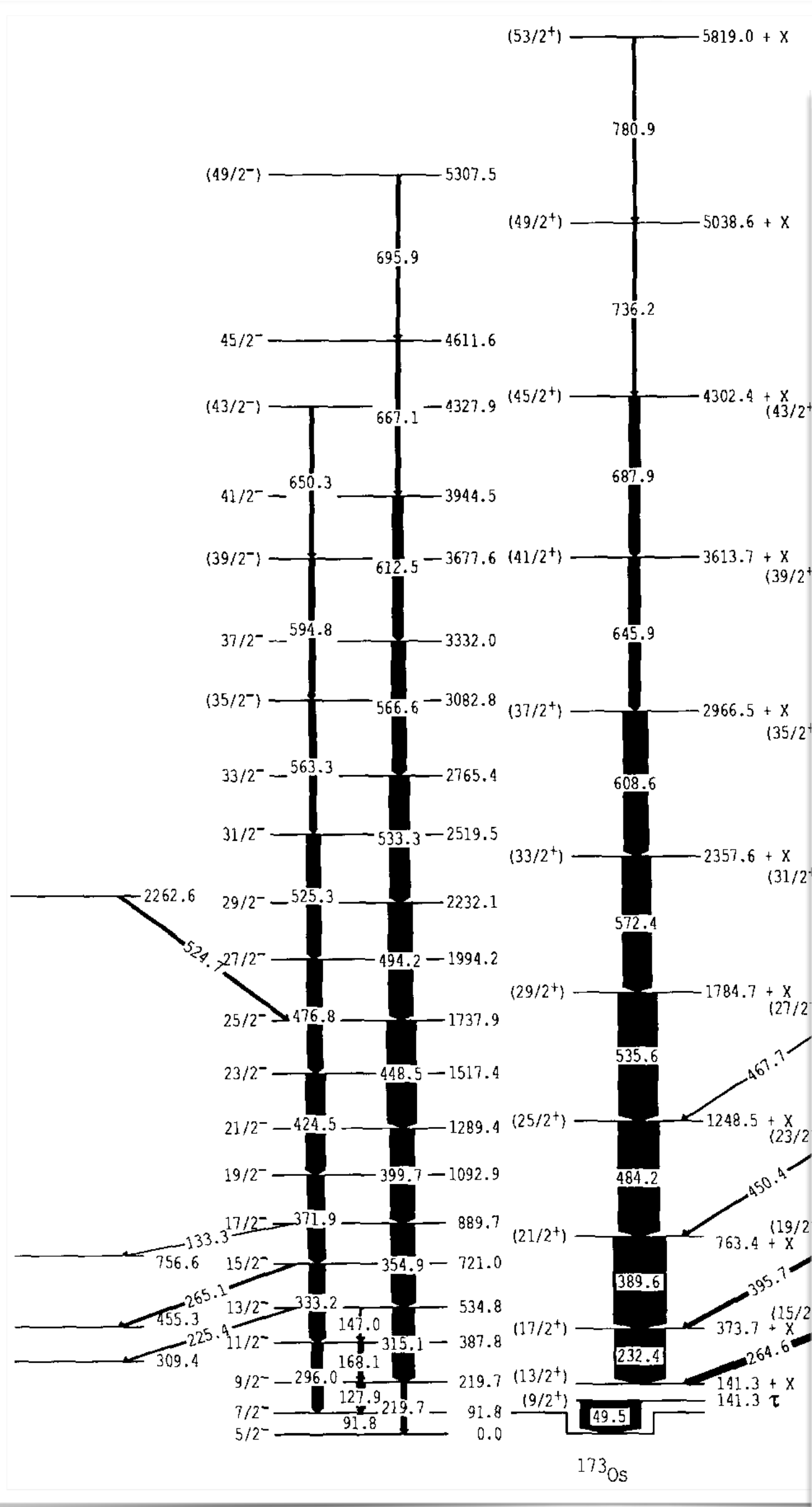
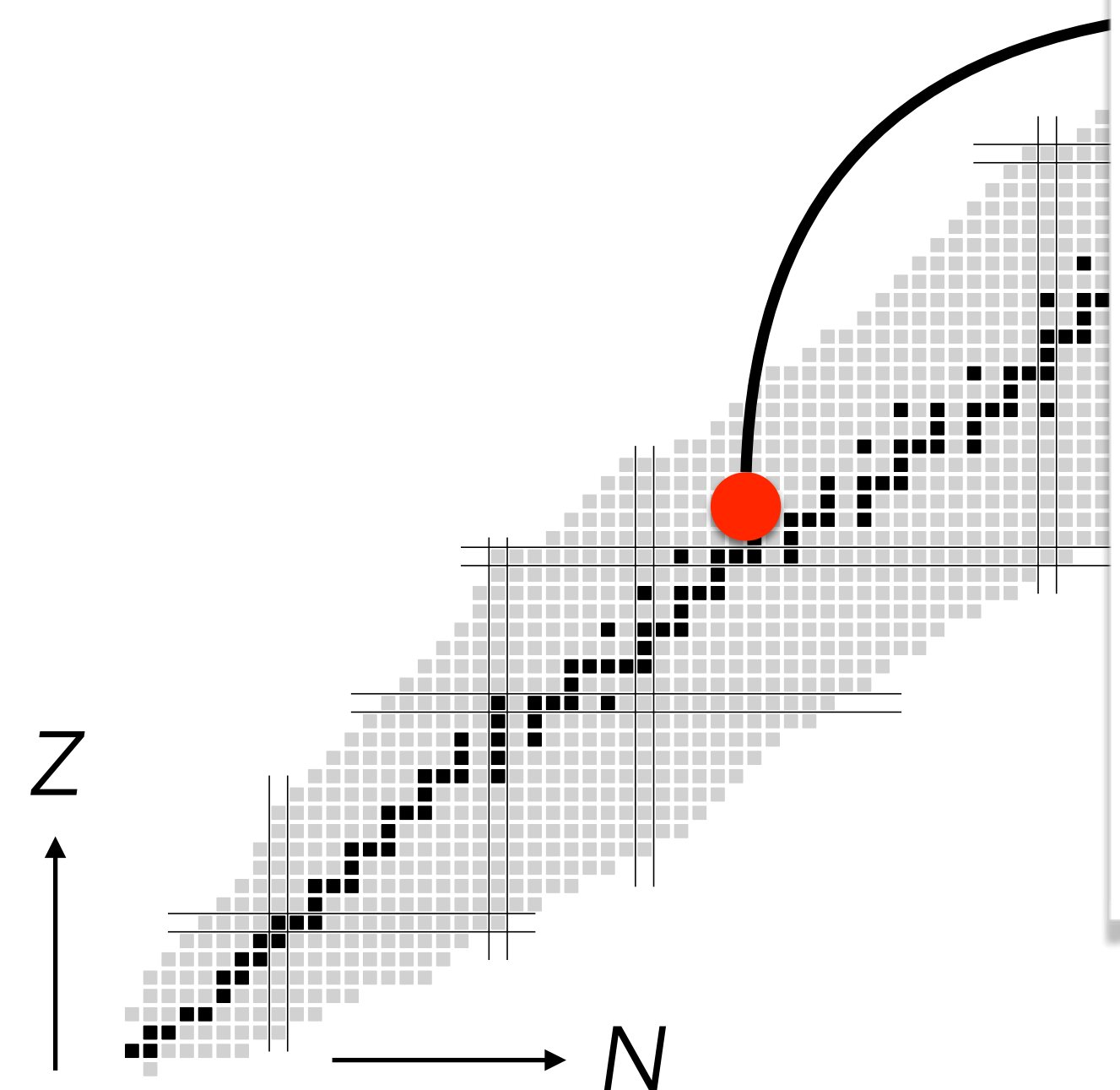
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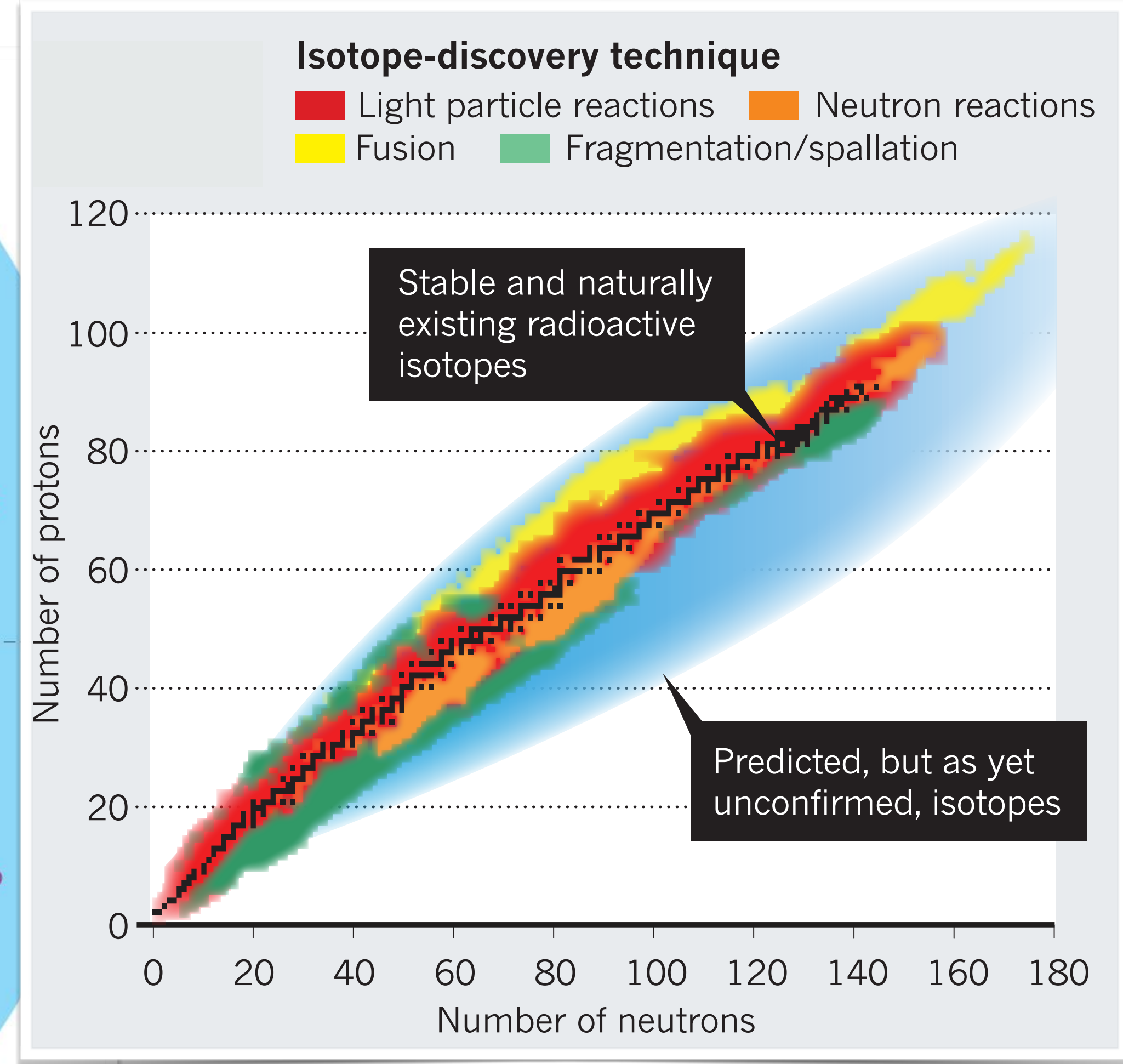
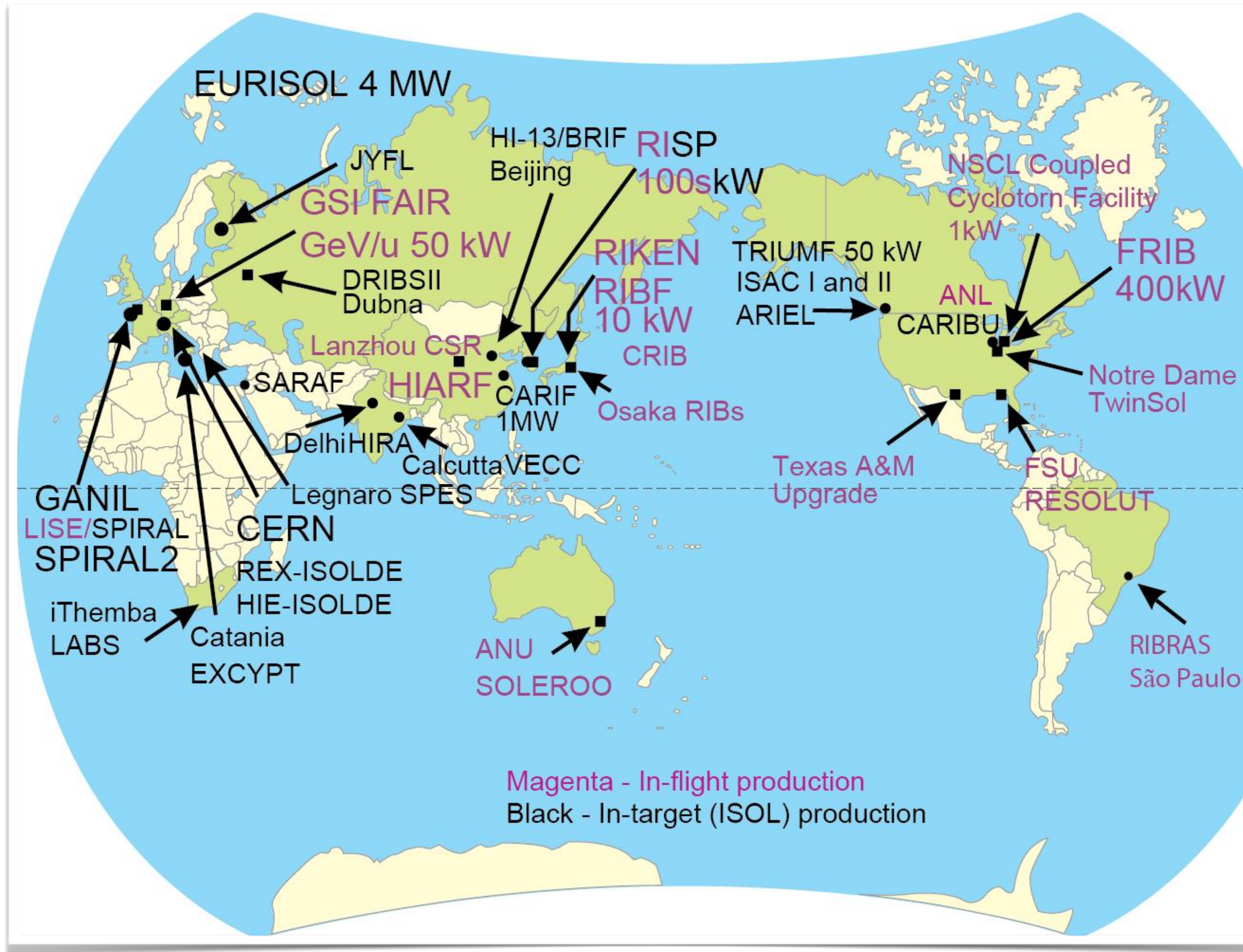


E.g. fusion(-ev)

^{122}Ce discovered via ^{64}Zn



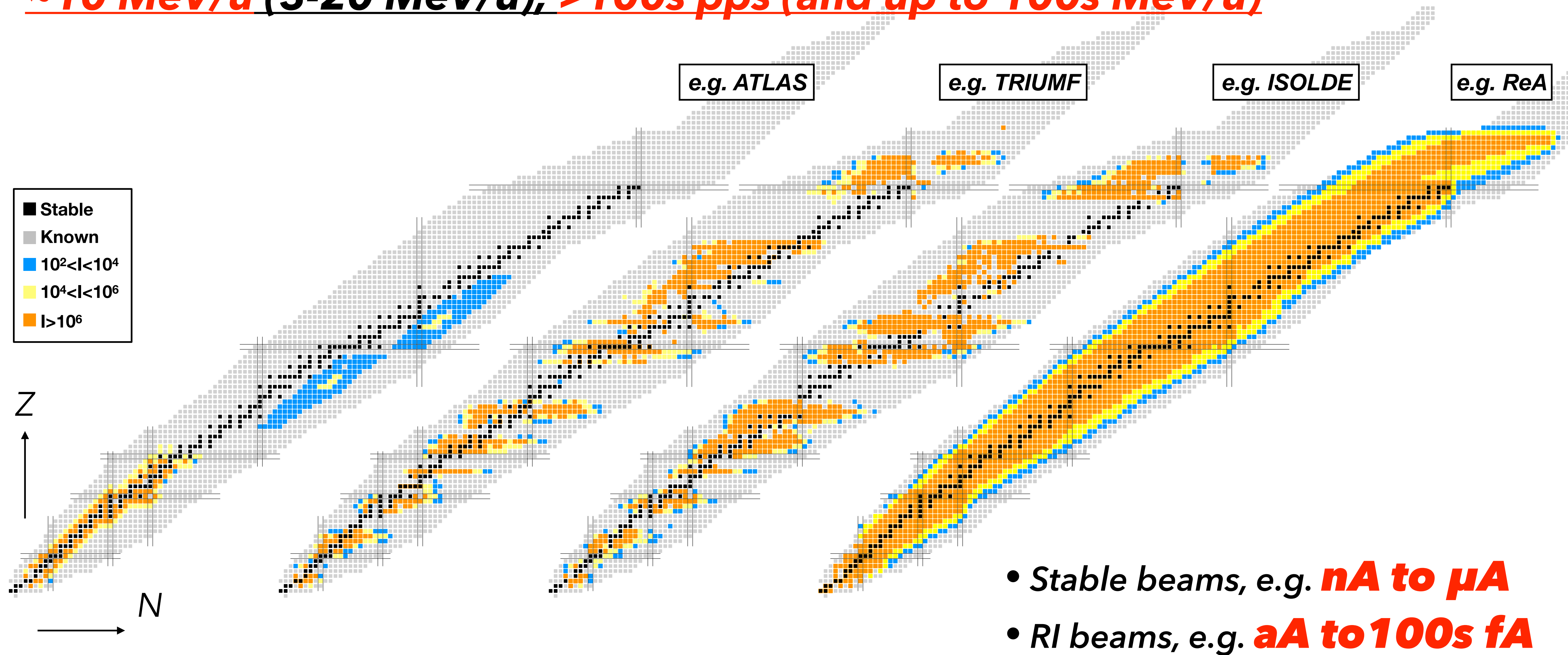
Radioactive ion beam facilities



(Left) Original source unsure (perhaps Brad Sherrill)

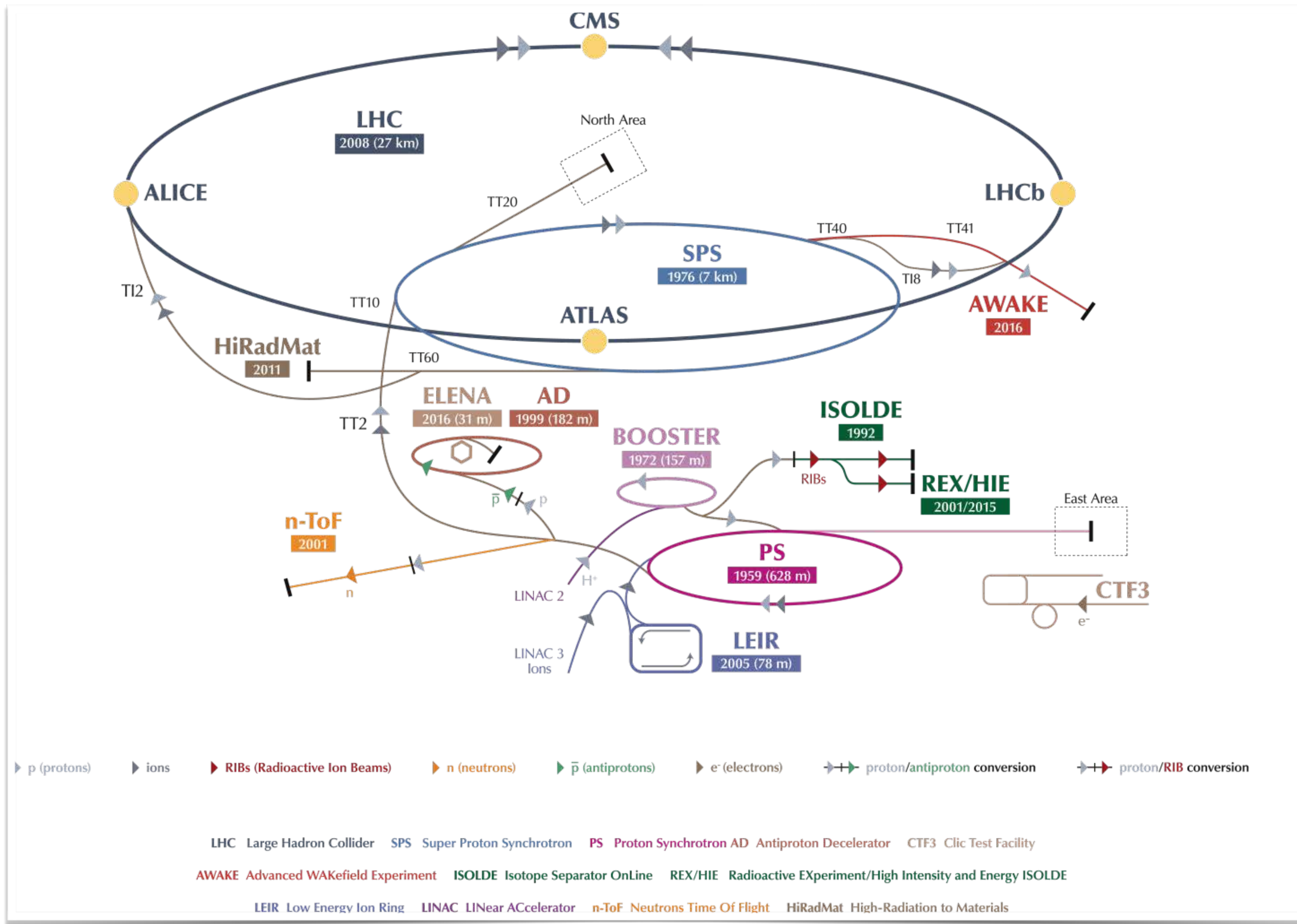
... RI beams

$\sim 10 \text{ MeV/u}$ (3-20 MeV/u), $> 100\text{s pps}$ (and up to 100s MeV/u)



(Beam rates are very crude estimates from various sources, illustrative, likely $\sim 1\text{-}2$ orders of mag. off)

RIB facilities, ISOL at e.g. CERN, TRIUMF, ...



Examples of ISOL facilities:

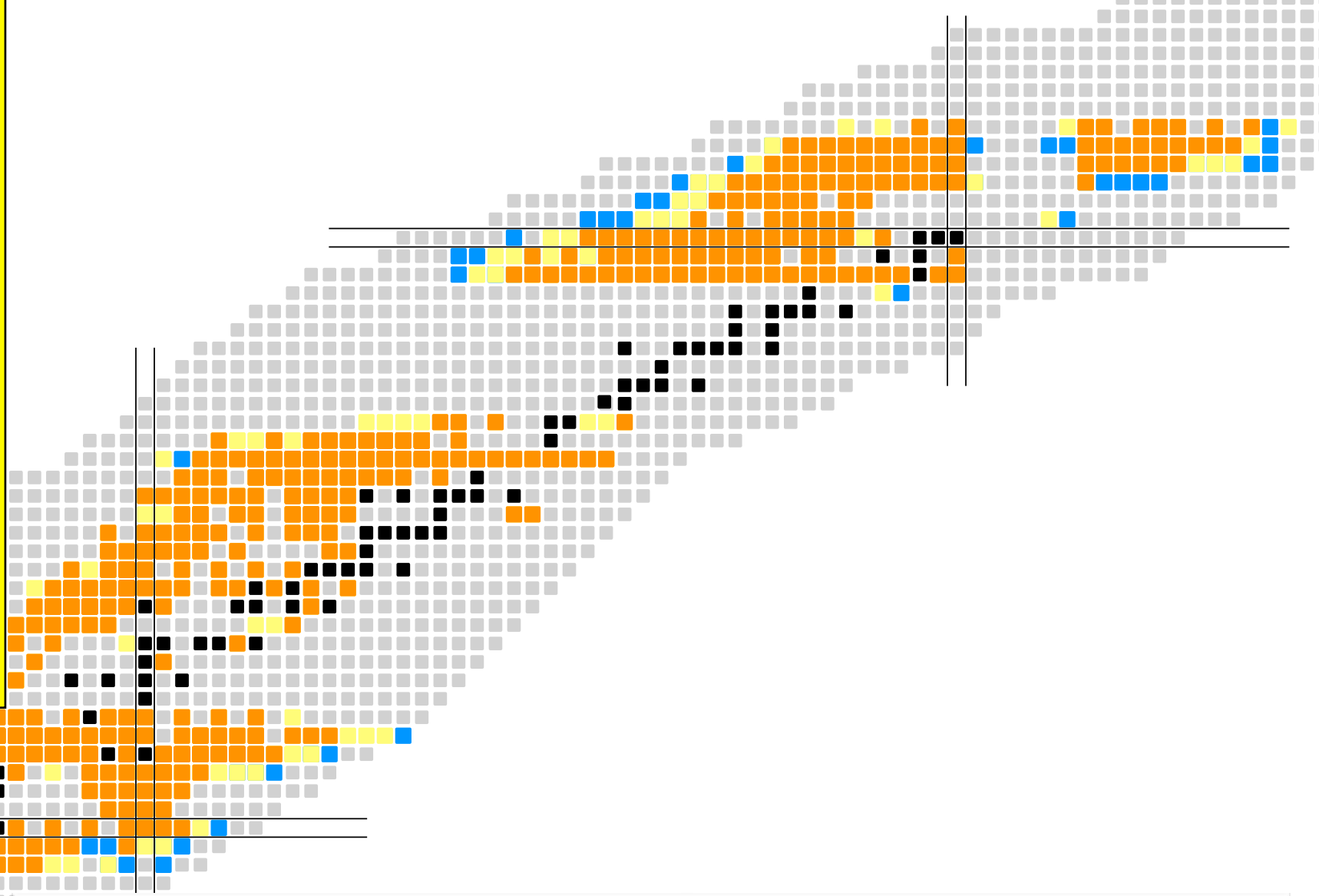
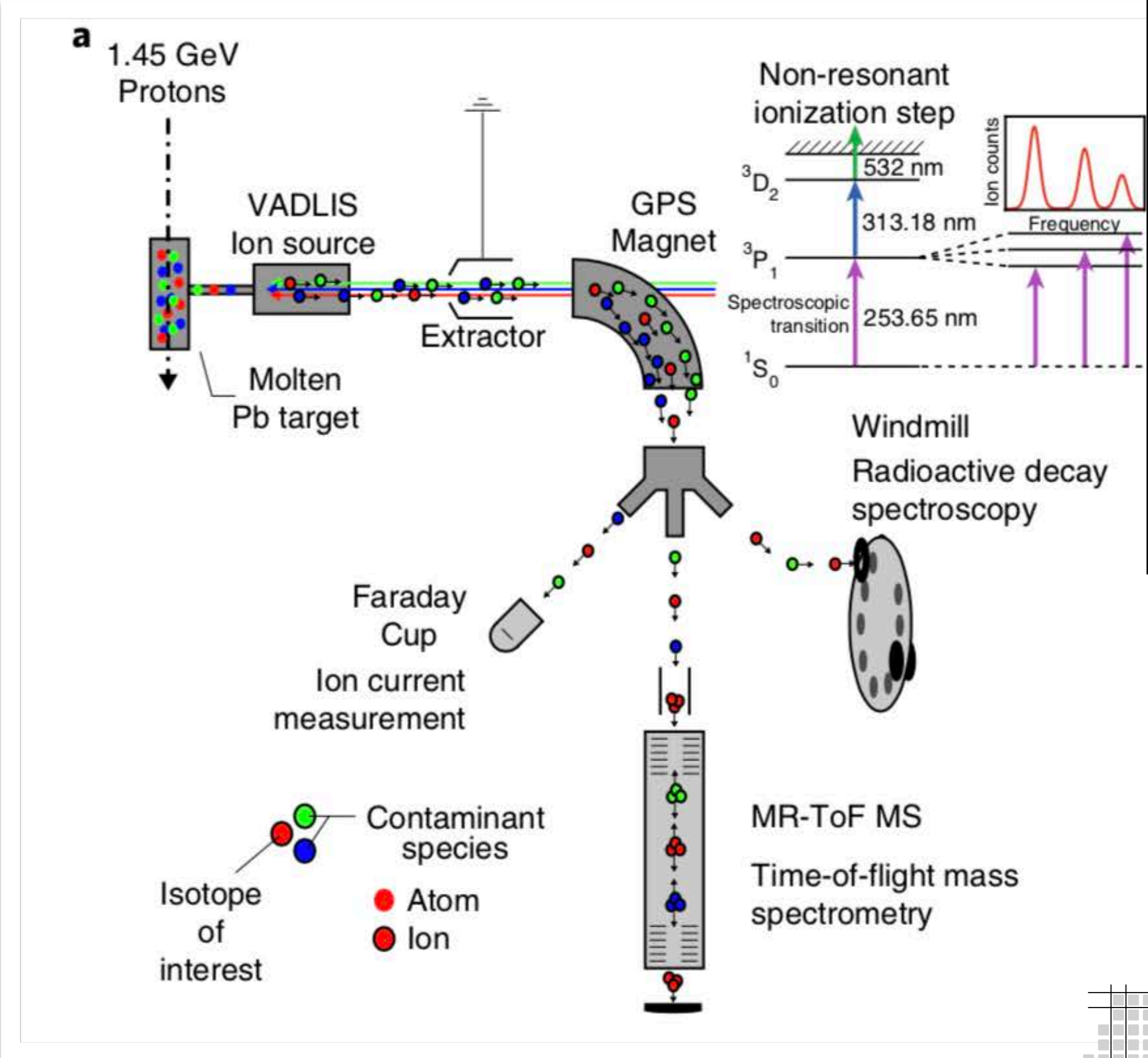
- TRIUMF (Canada)
- SPIRAL/SPIRAL2 (France)
- REX-ISOLDE/HIE-ISOLDE (CERN)
- iTHEMBA - future radioactive-beam facility (South Africa)
- JYFL (Finland) - IGISOL

Typically:

- *Light, energetic ion*
- *Thick, hot target*
- *Ion source*
- *Separator*
- *Post acceleration*
- *Lower beam energies*

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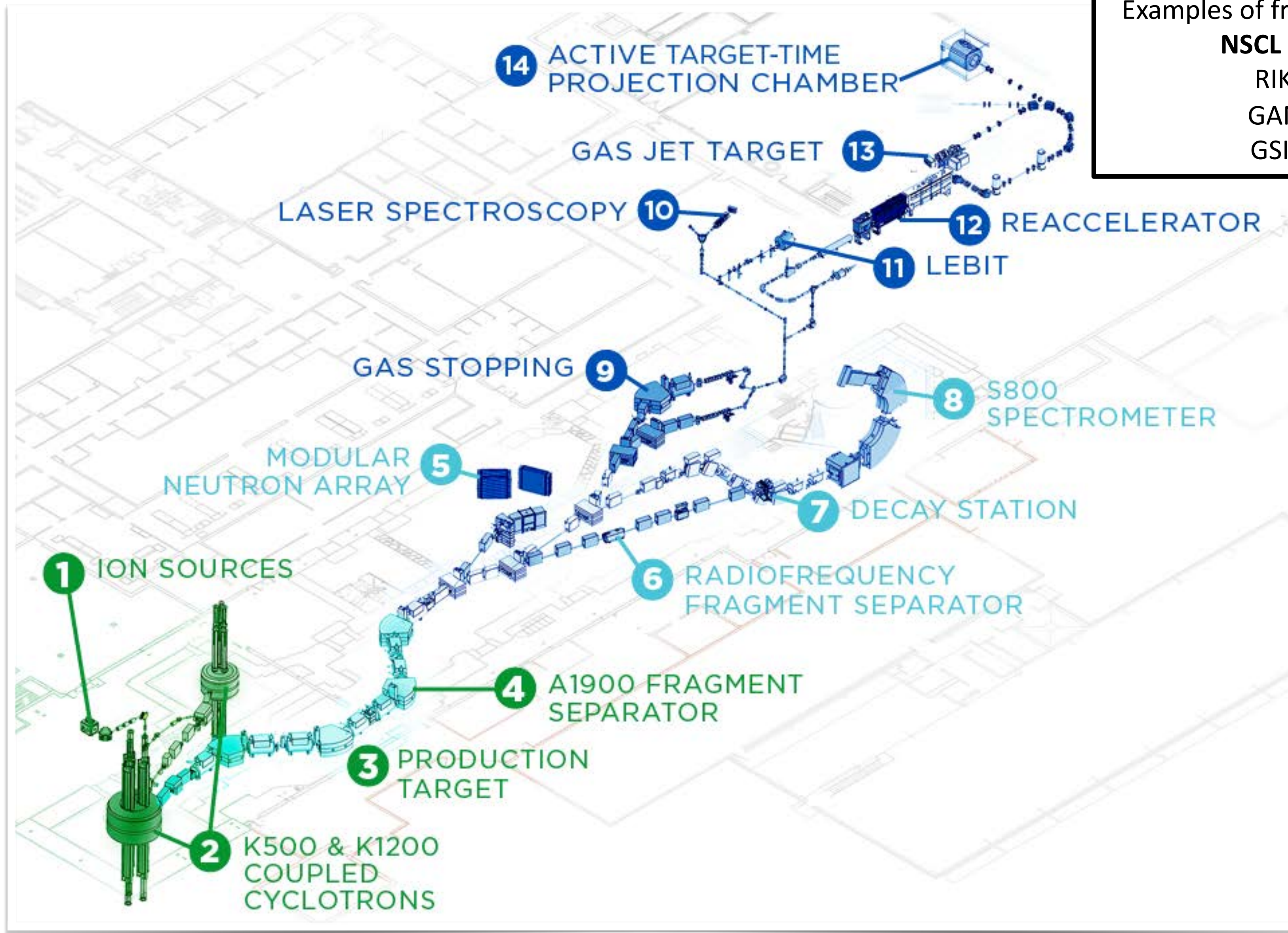
Target Elements

e.g. ISOL targets

H	He																	He																													
Li	Be	B	C	N	O	F	Ne																	Ne																							
Na	Mg	Al	Si	P	S	Cl	Ar																	Ar																							
K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr																	Kr													
Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe																	Xe													
Cs	Ba	La	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn																	Rn													
Fr	Ra	Ac	Rf	Db	Sg	Bh	Hs	Mt	Ds	Rg	Cn	Nh	Fl	Mc	Lv	Ts	Og																	Og													
																		Ce	Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm	Yb	Lu																
																		Th	Pa	U	Np	Pu	Am	Cm	Bk	Cf	Es	Fm	Md	No	Lr																

(Beam rates are very crude estimates from various sources, illustrative, likely ~1-2 orders of mag. off)

In-flight / fragmentation at e.g. NSCL



Examples of fragmentation facilities:

NSCL (USA) → FRIB

RIKEN (Japan)

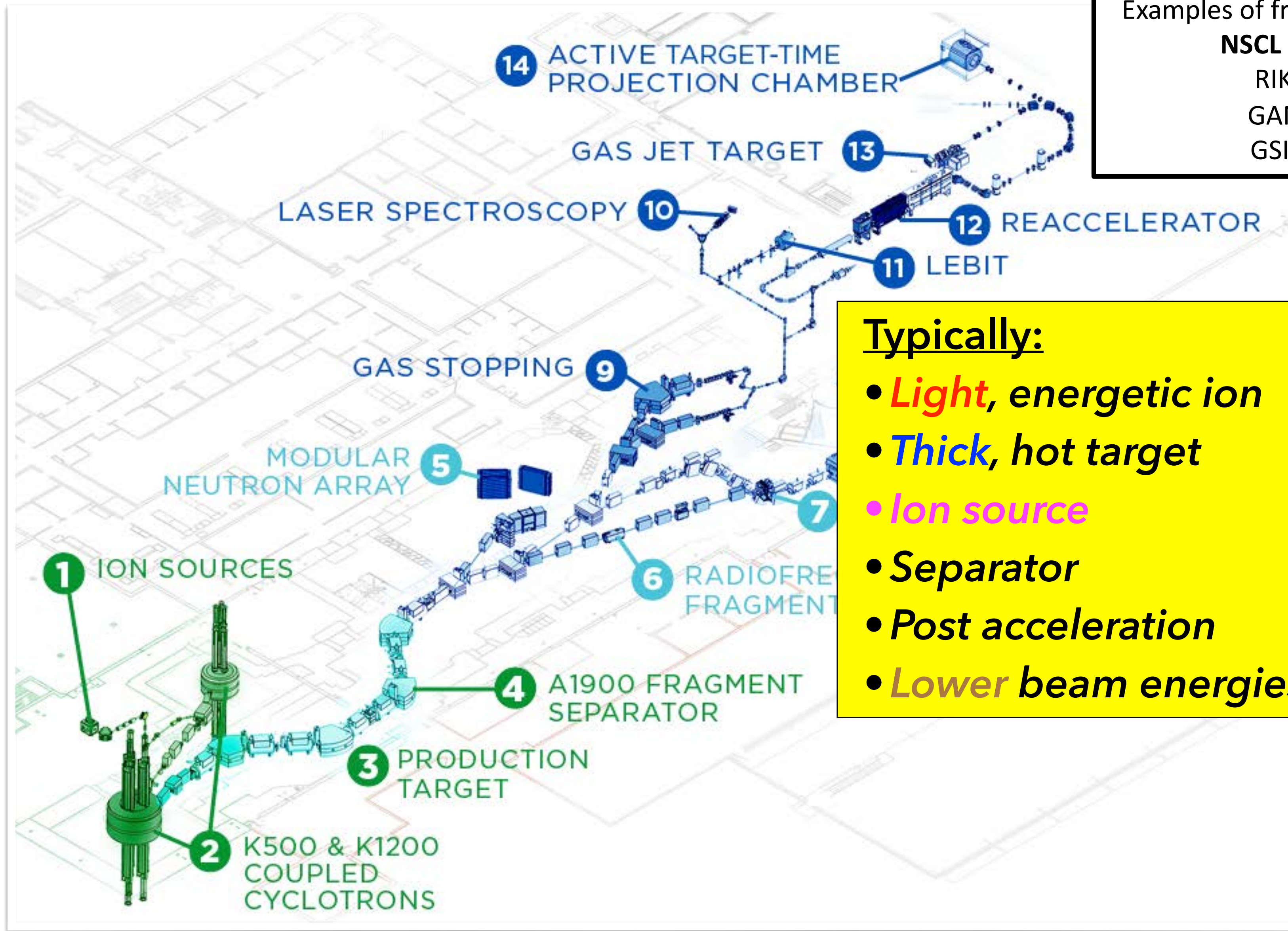
GANIL (France)

GSI (Germany)

Typically:

- *Ion source*
- *heavy, energetic ion*
- *Thin, light target*
- *Separator*
- *Higher beam energies*

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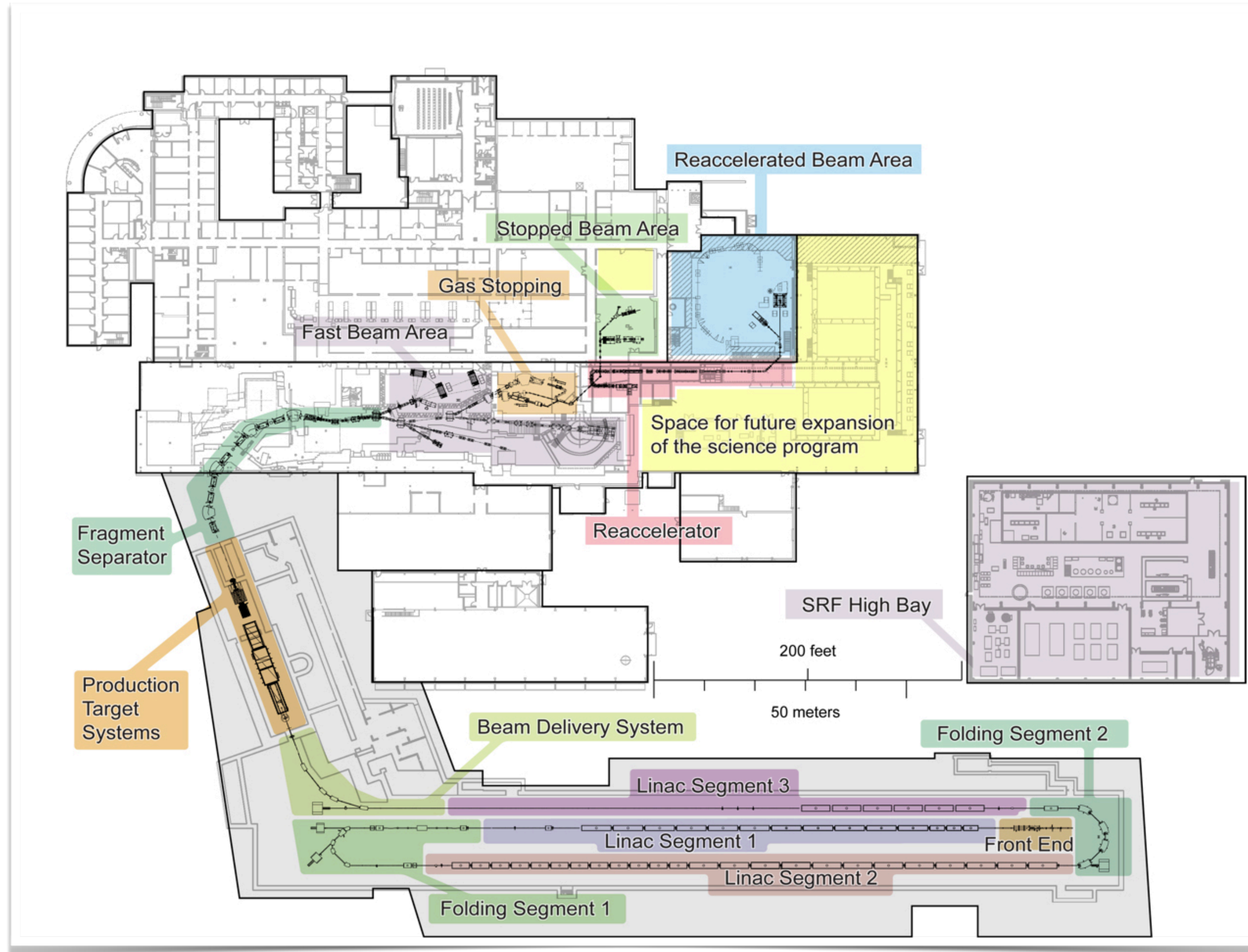
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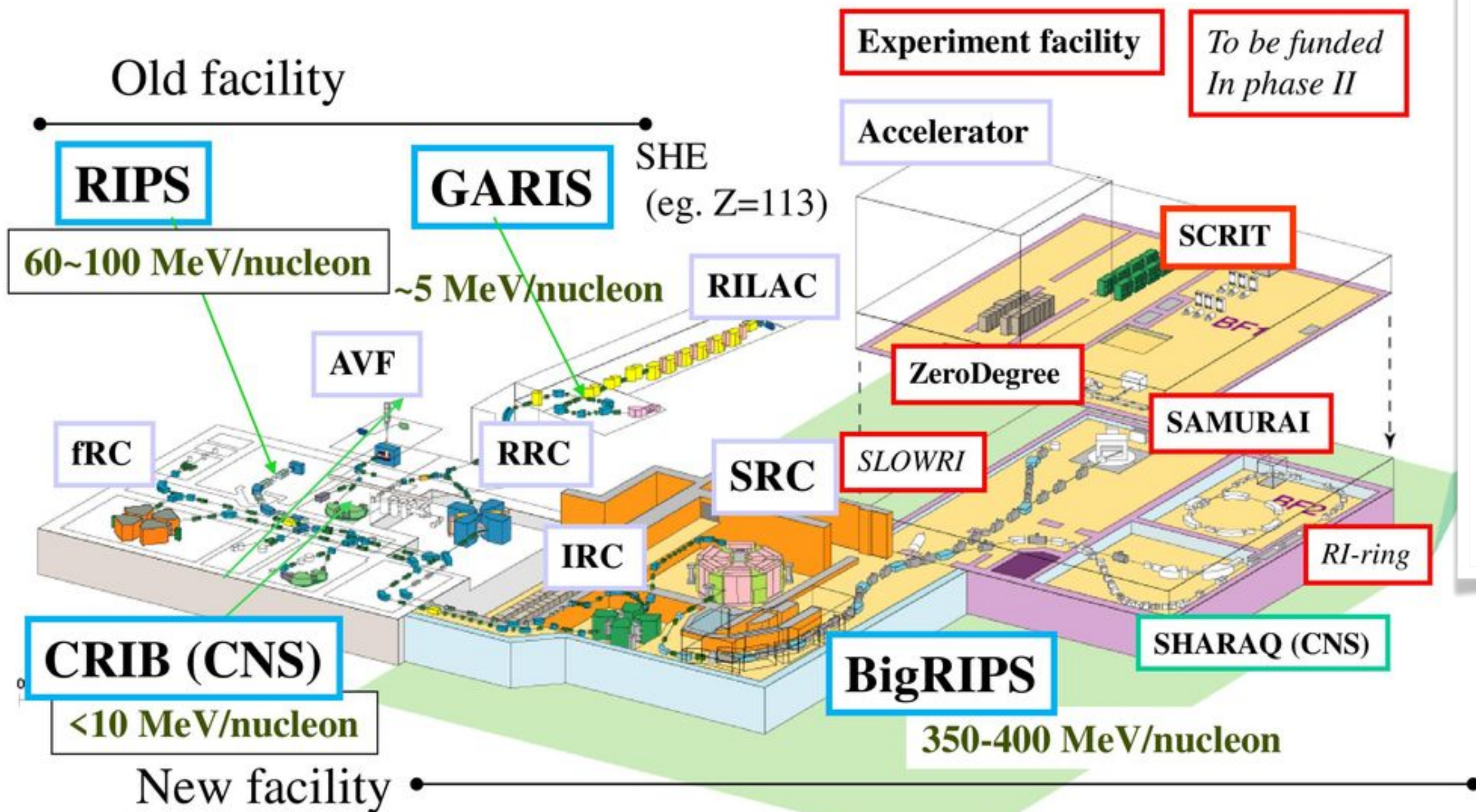


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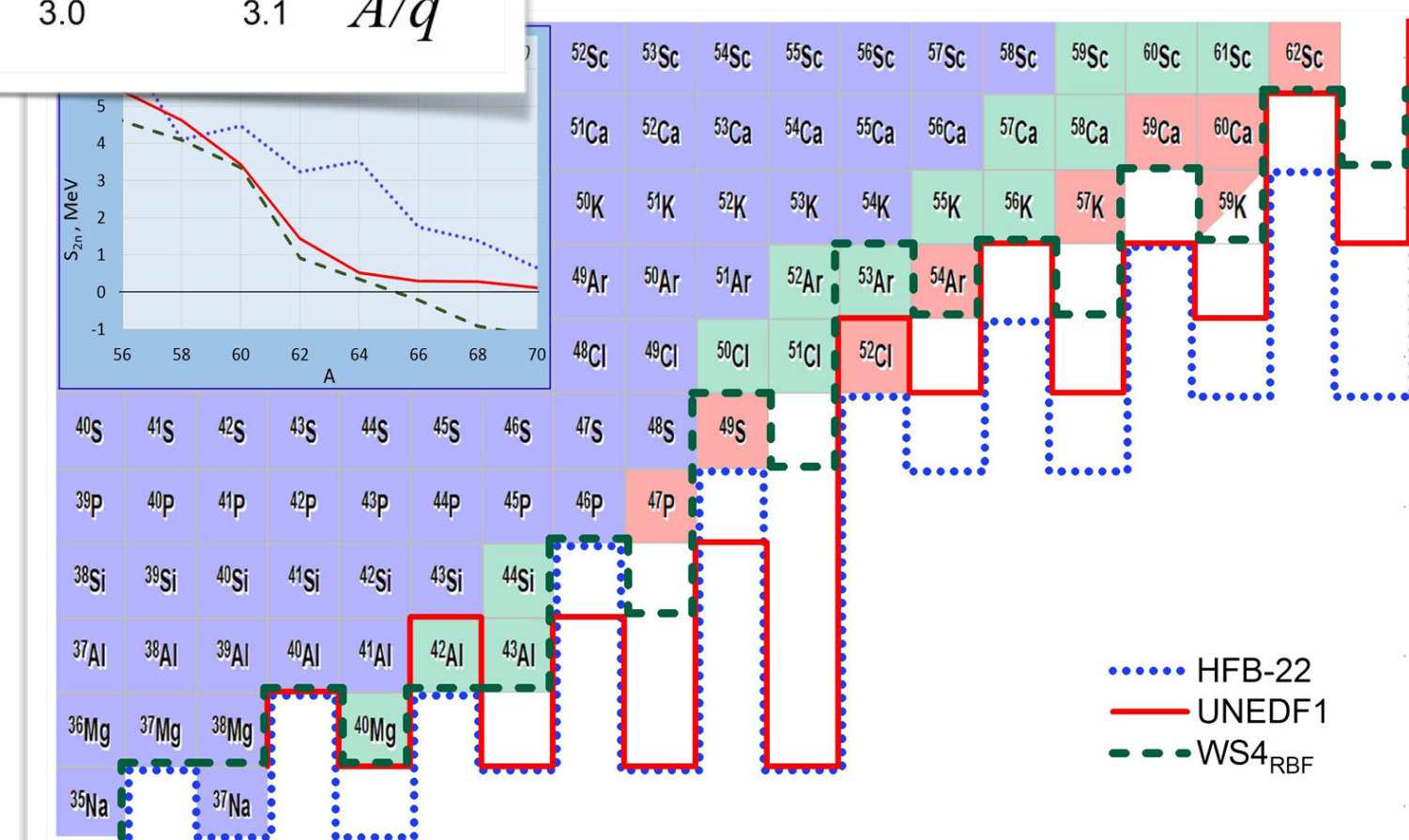
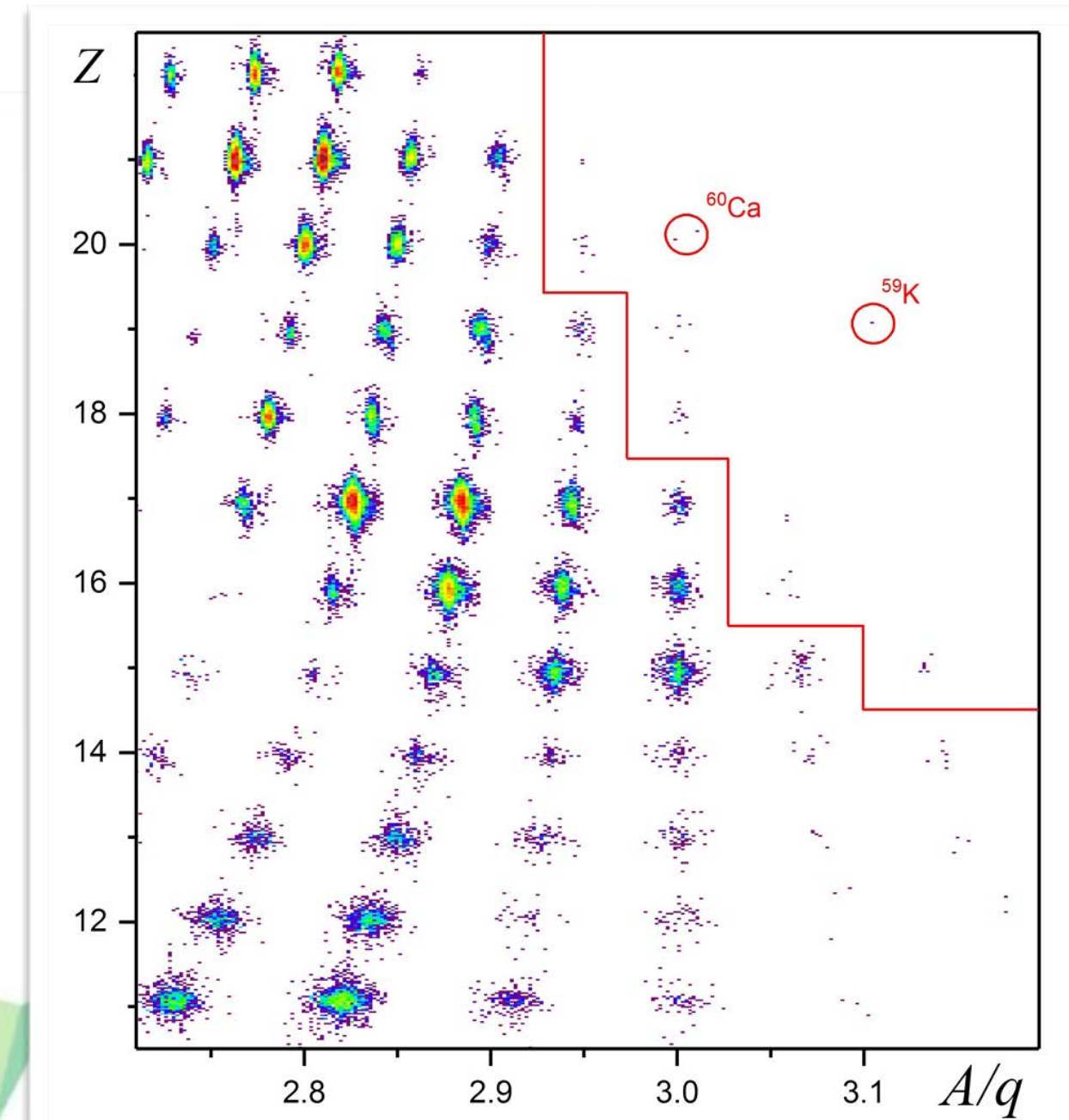
- *Ion source*
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In-flight / fragmentation at e.g. RIBF, RIKEN

RIKEN RI Beam Factory (RIBF)



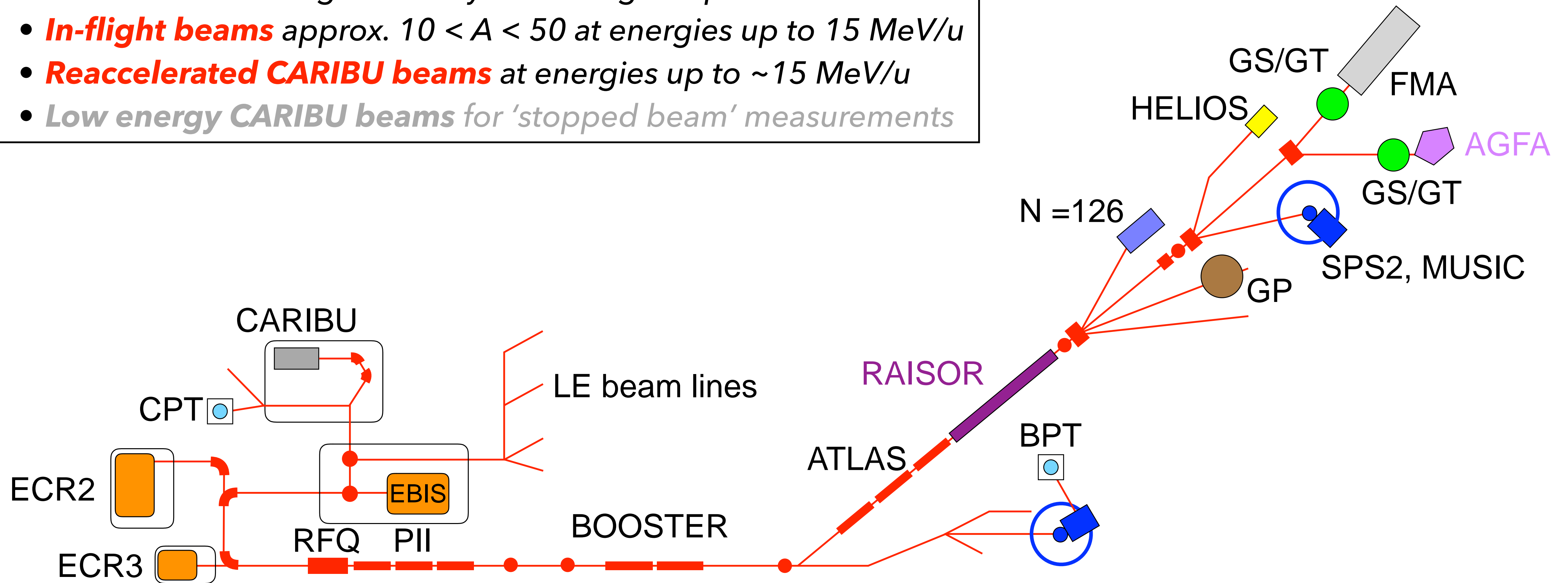
Intense (80 kW max.) H.I. beams (up to U) of 345 A MeV at SRC
 Fast RI beams by projectile fragmentation and U-fission at BigRIPS
 Operation since 2007



E.g. "ISOL" and in-flight at Argonne

An unrivaled combination for direct reaction studies

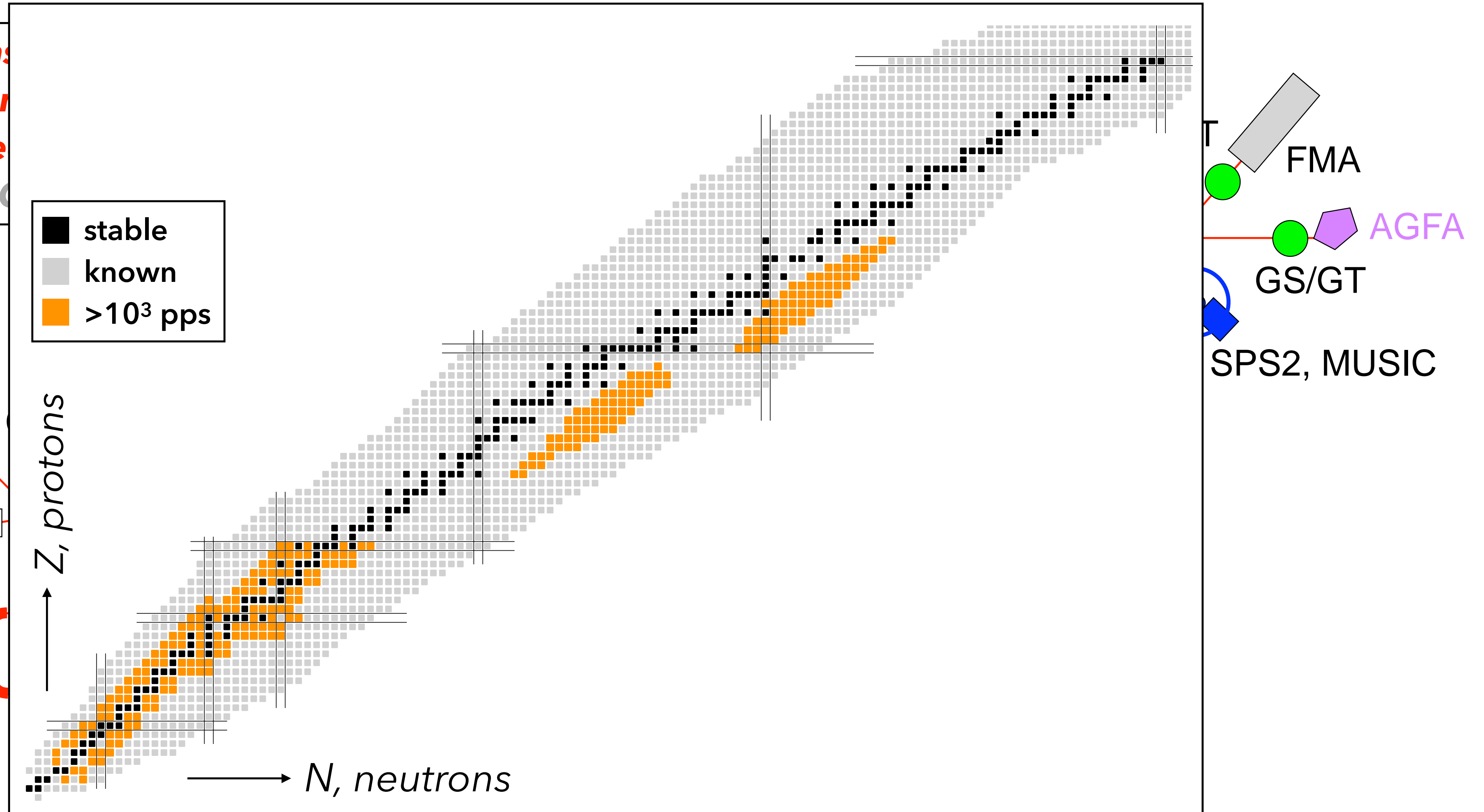
- **Stable beams** at high intensity and energies up to 18 MeV/u
- **In-flight beams** approx. $10 < A < 50$ at energies up to 15 MeV/u
- **Reaccelerated CARIBU beams** at energies up to ~15 MeV/u
- **Low energy CARIBU beams** for 'stopped beam' measurements



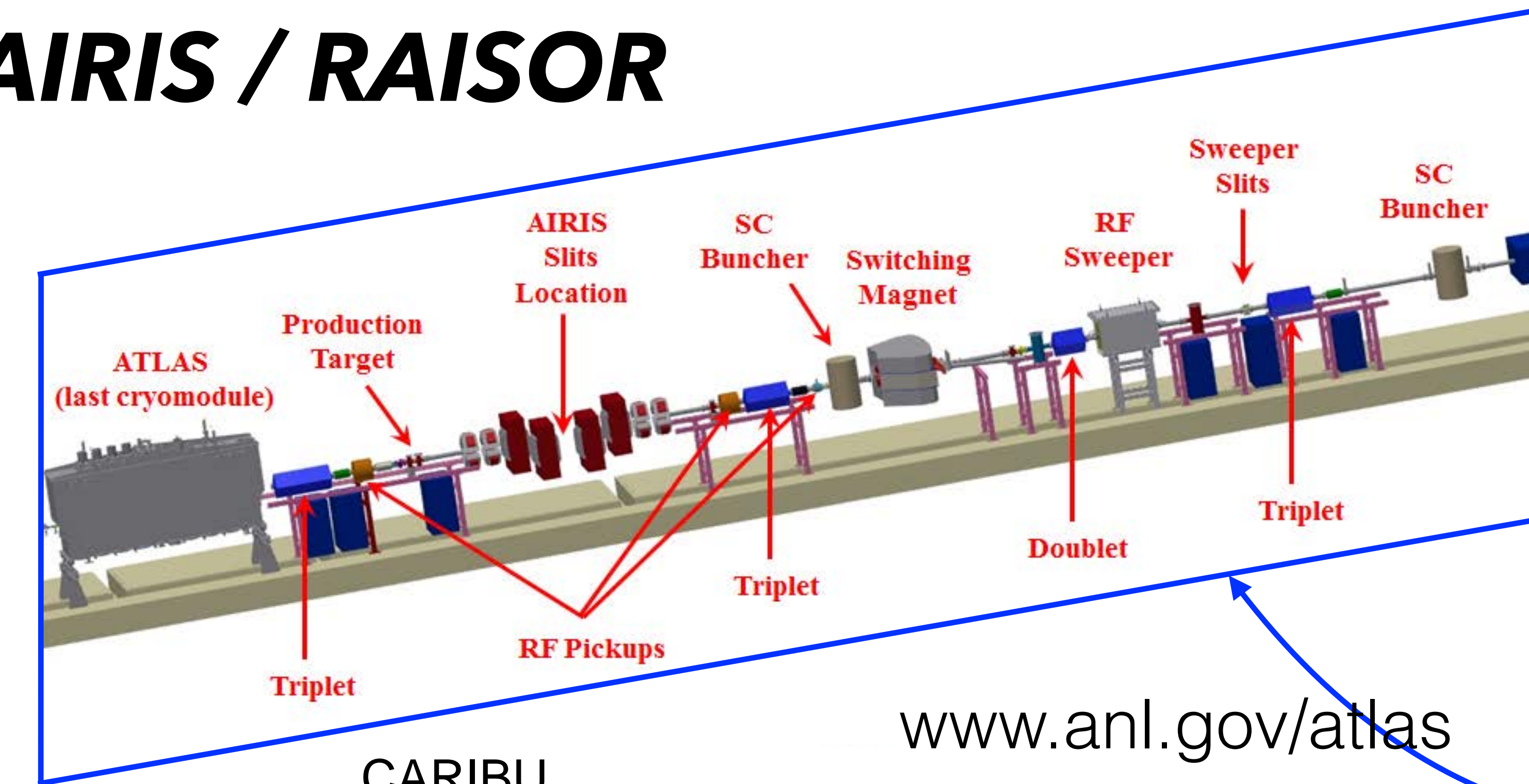
E.g. "ISOL" and in-flight at Argonne

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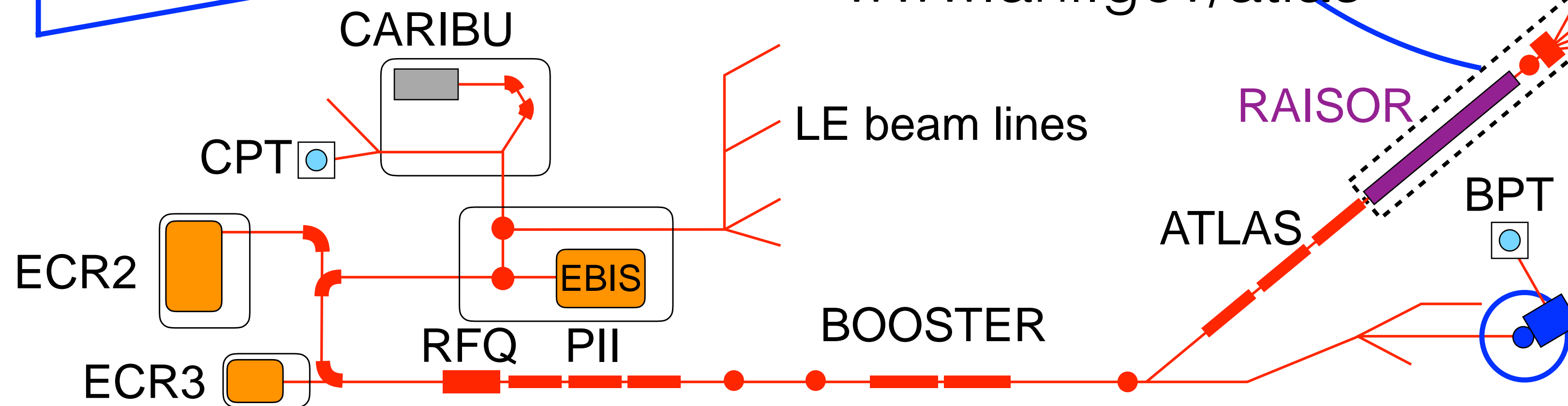
- **Stable beams**
- **In-flight beams**
- **Reaccelerate**
- **Low energy**



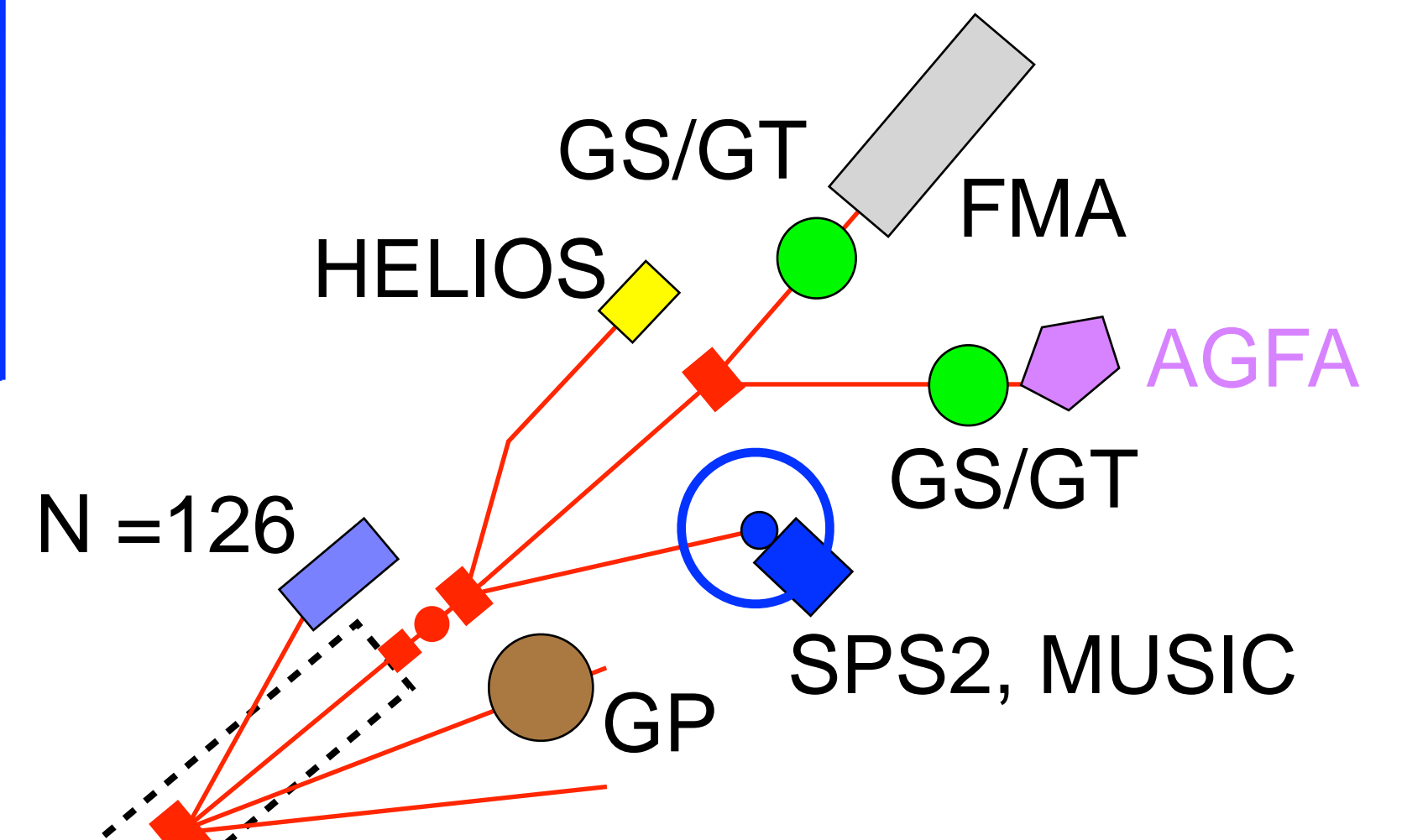
AIRIS / RAISOR



www.anl.gov/atlas



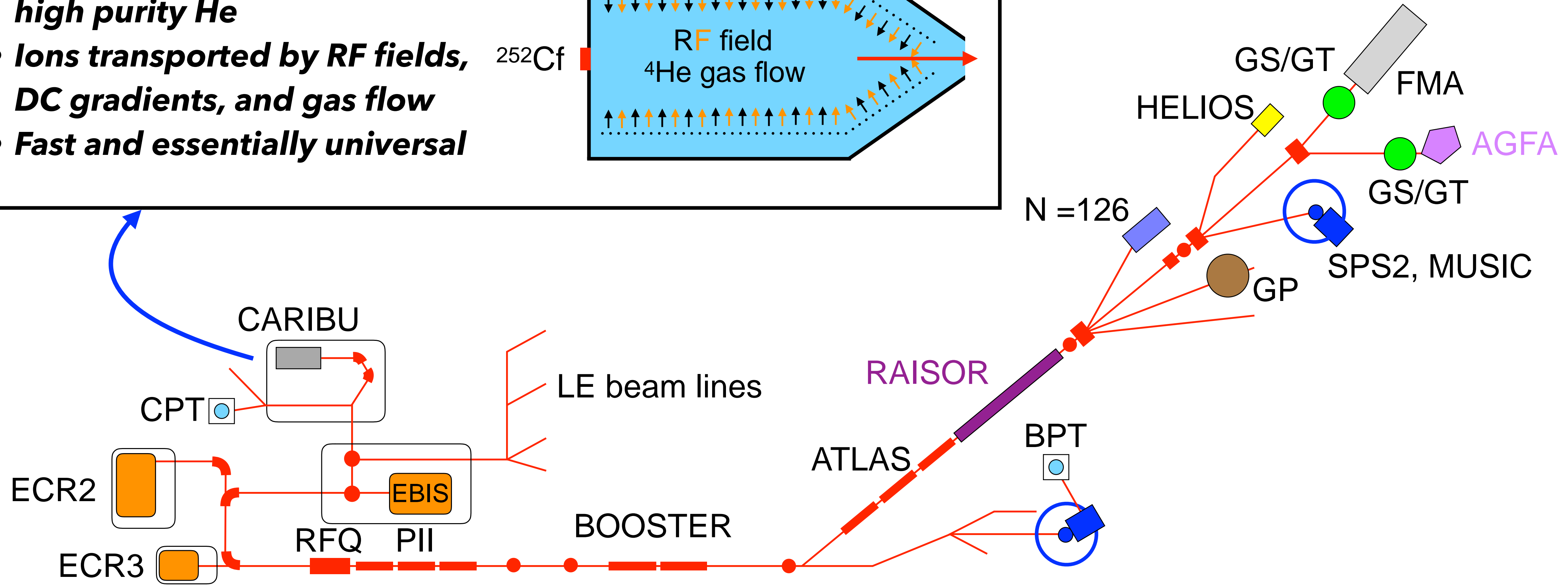
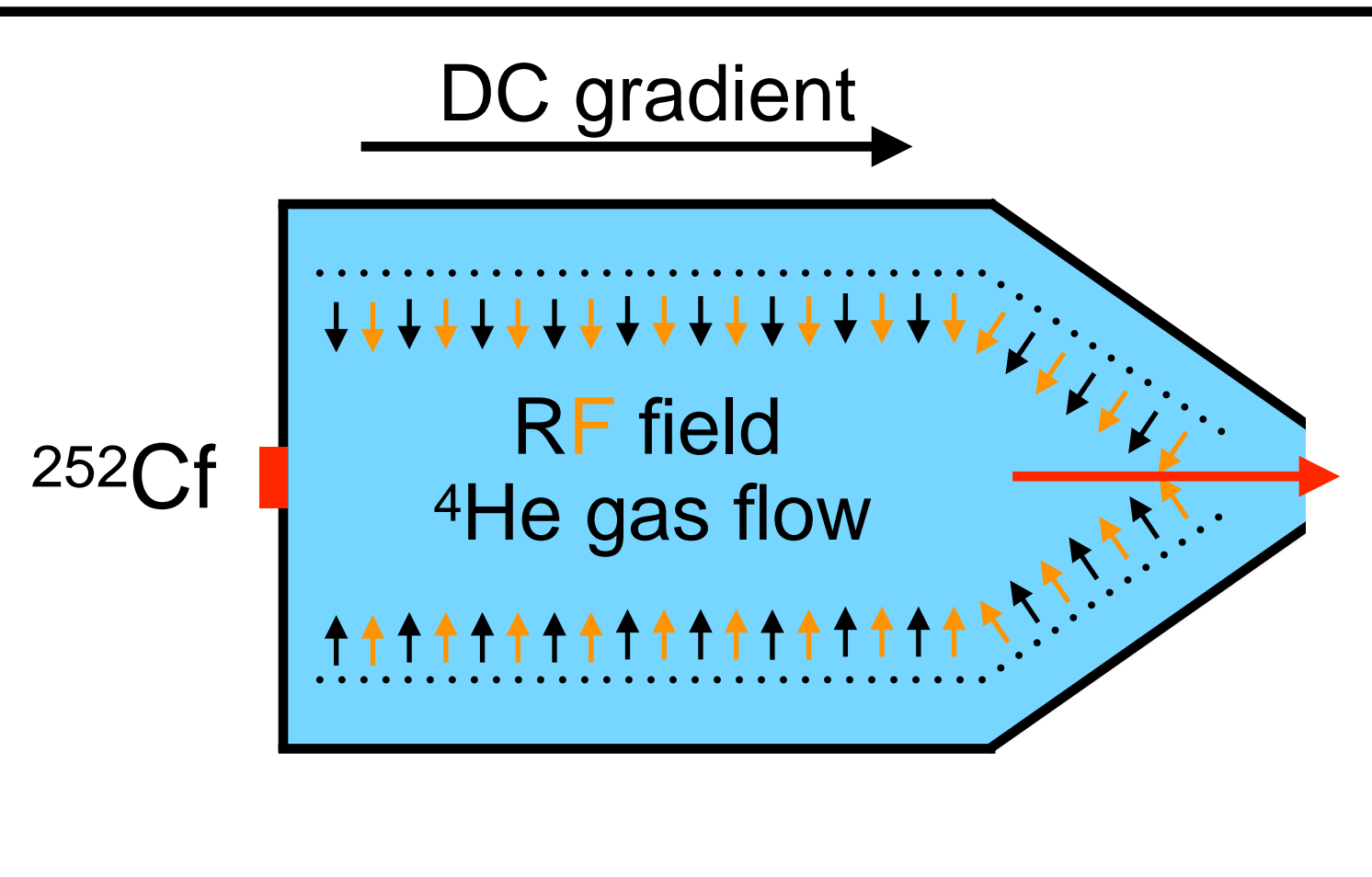
Primary beam from ATLAS, a few to 20 MeV/u, <few pμA



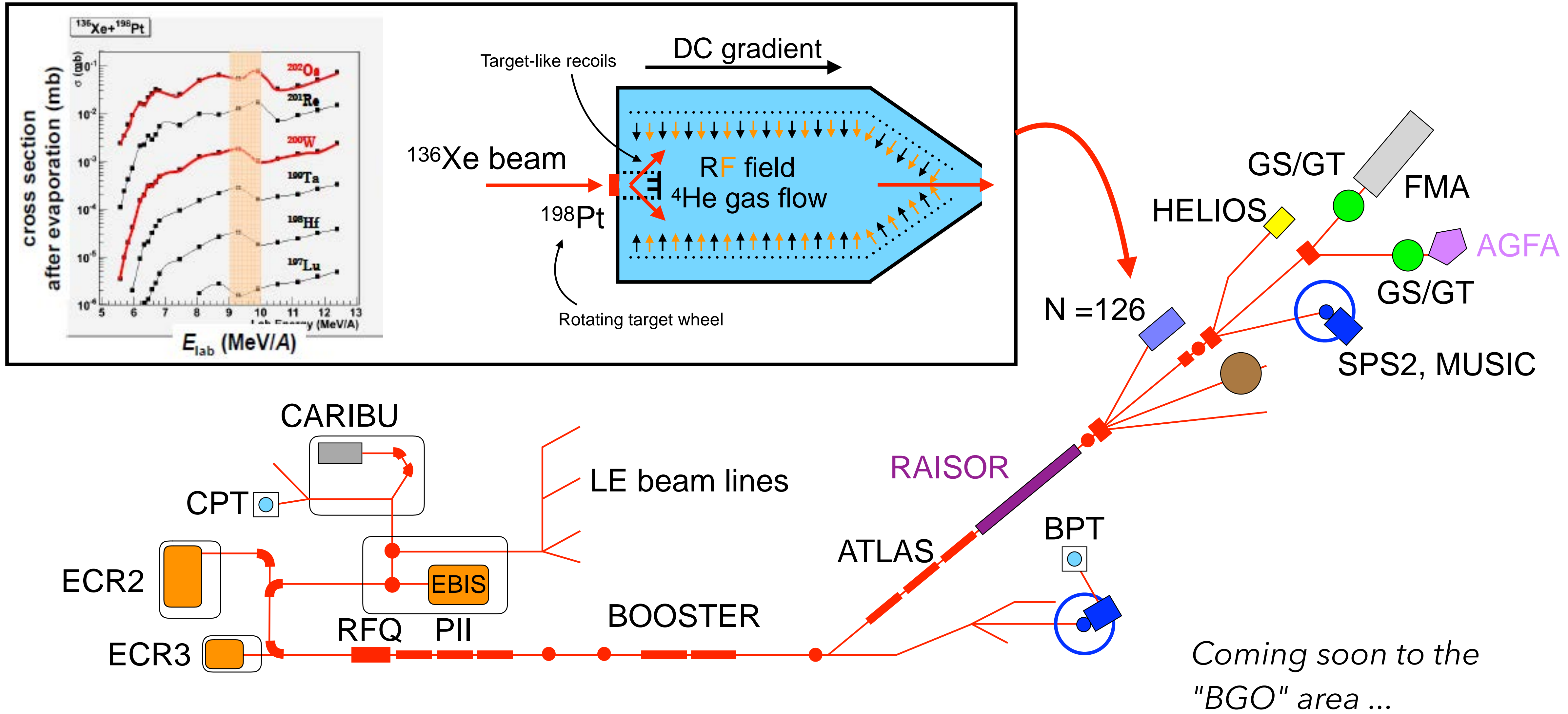
Provide in-flight beams to **all experimental areas downstream of ATLAS**

CARIBU

- **Fission fragments stopped in high purity He**
- **Ions transported by RF fields, DC gradients, and gas flow**
- **Fast and essentially universal**



Production of $N = 126$ nuclei

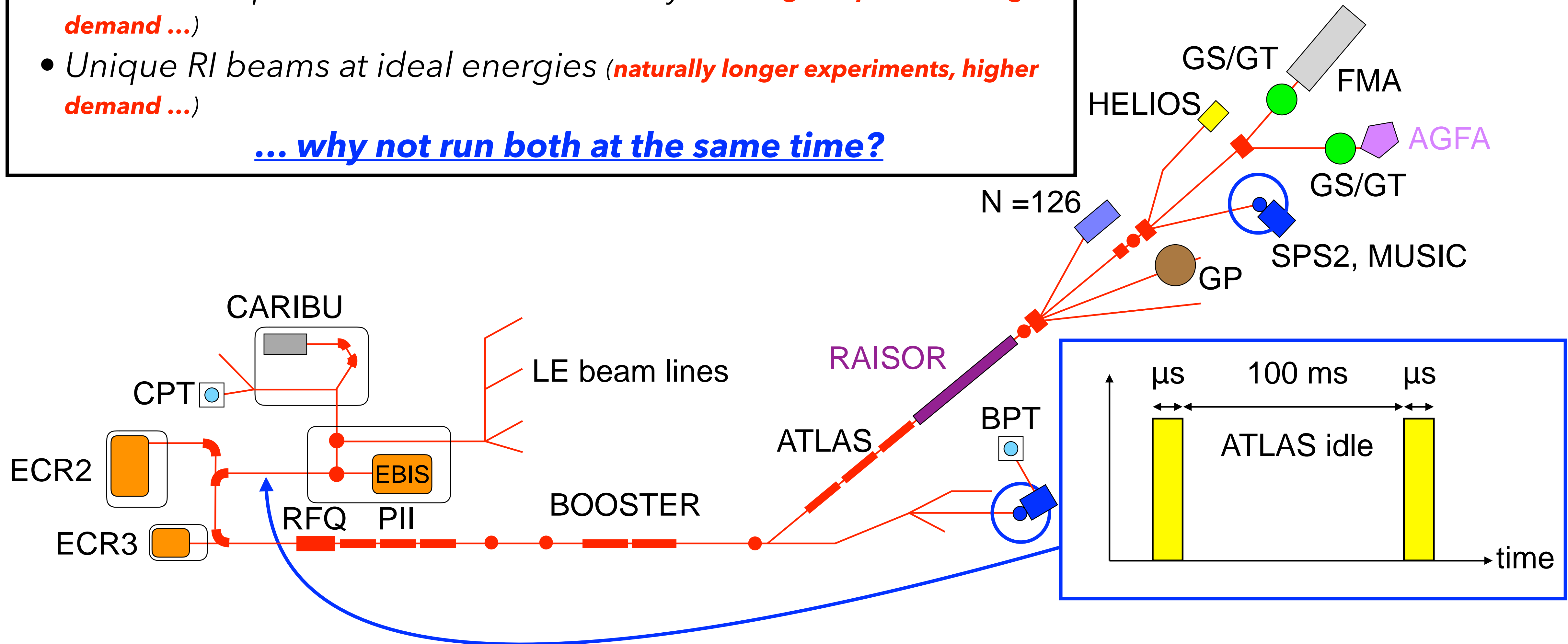


More than one beam!?

Higher demand for ATLAS beam time resolved ...

- The nations premier stable beam facility (**but longer experiments, higher demand ...**)
- Unique RI beams at ideal energies (**naturally longer experiments, higher demand ...**)

... why not run both at the same time?

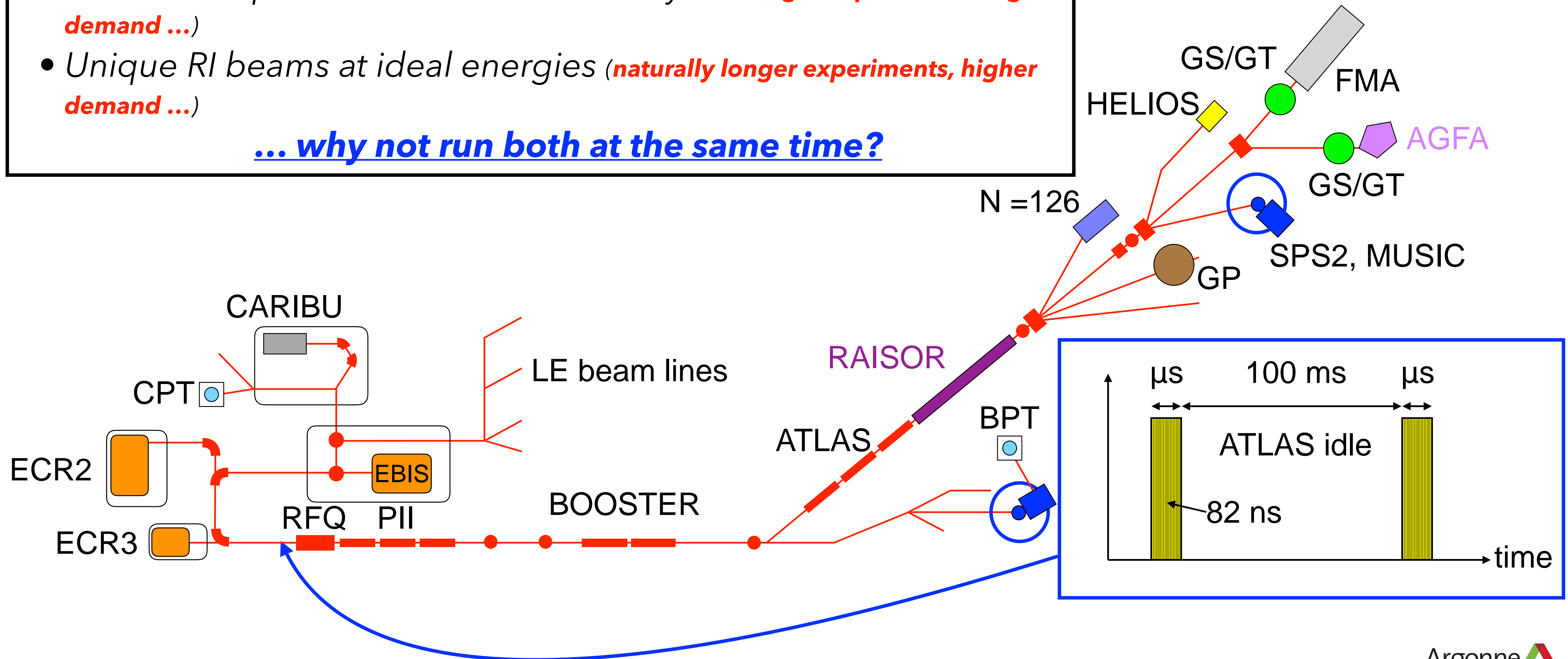


"Multi-user" facility

Higher demand for ATLAS beam time resolved ...

- The nations premier stable beam facility (**but longer experiments, higher demand ...**)
- Unique RI beams at ideal energies (**naturally longer experiments, higher demand ...**)

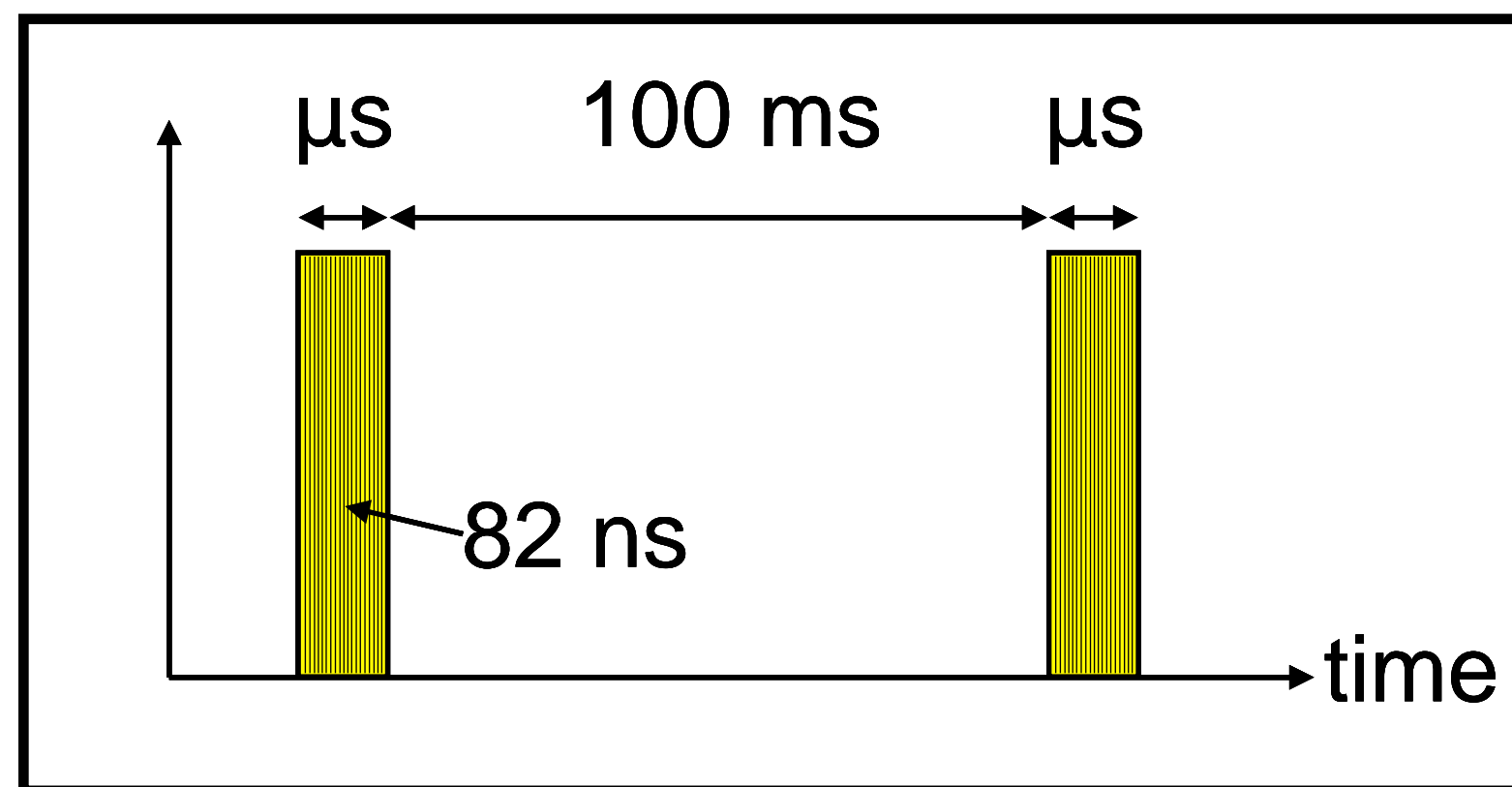
... why not run both at the same time?



...AMUU

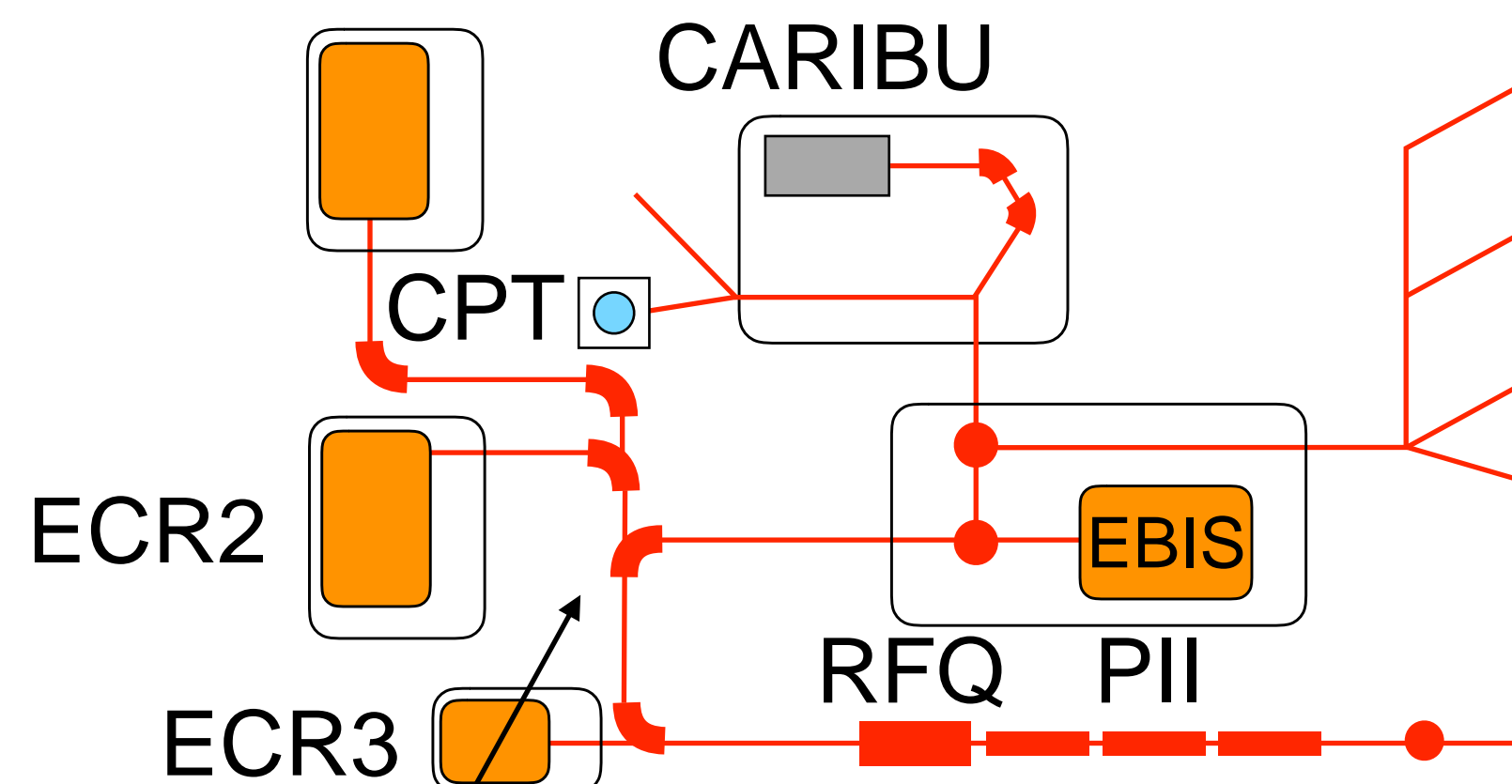
Key point:

No compromise made – each beam is optimal



e.g. CARIBU beam at 10 MeV/u

Future alt. ECR



Modified injection and LEBT

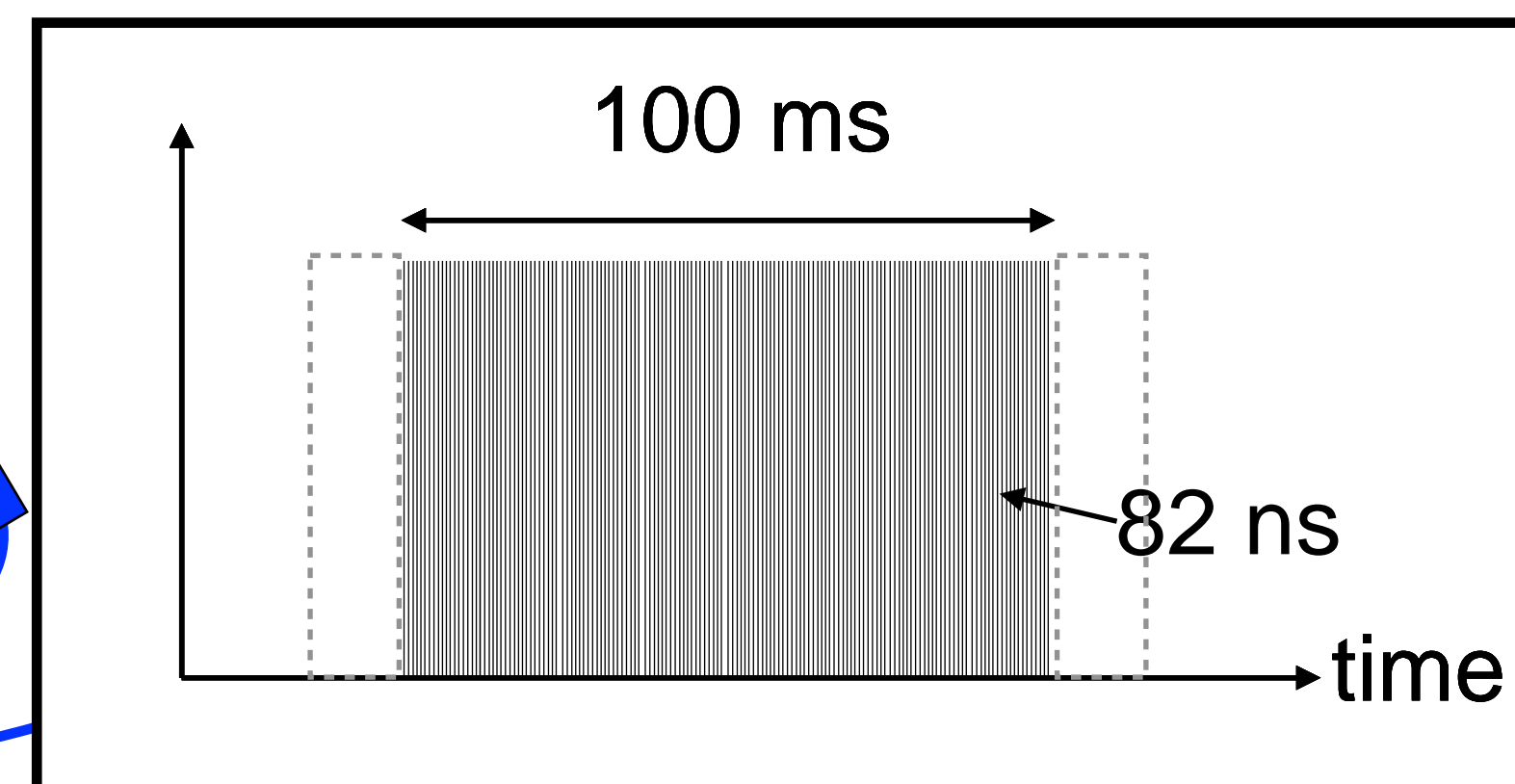
BOOSTER

Pulsed switchyard

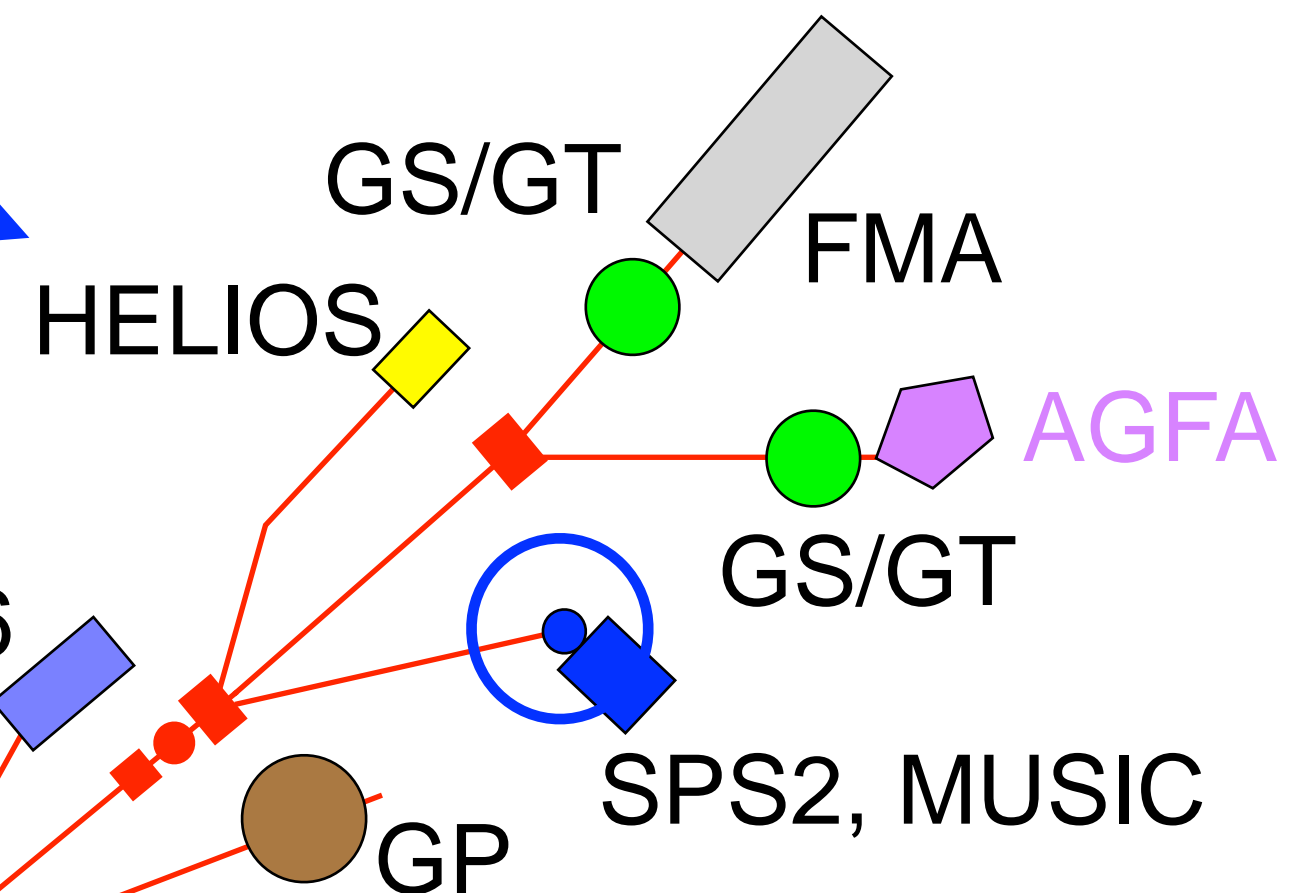
ATLAS

RAISOR

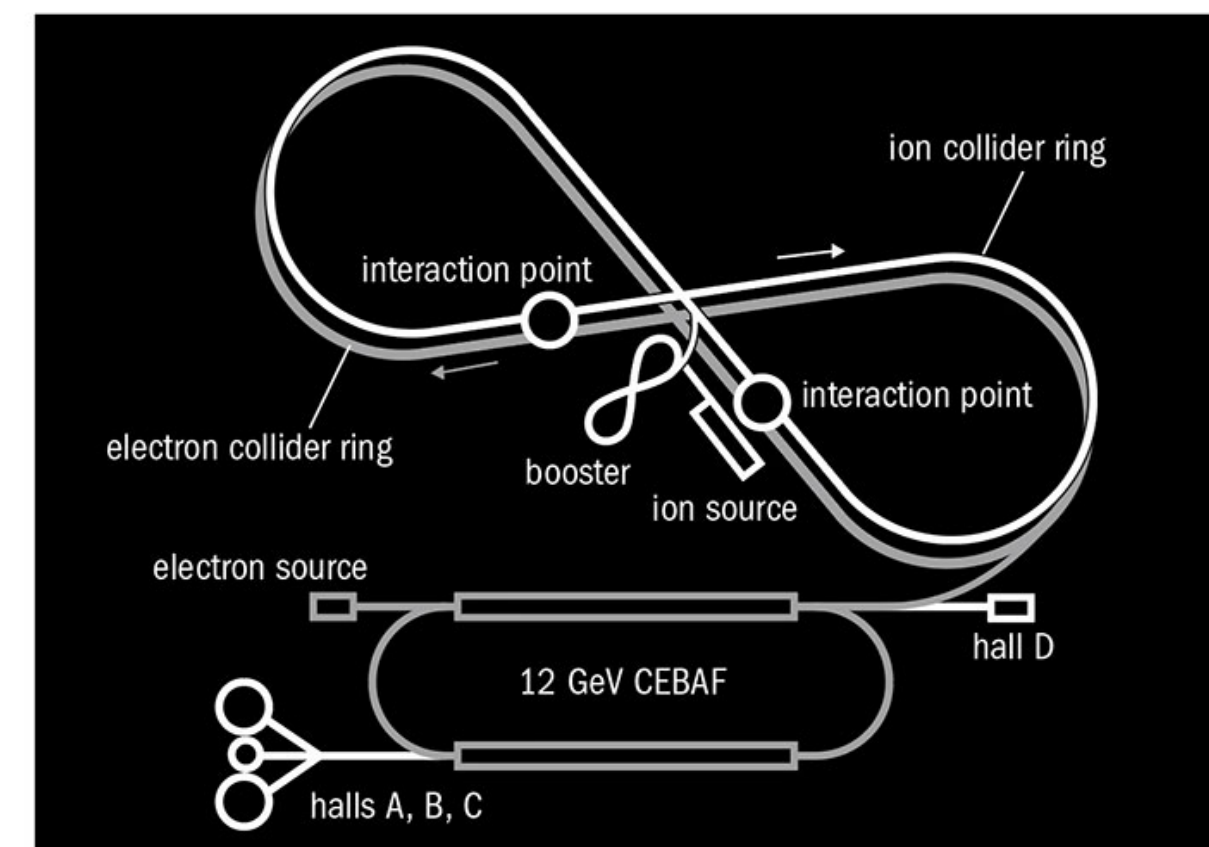
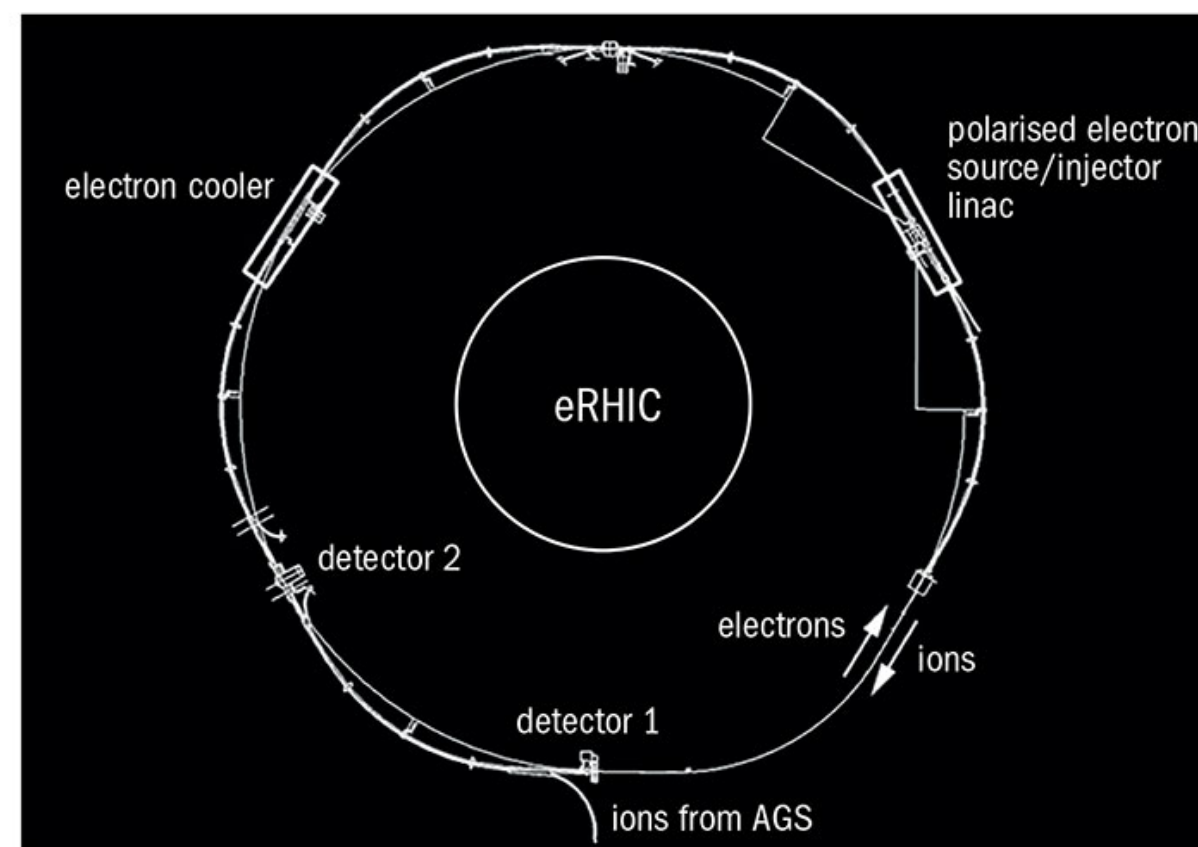
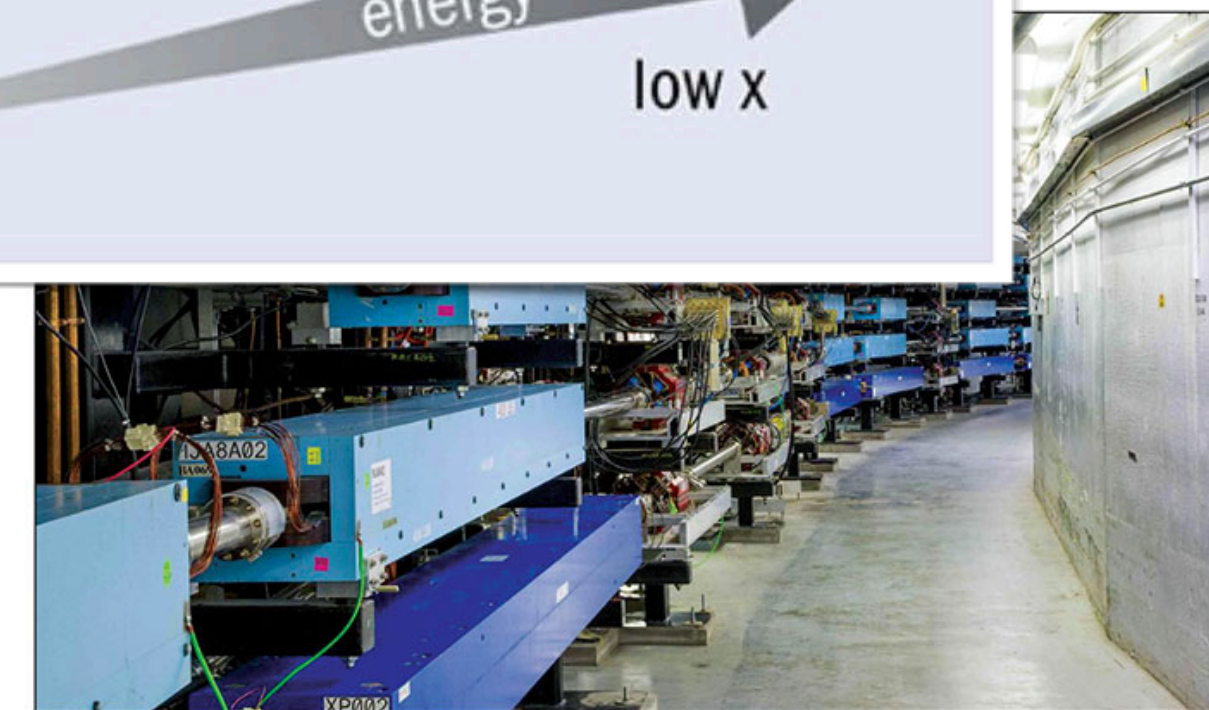
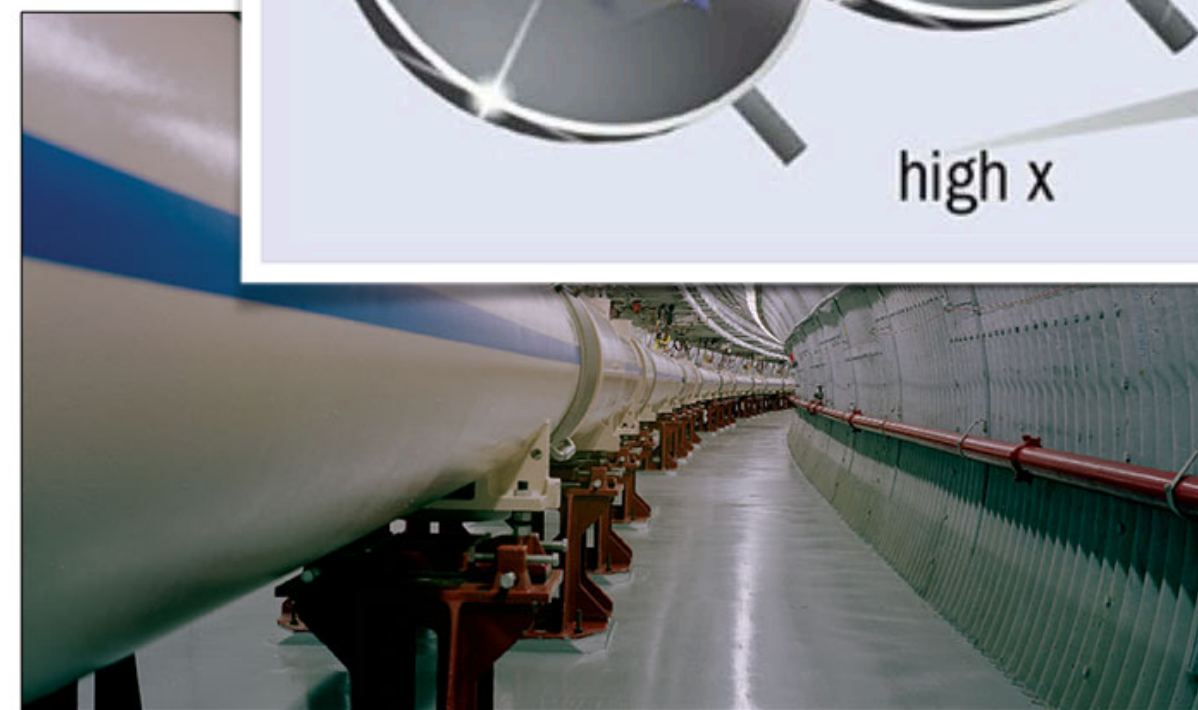
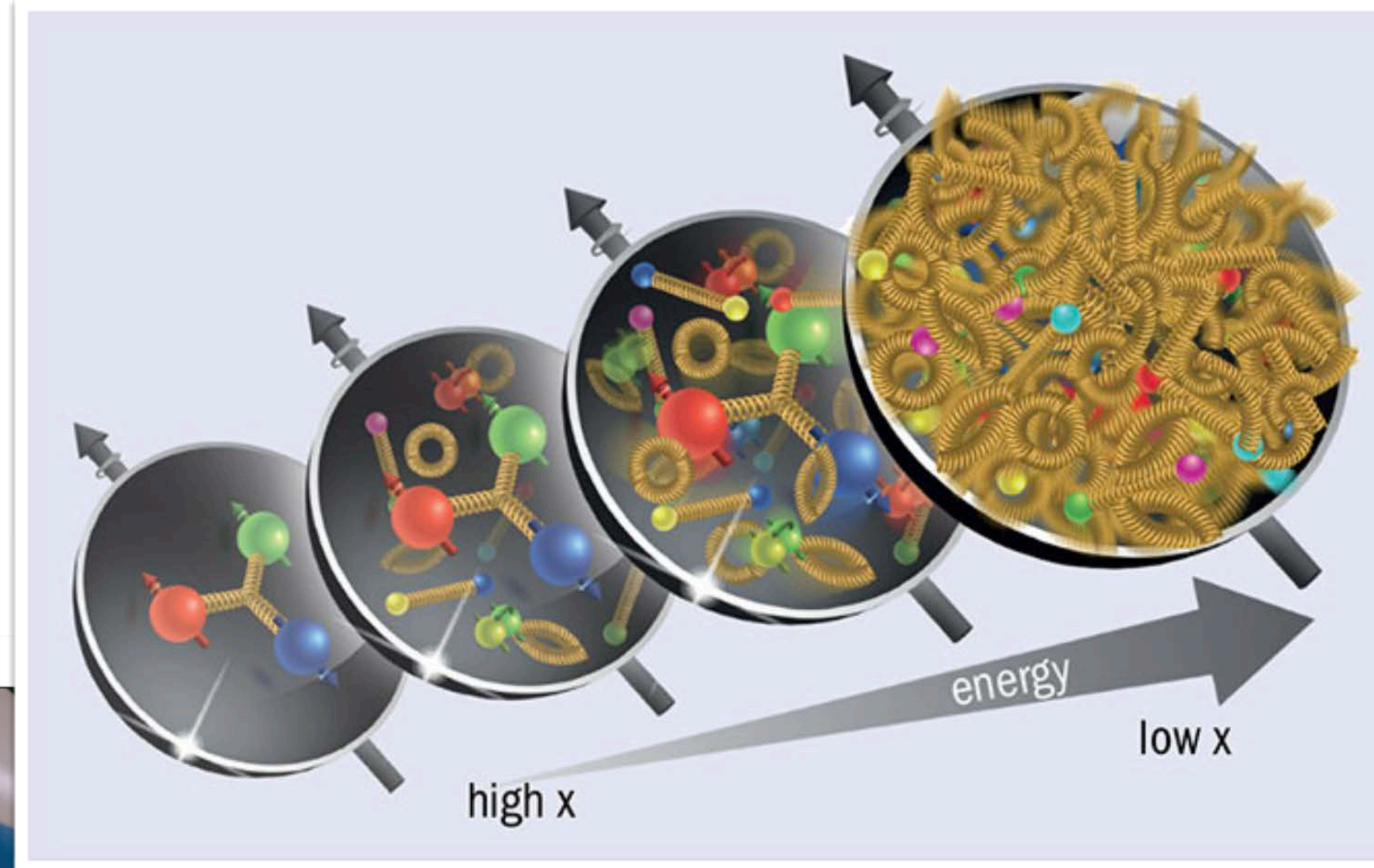
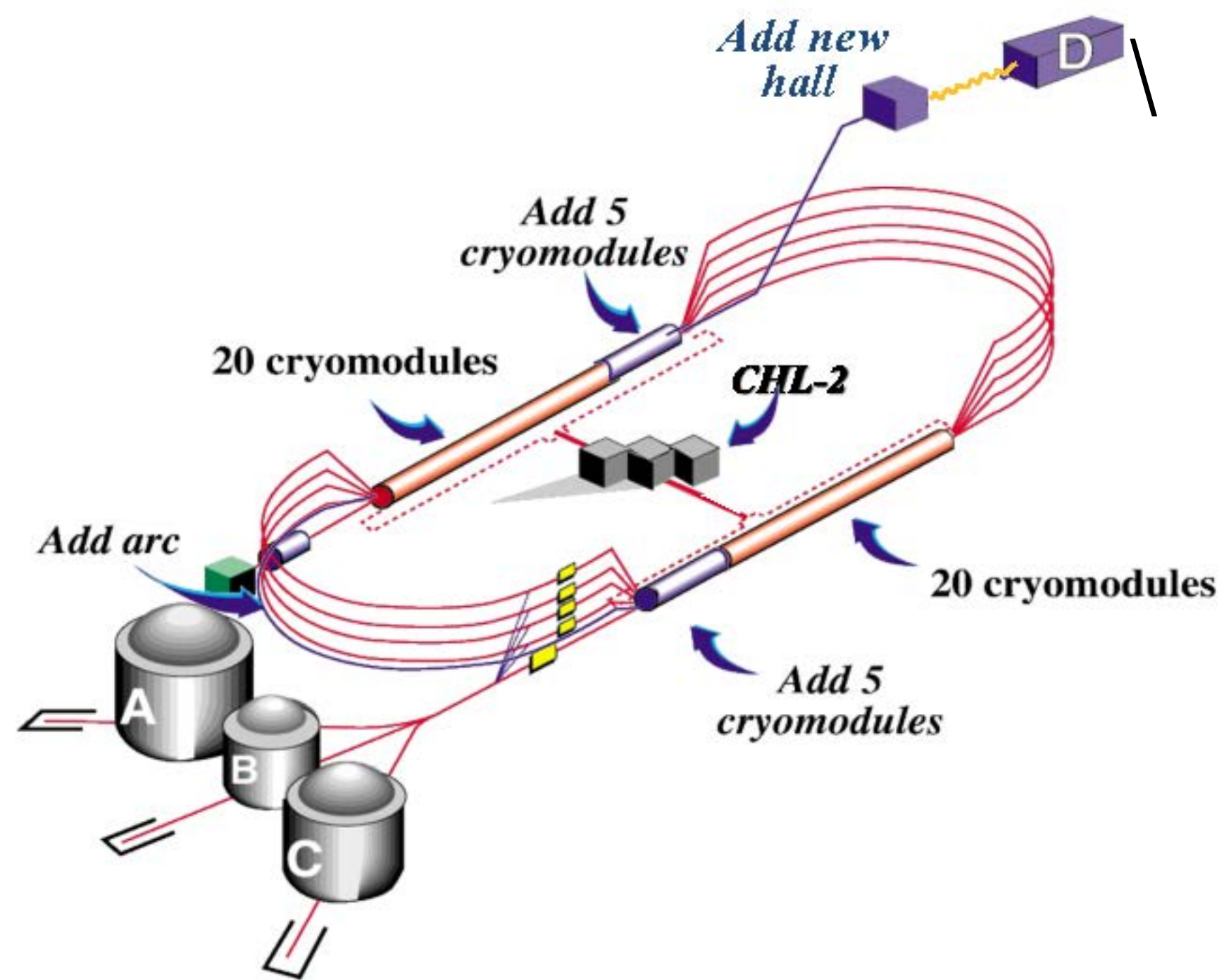
BPT



e.g. stable beam at 6 MeV/u



Non-hadronic probes ...



Nuclear cartography

PAPER

Phys. Educ. **52** (2017) 064002 (9pp)

iopscience.org/ped

Nuclear cartography: patterns in binding energies and subatomic structure

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CrossMark

Abstract

Nuclear masses and binding energies are some of the first nuclear properties met in high school physics, and can be used to introduce radioactive decays, fusion, and fission. With relatively little extension, they can also illustrate fundamental concepts in nuclear physics, such as shell structure and pairing, and to discuss how the elements around us were formed in stars. One way of visualising these nuclear properties is through the nuclide chart, which maps all nuclides as a function of their proton and neutron numbers. Here we use the nuclide chart to illustrate various aspects of nuclear physics, and present 3D visualisations of it produced as part of the binding blocks project.

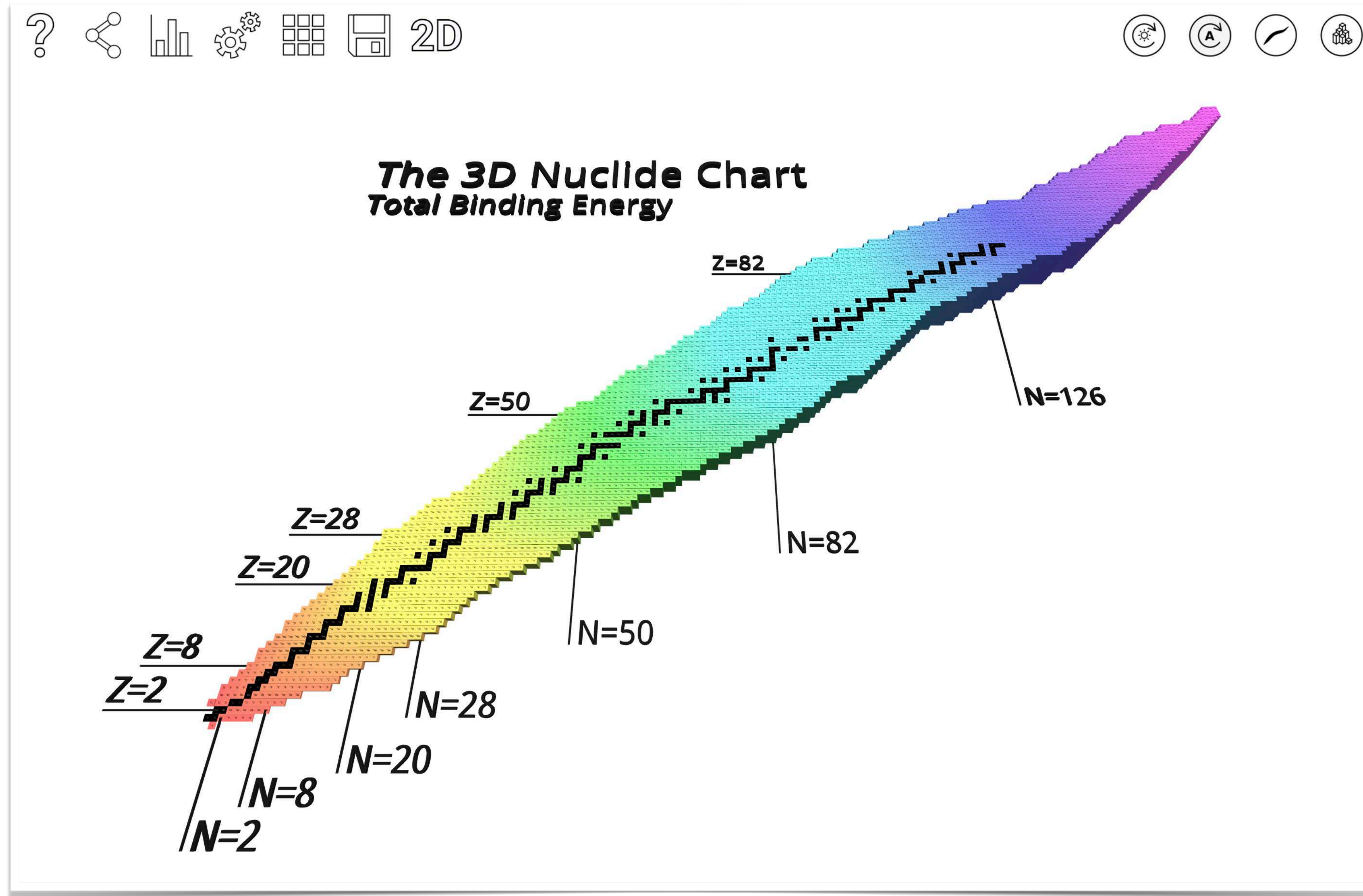
The 3D Nuclide Chart

Putting the joy back into nuclear data.



WARNING: support for touch devices is currently experimental!

Nuclear playground

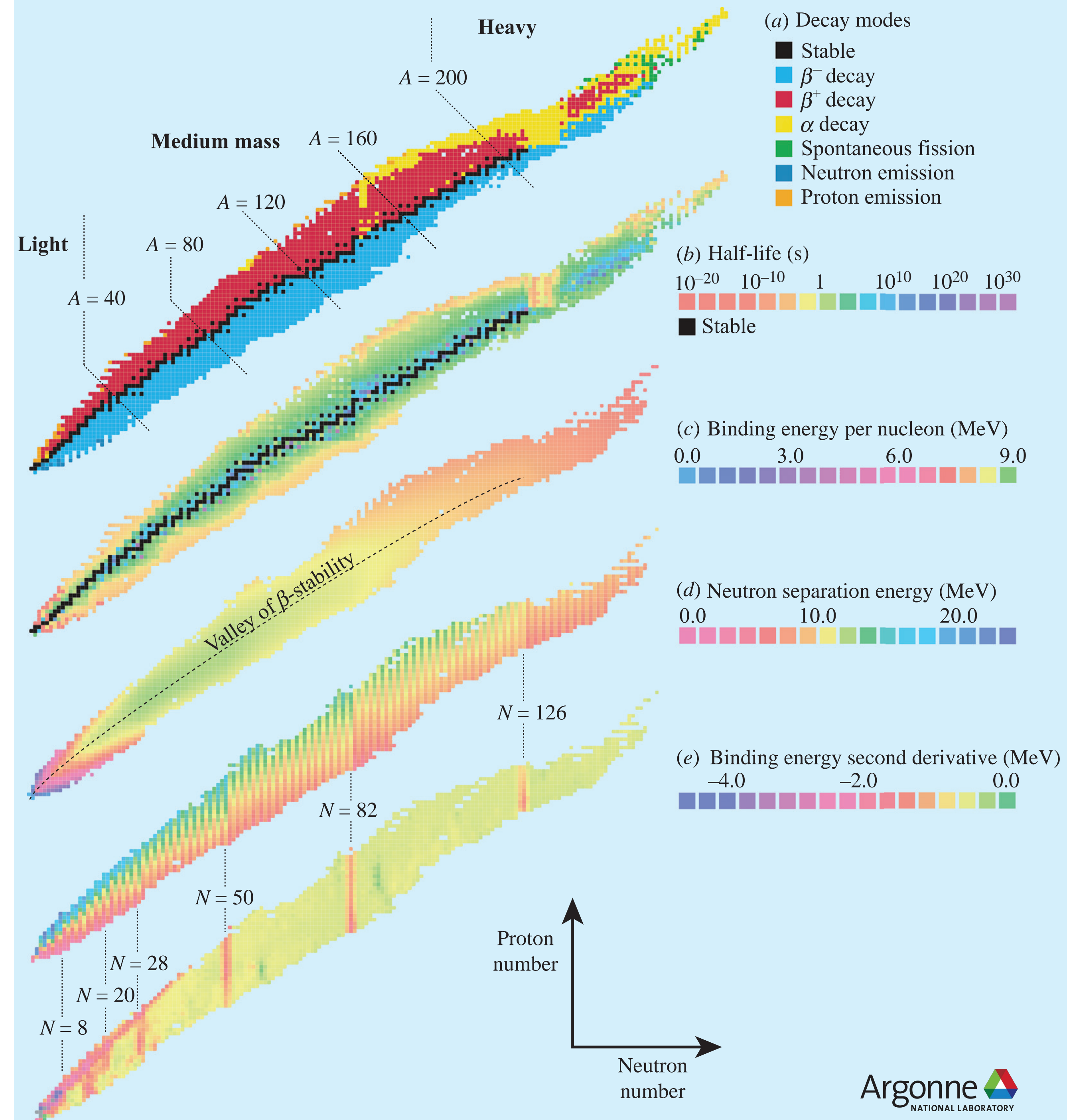


<https://people.physics.anu.edu.au/~ecs103/chart3d/>

Nuclear playground

As experimentalists we can:

- Determine the decay mode
- Determine the half life
- Determine the mass, binding
- Reaction cross sections
- Moments
- Transition rates / energies

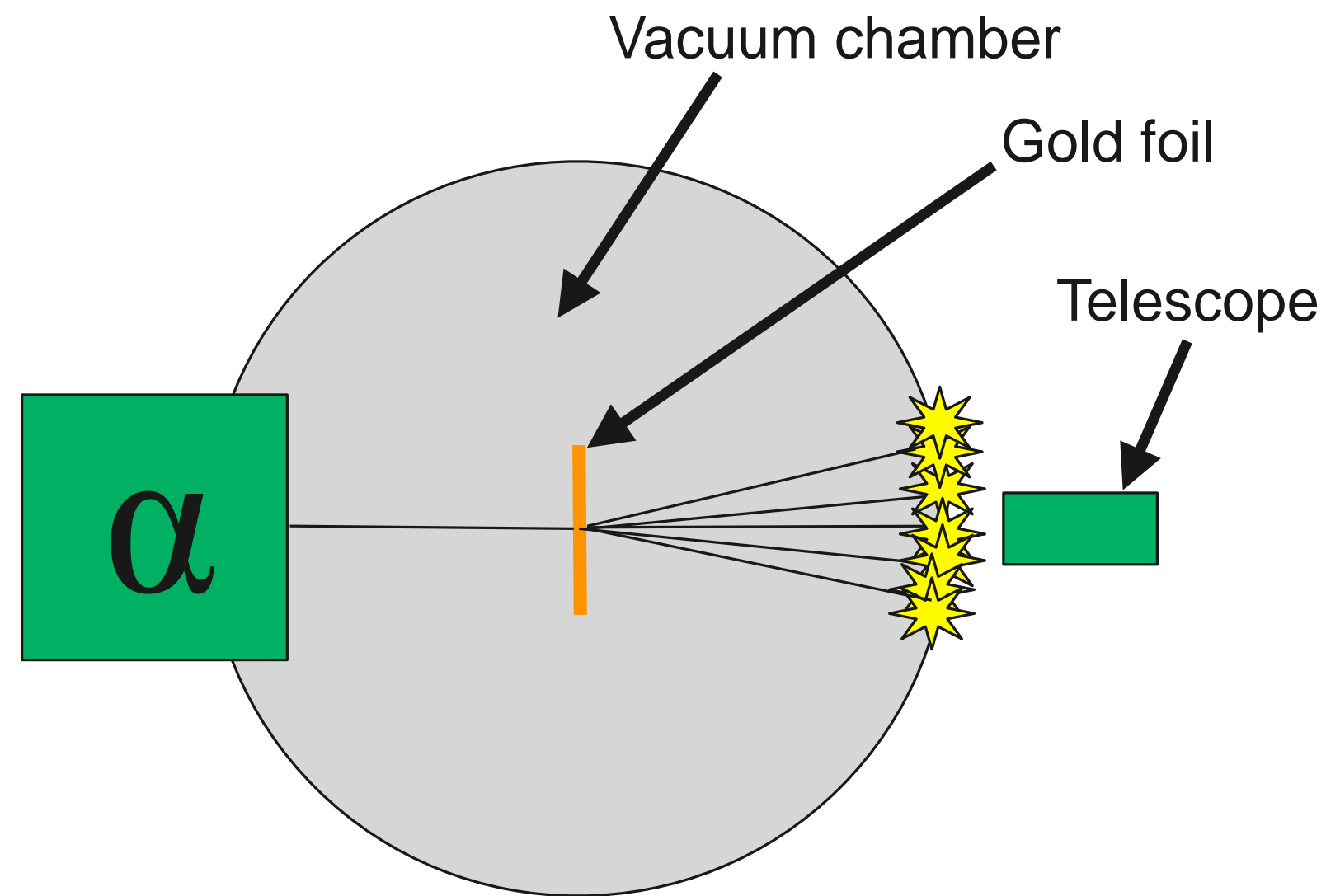


What can one observe, why is it tough?

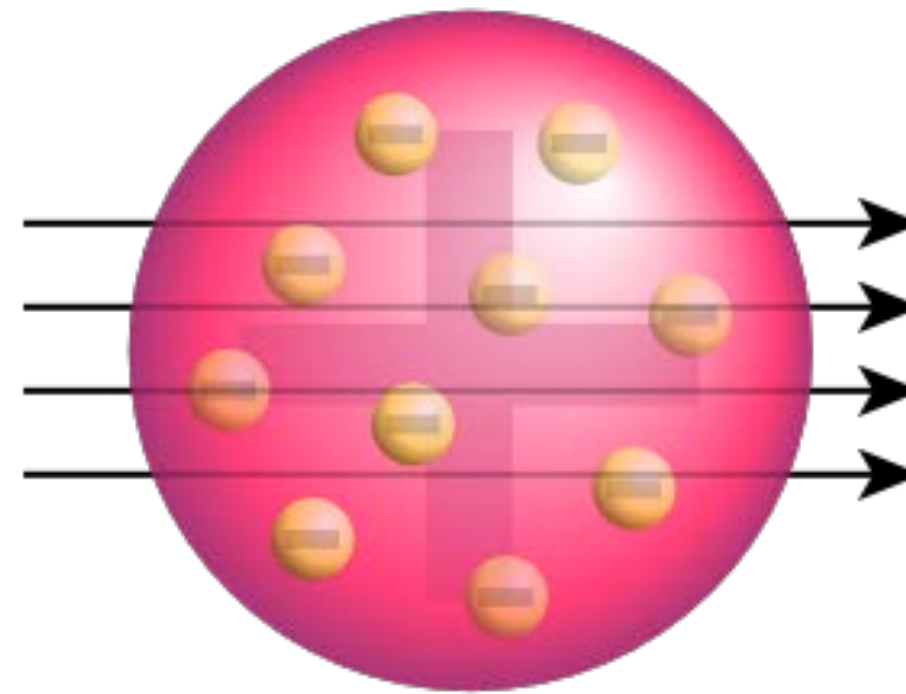
- **As you will have heard over the work shop, there are only a handful of physical properties of nuclei one can probe and link to models**
- **For a low-energy experimentalists the challenges are many ... as we've just seen only about 4% of the nuclei predicted to be bound are stable, the rest we have to make ...**
- **The best probes are typically nuclei themselves ...**
- **Then there is the connection to theory and understanding, what we measure are not always instructive without model-dependent conversions (plenty of discussions in lectures 2 and 3)**

To begin at the beginning ...

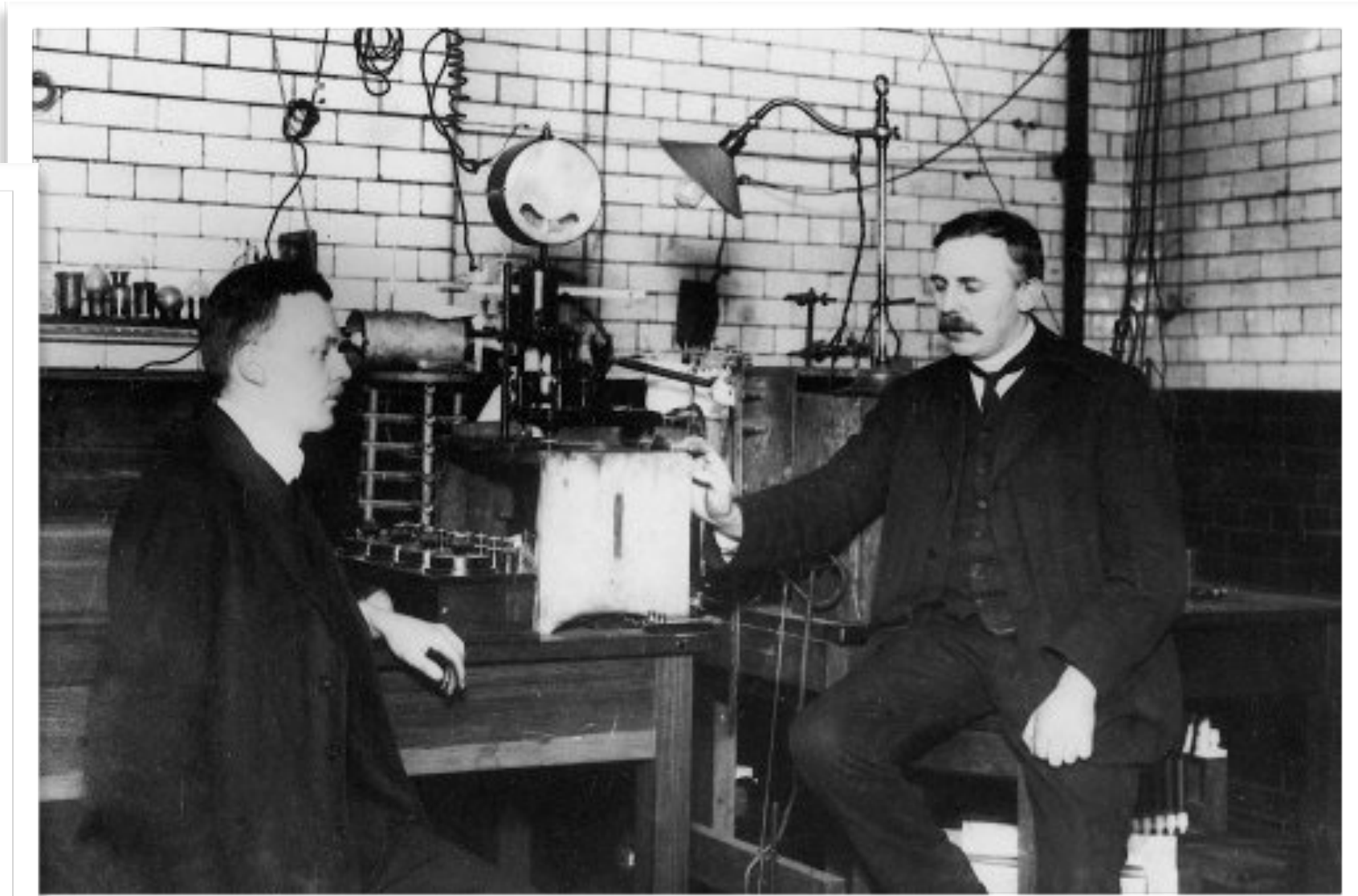
The Geiger-Marsden experiment



The plum-pudding idea seemed reasonable: this result would fit expectations

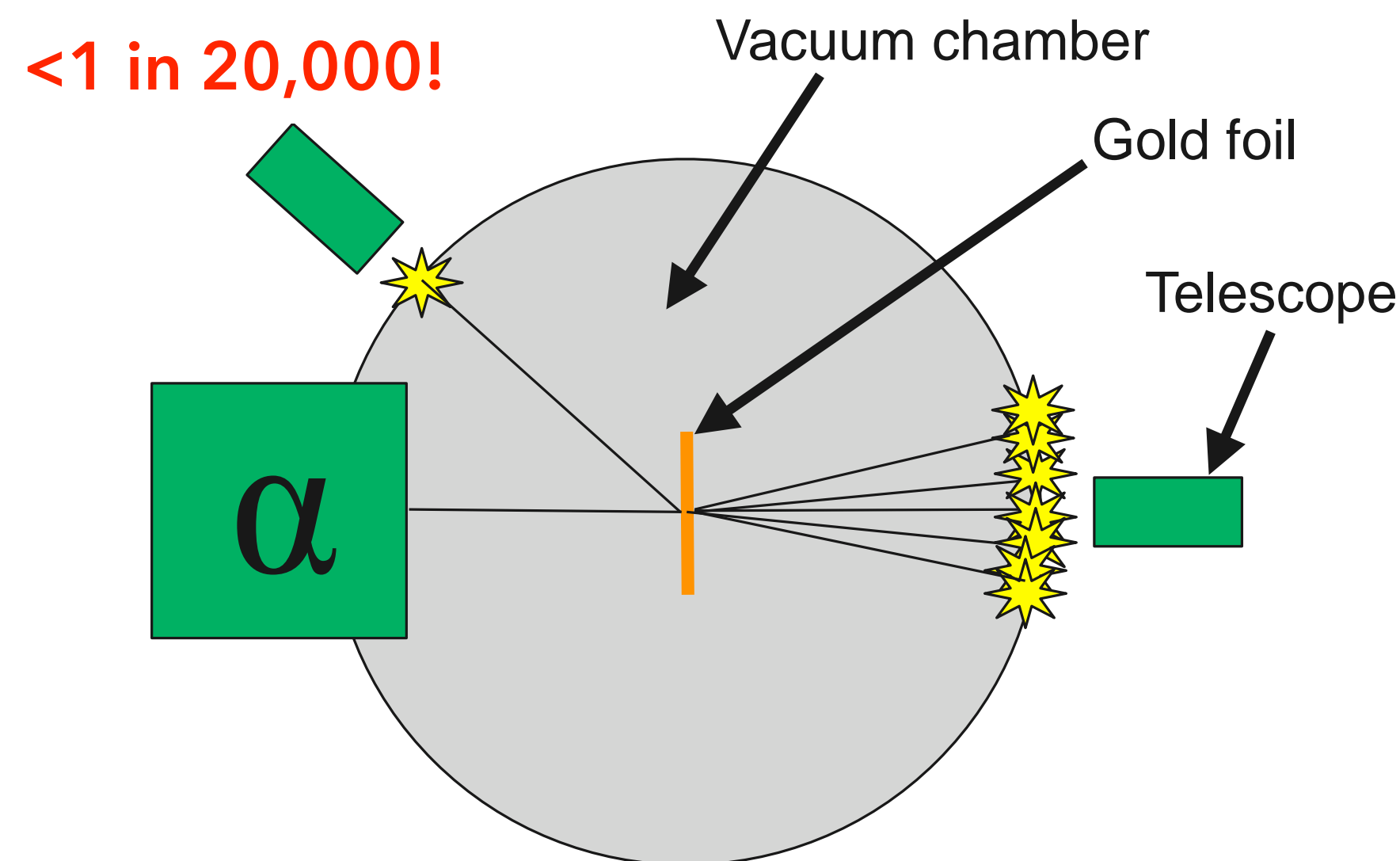


- A 0.1 Ci radium source
- $\sim 10^{10}$ α particles per second (~ 1 nA of ^4He)
- α particles of 7.7 MeV (~ 1.9 MeV/u)
- A gold foil of 0.00004 cm thick (~ 0.8 mg/cm²)
- A telescope was used to look at flashes of light on a zinc sulphide screen

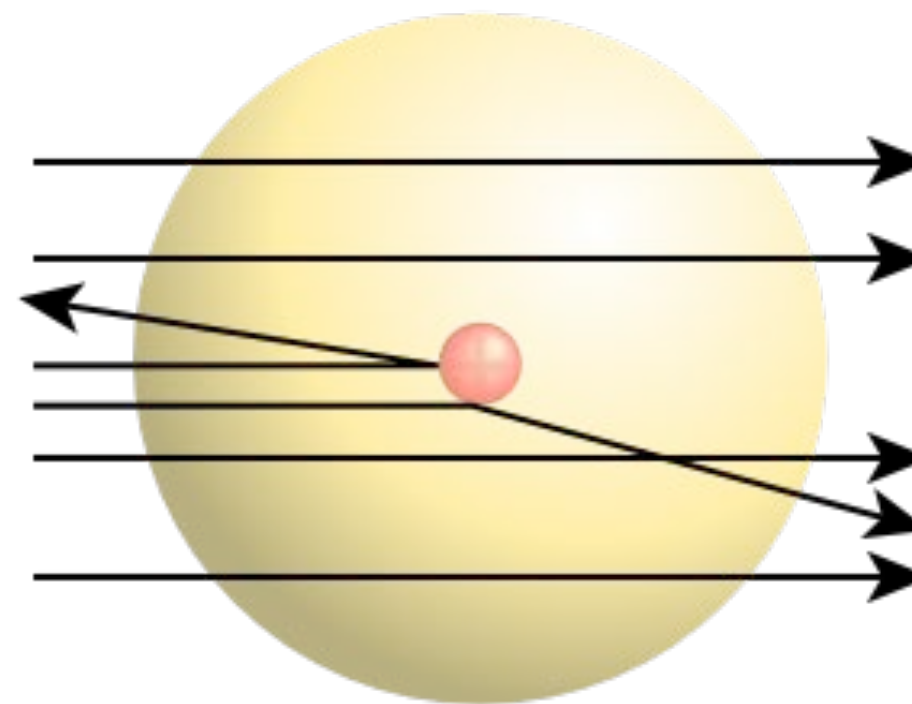


Reaction ... information

The Geiger-Marsden experiment



The atom has a dense, positive mass at the centre ... it is mostly empty!



- A 0.1 Ci radium source
- $\sim 10^{10}$ α particles per second (~ 1 nA of ^4He)
- α particles of 7.7 MeV (~ 1.9 MeV/u)
- A gold foil of 0.00004 cm thick (~ 0.8 mg/cm²)
- A telescope was used to look at flashes of light on a zinc sulphide screen

"It was almost as incredible as if you fired a 15-inch shell at a piece of tissue paper and it came back and hit you." – E. Rutherford.

This has *all the same* ingredients a modern nuclear reaction experiment:

- A beam
- A target
- A chamber
- Reaction products
- A detector
- ... deduce something about the gold ... its

SIZE, that's

it's bound,...

Neutrons, strong force, shell structure

Still not quite the nucleus ... what's missing

- * What is the positive charge? Protons (Rutherford)
- * What else is in there? Neutrons (Chadwick)
- * How does it stick together? (A strong nuclear potential)



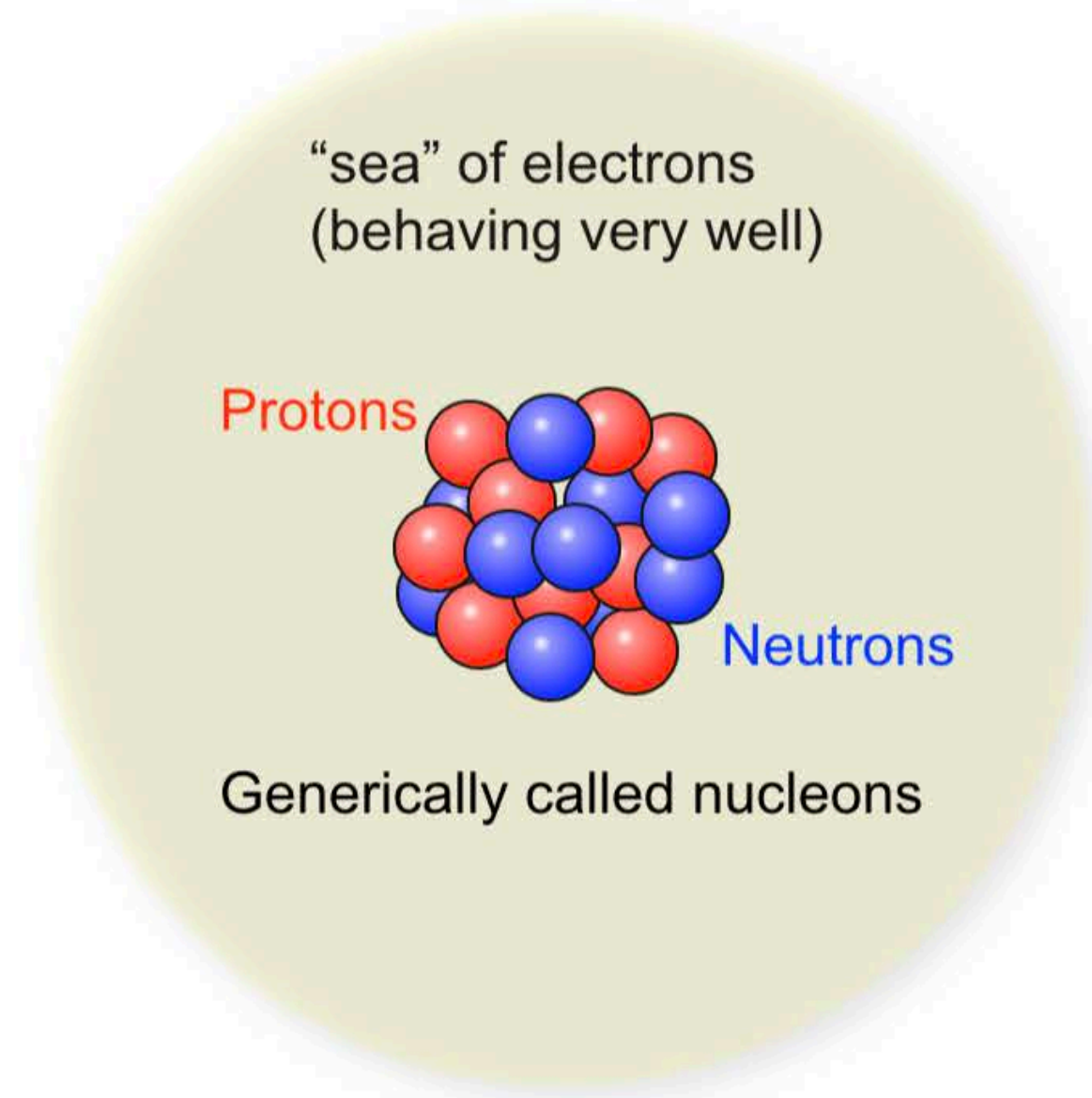
Neils Bohr,
Atomic theory



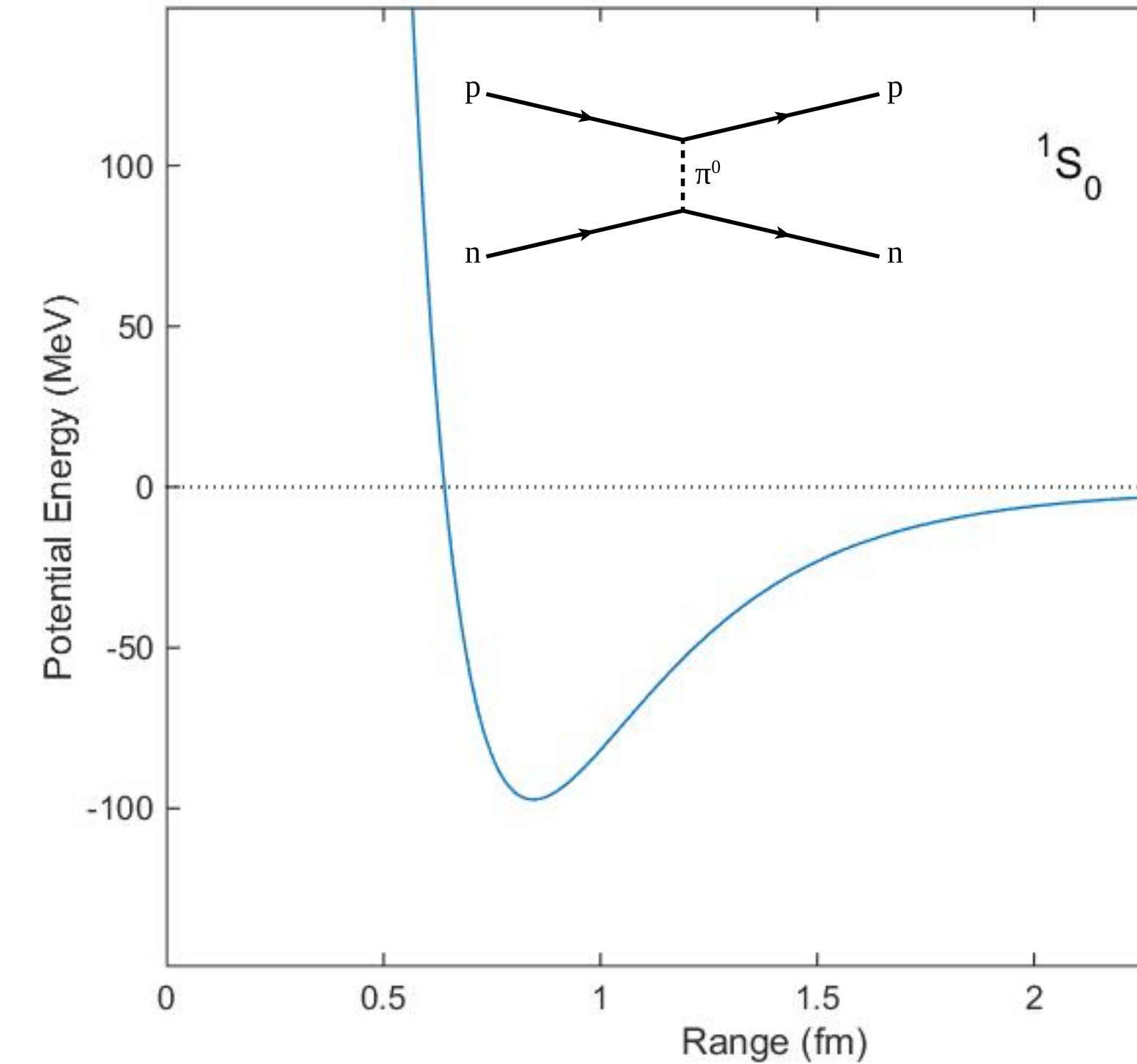
James
Chadwick,
neutrons



Hideki Yukawa,
potential



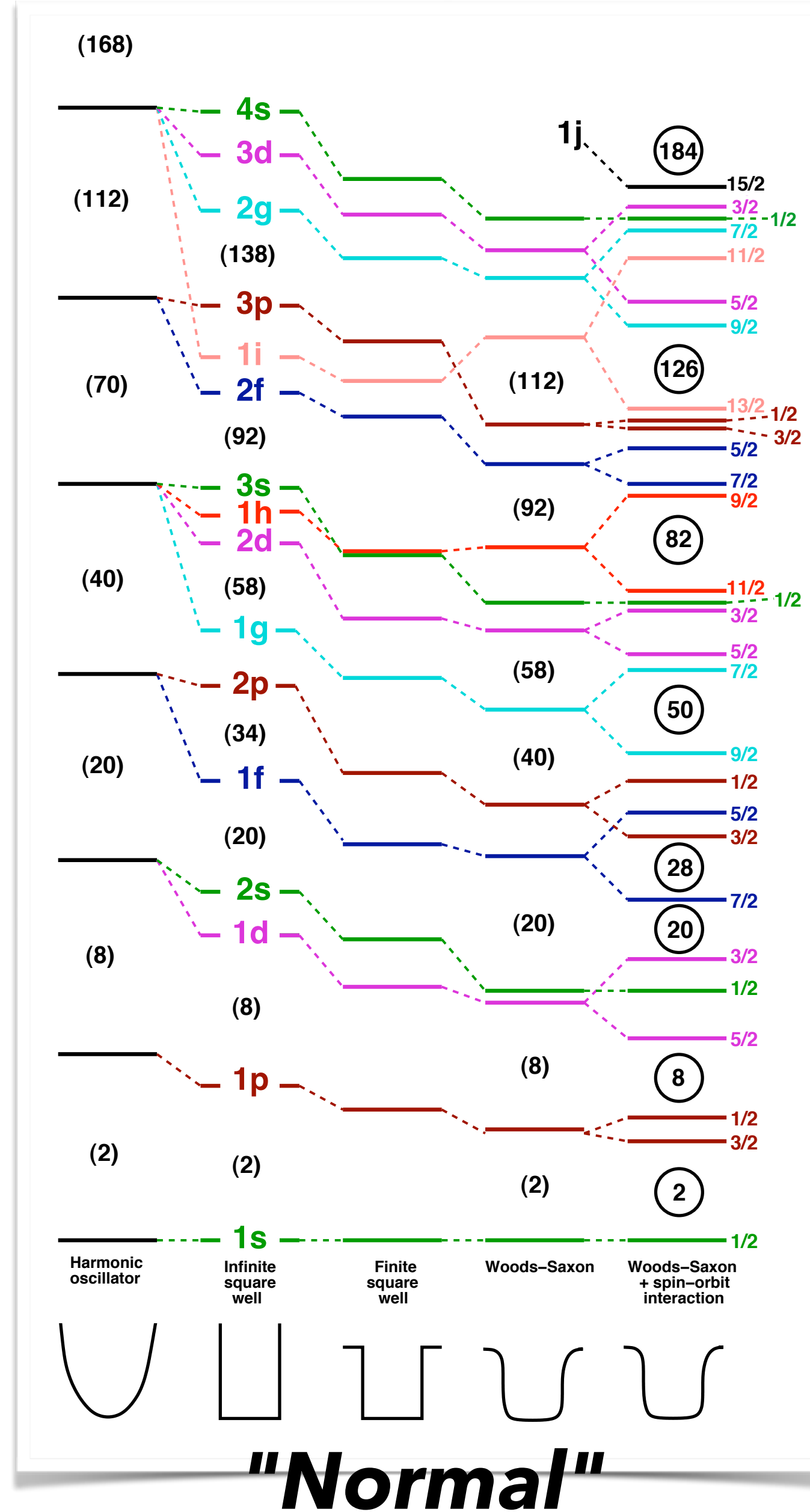
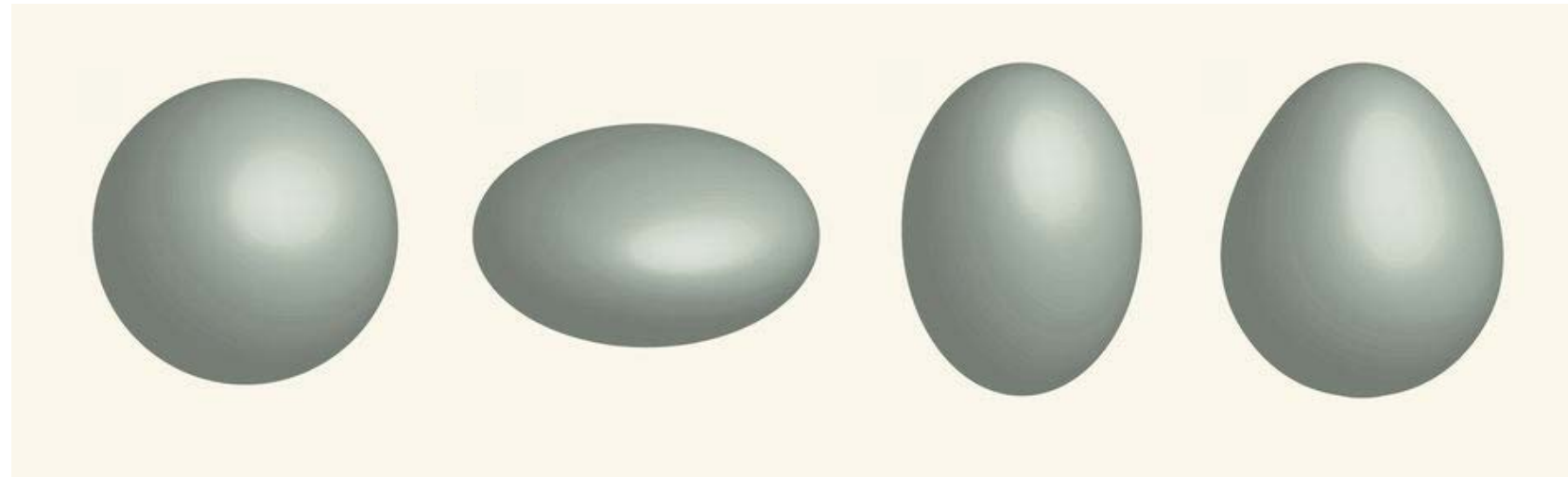
Hugely exaggerated cartoon on the atom. On this scale the nucleus would be a tiny period at the center.


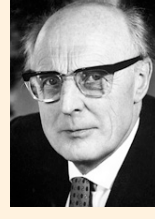


Neutrons, strong force, shell structure, shapes

• Shapes (all nucleons, **collective**)

• Single-particle (inert cores, valence nucleons)



Nuclear Shell Structure

126

$P_{1/2}$

$f_{5/2}$

$i_{13/2}$

$P_{3/2}$

$h_{9/2}$

$f_{7/2}$

➔

82

$d_{3/2}$

$h_{11/2}$

$s_{1/2}$

$g_{7/2}$

$d_{5/2}$

50

$g_{9/2}$

✓

Around stability
N/Z ~1-1.6

?

Neutron-rich nuclei
N/Z ~3

"Exotic"

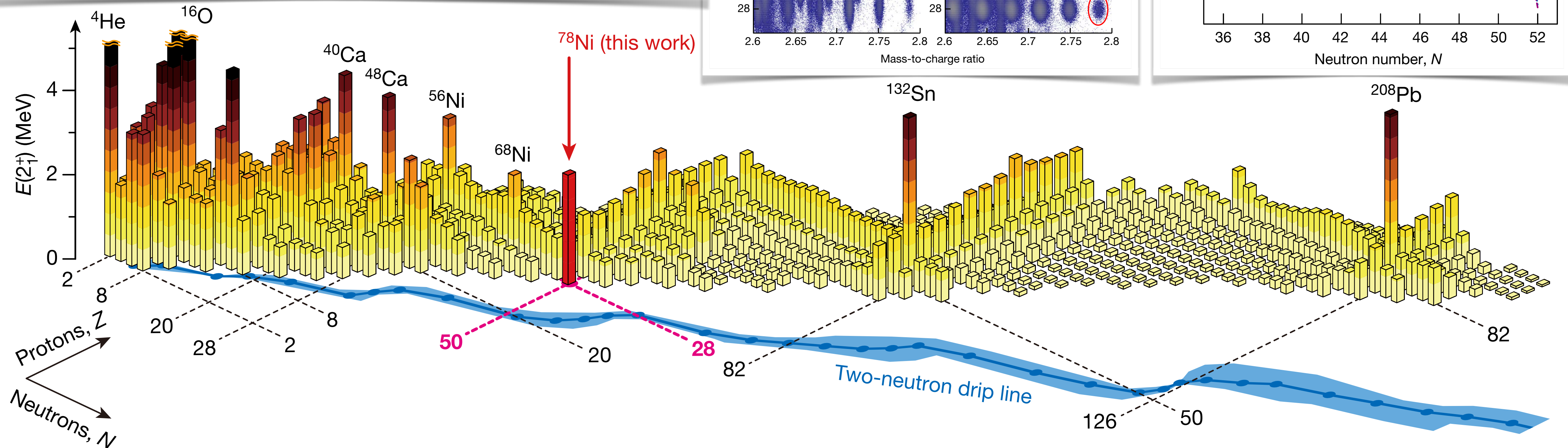
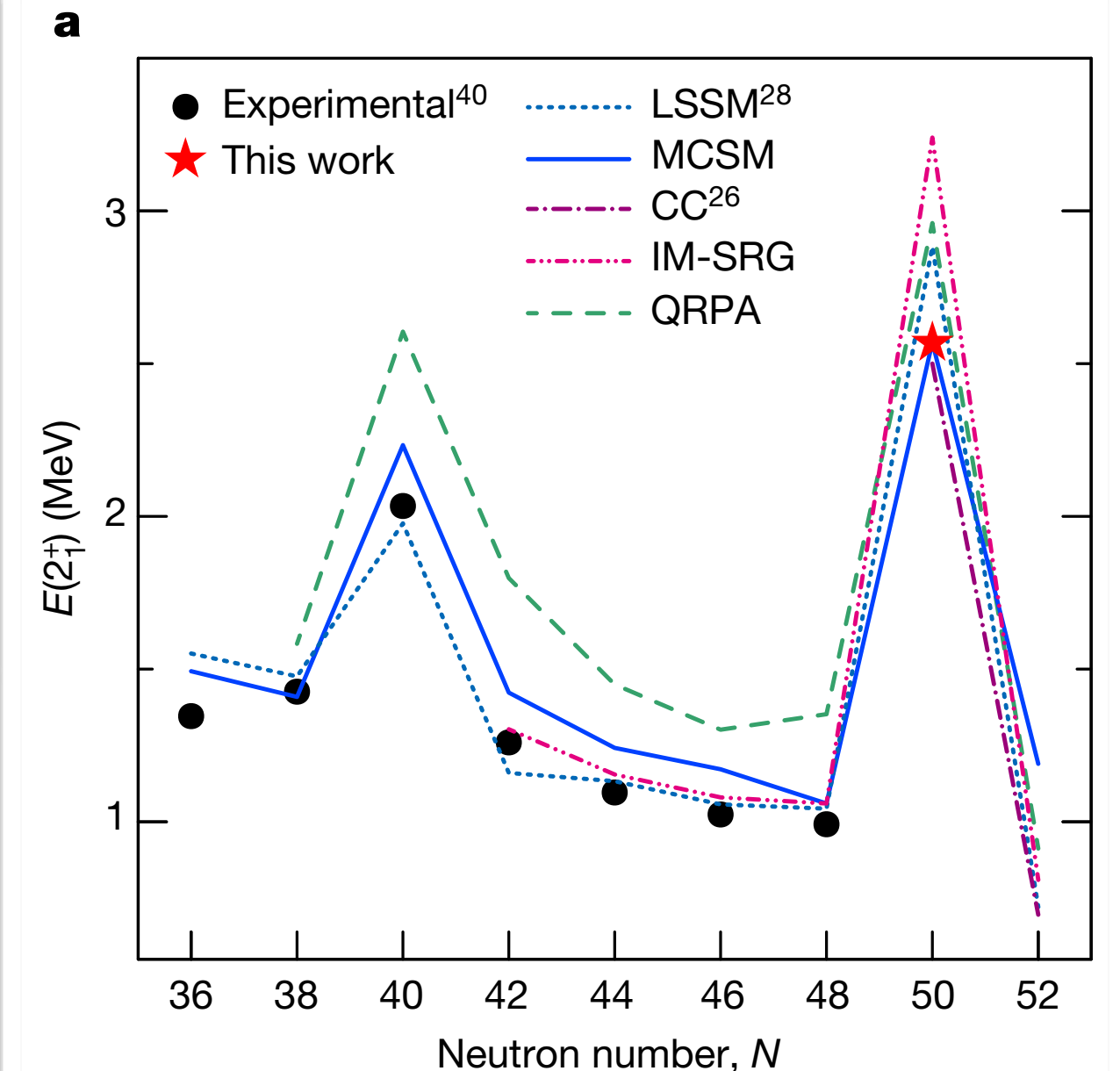
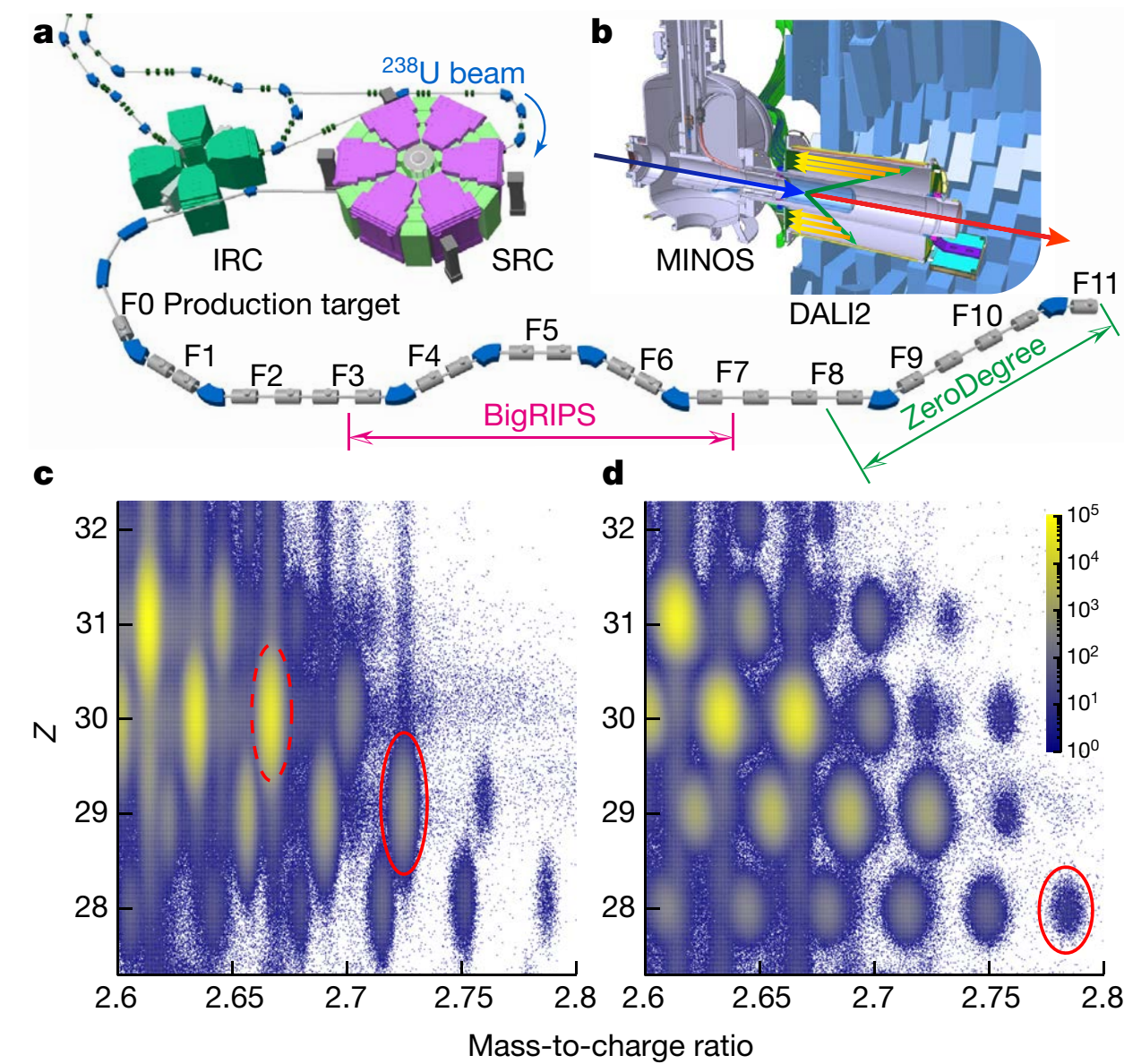
Magic systems still the pillars of our understanding

ARTICLE

<https://doi.org/10.1038/s41586-019-1155-x>

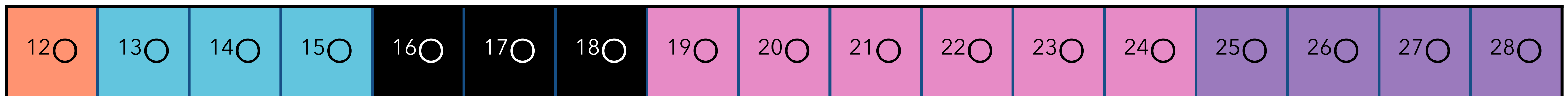
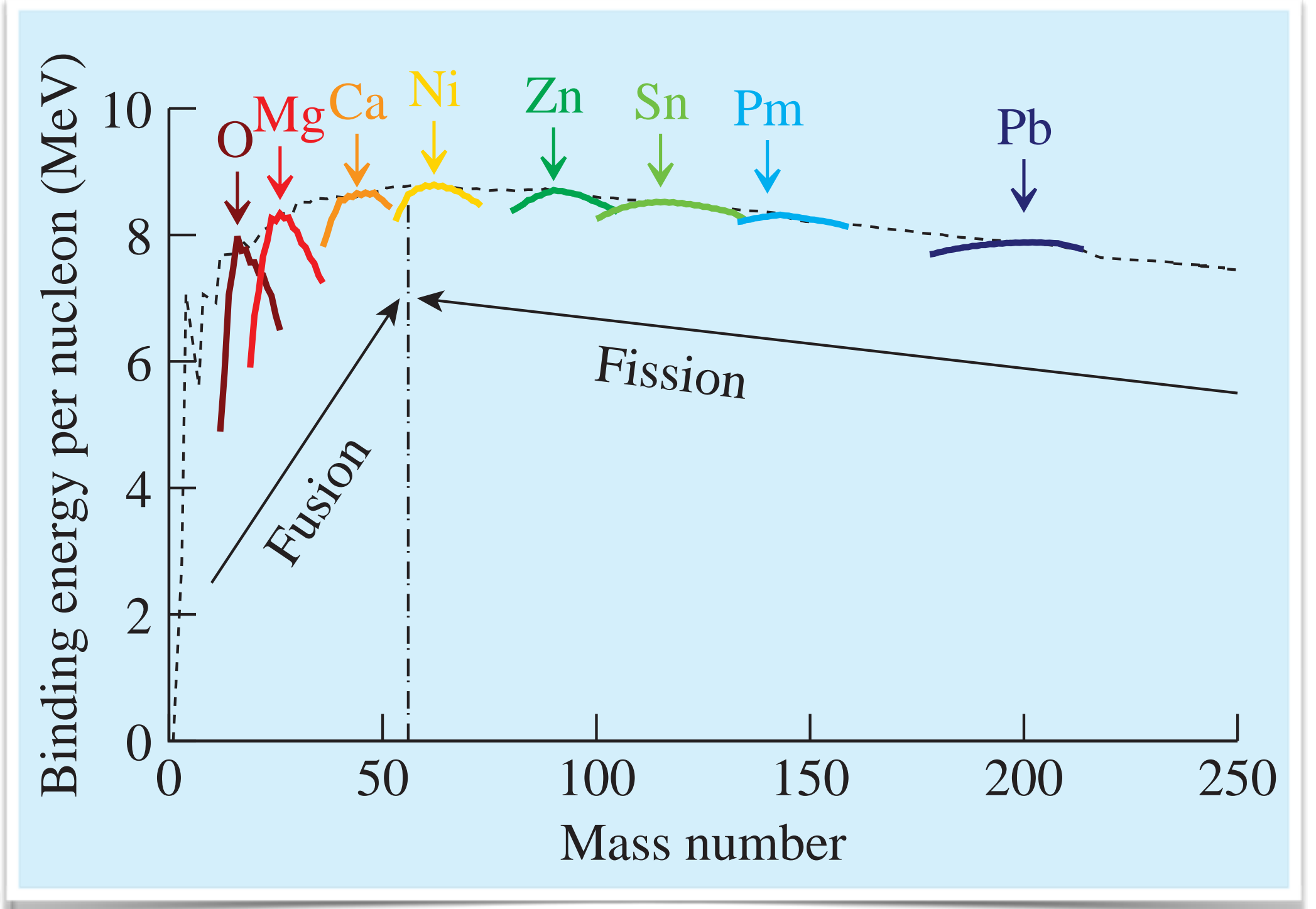
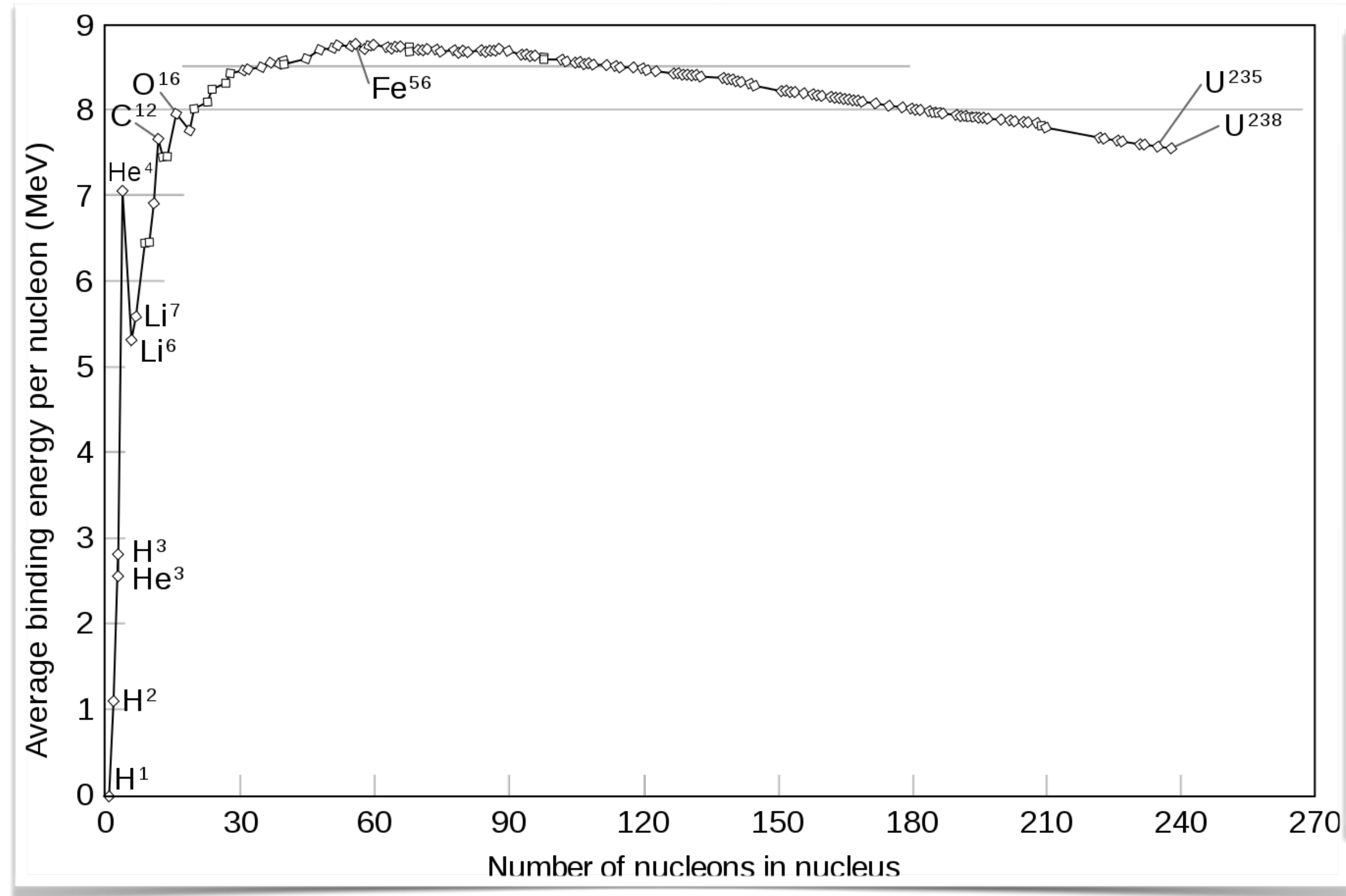
^{78}Ni revealed as a doubly magic stronghold against nuclear deformation

R. Taniuchi^{1,2}, C. Santamaria^{2,3}, P. Doornenbal^{2*}, A. Obertelli^{2,3,4}, K. Yoneda², G. Authélet³, H. Baba², D. Calvet³, F. Château³, A. Corsi³, A. Delbart³, J.-M. Gheller³, A. Gillibert³, J. D. Holt⁵, T. Isobe², V. Lapoux³, M. Matsushita⁶, J. Menéndez⁶, S. Momiyama^{1,2}, T. Motobayashi², M. Niikura¹, F. Nowacki⁷, K. Ogata^{8,9}, H. Otsu², T. Otsuka^{1,2,6}, C. Péron³, S. Péru¹⁰, A. Peyaud³, E. C. Pollacco³, A. Poves¹¹, J.-Y. Rousse³, H. Sakurai^{1,2}, A. Schwenk^{4,12,13}, Y. Shiga^{2,14}, J. Simonis^{4,12,15}, S. R. Stroberg^{5,16}, S. Takeuchi², Y. Tsunoda⁶, T. Uesaka², H. Wang², F. Browne¹⁷, L. X. Chung¹⁸, Z. Dombradi¹⁹, S. Franchoo²⁰, F. Giacoppo²¹, A. Gottardo²⁰, K. Hadyńska-Kle²¹, Z. Korkulu¹⁹, S. Koyama^{1,2}, Y. Kubota^{2,6}, J. Lee²², M. Lettmann⁴, C. Louchart⁴, R. Lozeva^{7,23}, K. Matsui^{1,2}, T. Miyazaki^{1,2}, S. Nishimura², L. Olivier²⁰, S. Ota⁶, Z. Patel²⁴, E. Şahin²¹, C. Shand²⁴, P.-A. Söderström², I. Stefan²⁰, D. Steppenbeck⁶, T. Sumikama²⁵, D. Suzuki²⁰, Z. Vajta¹⁹, V. Werner⁴, J. Wu^{2,26} & Z. Y. Xu²²

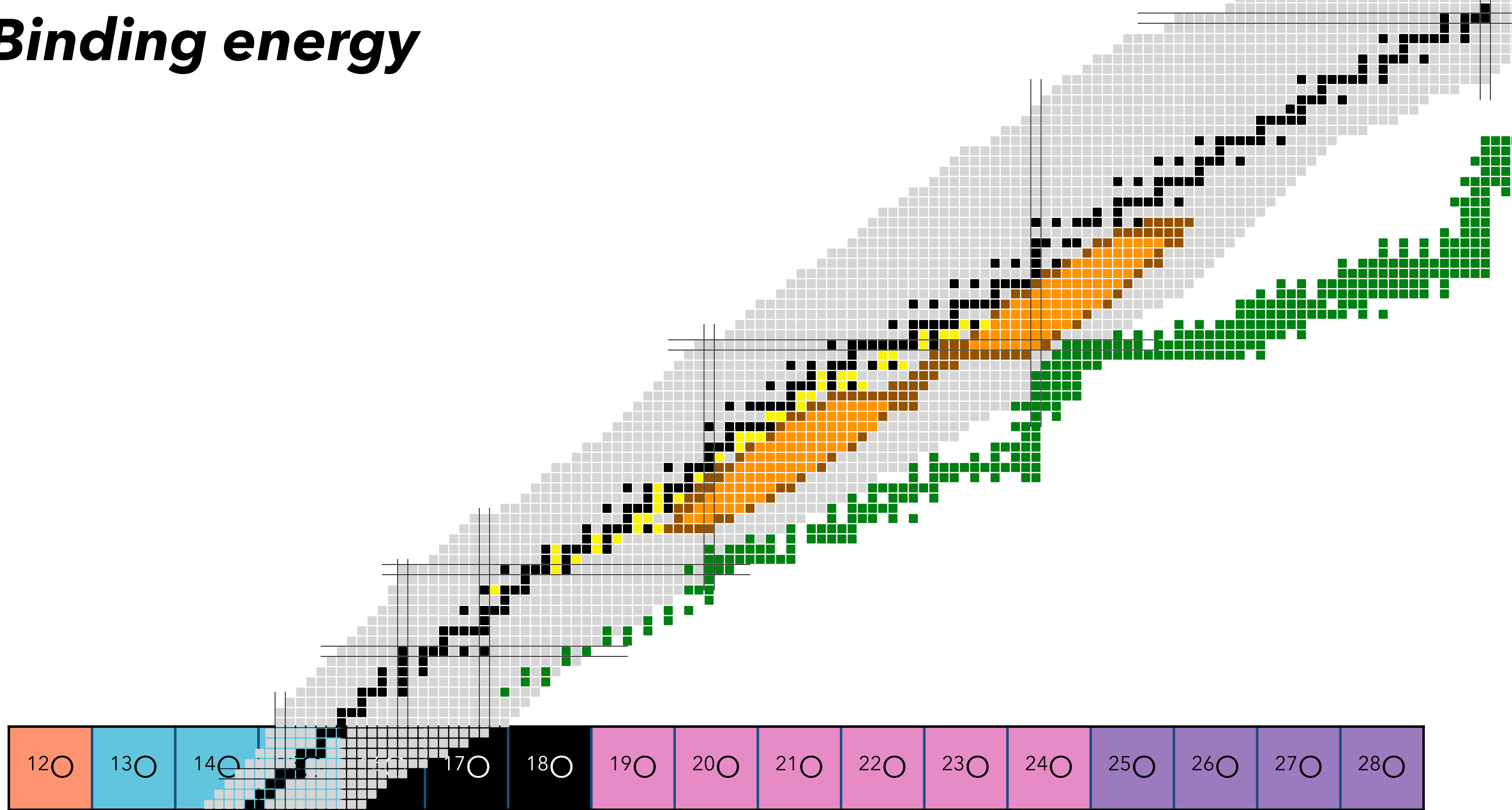


Masses

Binding energy



Binding energy



Wikimedia Commons (left) and NuDat2.0 (right)

Terms and data

International Atomic Energy Agency
Nuclear Data Services
Provided by the Nuclear Data Section

IAEA.org | NDS Mission | About Us | Mirrors: India | China | Russia

Databases » ENSDF | XUNDL | NuDat | LiveChart | NSR | Nuclear Wallet Cards | Related » ENSDF Manuals | Codes | Nuclear Data Sheets | EXFOR

Links
AMDC web site
Previous data
LiveChart
Q-values calculator

AMDC Atomic Mass Data Center

This page contains data provided by the **Atomic Mass Data Center**, located at the Institute of Modern Physics, Chinese Academy of Sciences (IMP), Lanzhou, China. Please refer to that web-site for further information about AME and NUBASE.

Atomic Mass Evaluation - AME2016

The evaluation has been published in *Chinese Physics C* 41 (2017) 030002 (PDF), 030003 (PDF).

The four main ASCII files of AME2016 are

1. [mass16.txt](#) - atomic masses
2. [mass16round.txt](#) - atomic masses "rounded" version
3. [rct1-16.txt](#) - reaction energies, table 1
4. [rct2-16.txt](#) - reaction energies, table 2

Additionally, the [covariance](#) zip file is downloadable.

Any work that will use these files should make reference to these paper and not to the electronic files.

A [Q-Value](#) Calculator is available, using AME2016 masses.

Nuclear mass: $m(Z,N)$

Binding energy per nucleon

Mass excess: $\{Zm_p + Nm_n - m(Z,N)\}c^2$

Separation energies, S_p and S_n ($S_{2p} + S_{2n}$)

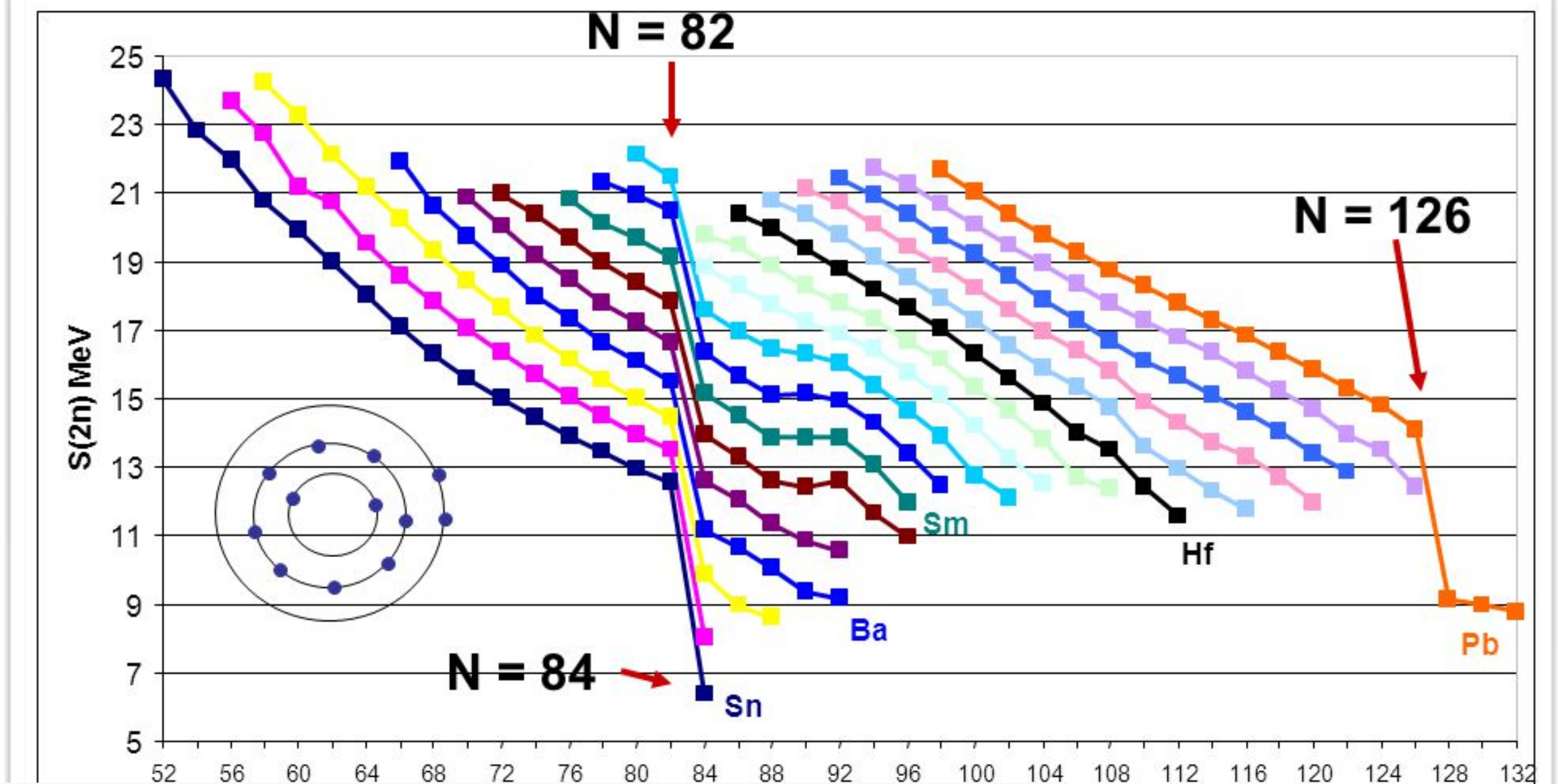
[difference in binding energies e.g. $S_n = \{m(Z,N-1) + m_n - m(Z,N)\}c^2 = B(Z,N) - B(Z,N-1)$]

Atomic mass unit (mass of nucleon): $931.49410242(28) \text{ MeV}/c^2$

Mass databases compile these ...

e.g. <https://www-nds.iaea.org/amdc/>

Terms and data



MASS LIST
for analysis

1N-Z	N	Z	A	EL	O	MASS EXCESS (keV)	BINDING ENERGY/A (keV)	BETA-DECAY ENERGY (keV)	ATOMIC MASS (micro-u)	V/S
0 1	1	0	1	n		8071.31714	0.00046	0.0	1 008664.91585	0.00049
-1	0	1	1	H		7288.97059	0.00009	0.0	1 007825.03223	0.00009
0 0	1	1	2	H		13135.72174	0.00011	1112.283	2 014101.77812	0.00012

LINEAR COMBINATIONS

ADAM STYLE(1)
for analysis

1 A	elt	Z	S(2n)	S(2p)	Q(a)	Q(2B-)	Q(ep)	Q(B- n)
0 1	n	0	*	*	*	*	*	*
1	H	1	*	*	*	*	*	*
0 2	H	1	*	*	*	*	*	*
0 3	H	1	8481.80	0.00	*	-13717#	2000#	*

Various techniques to determine mass

- **Q-values**

- Decays
- Kinematics

$$Q = \sum m_i - \sum m_f$$

- **ToF**

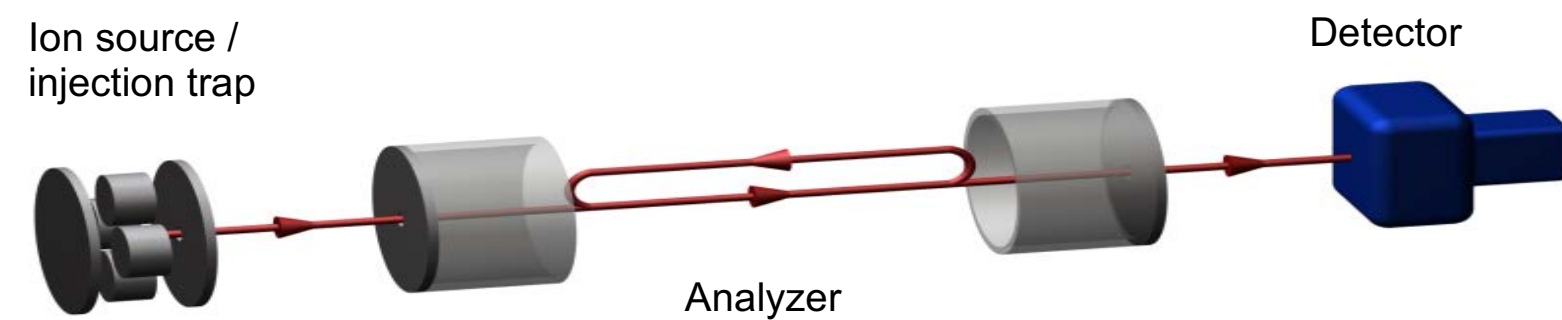
- Spectrograph

$$qvB = \frac{(\gamma m_0)(L_{path}/ToF)^2}{\rho}$$

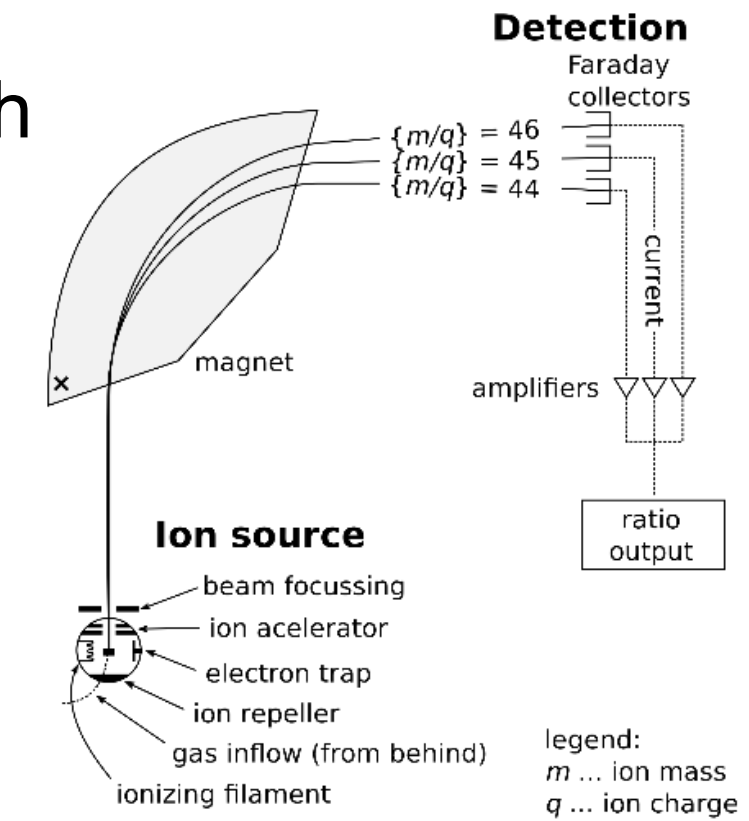
$$\Rightarrow m_0 = \frac{ToF}{L_{path}} \frac{qB\rho}{\gamma}$$

- Multi-reflection (MR-TOF)

$$E = \frac{1}{2}mv^2 = qeU \Rightarrow \frac{m}{q} \propto t^2$$



- **Dispersion**
– Spectrograph

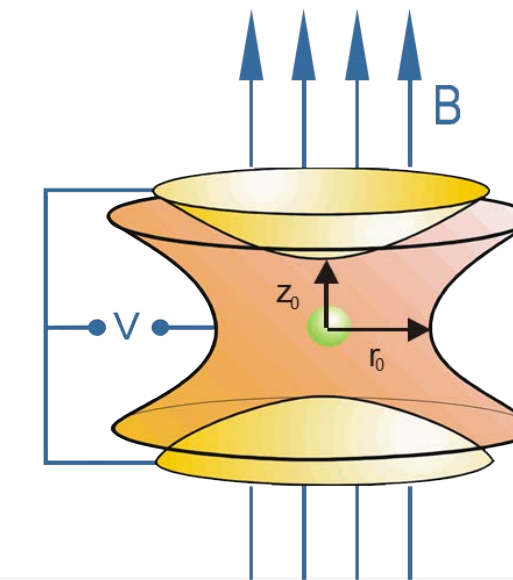
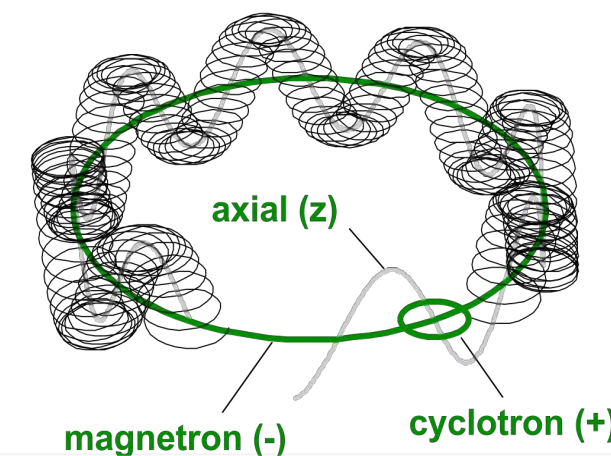


- **Frequency**

- Penning trap
- Storage rings

Cyclotron Frequency

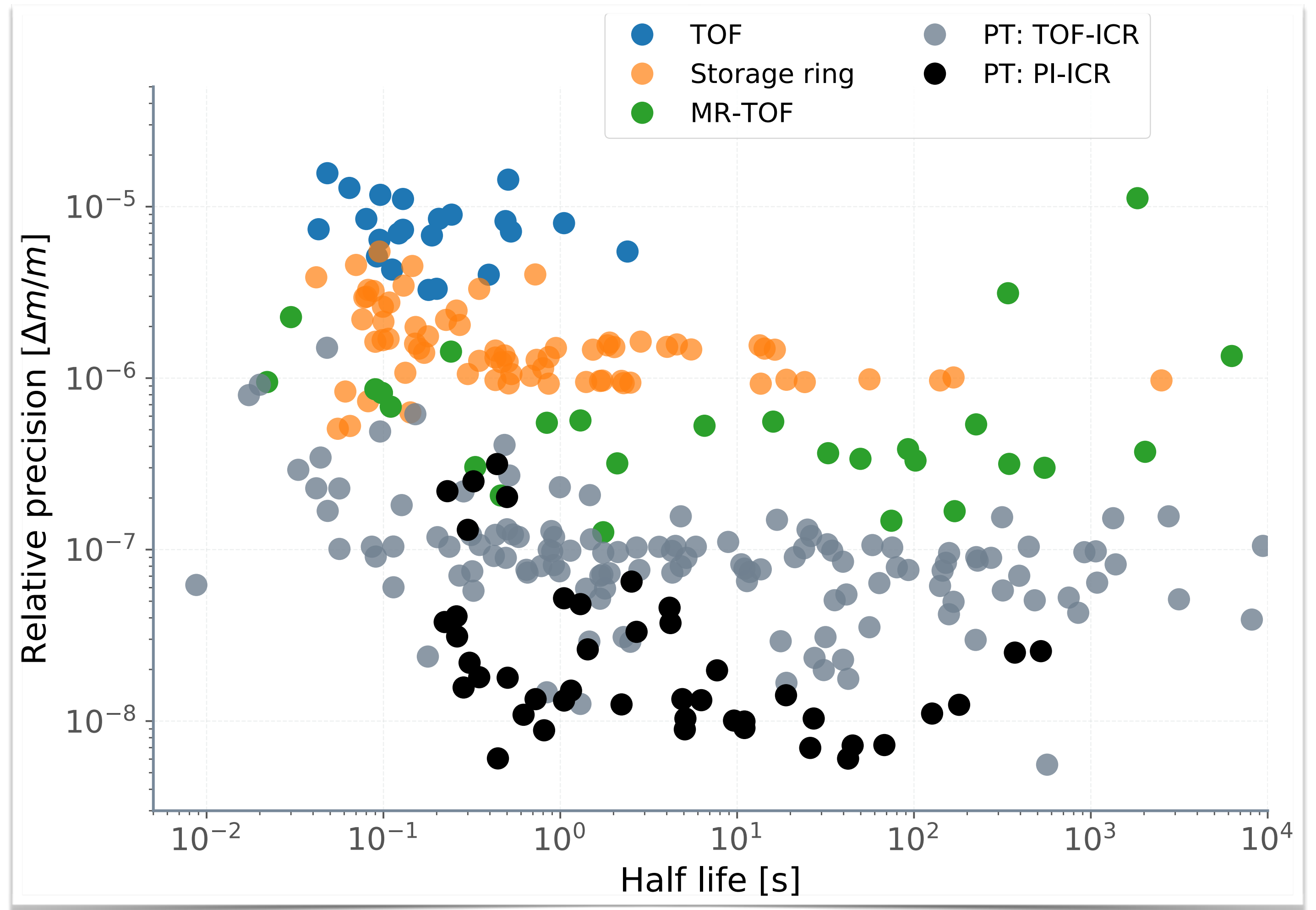
$$\omega_c = \frac{1}{2\pi} \frac{q}{m} B$$



Precision, time, resolution

Modern techniques:

- *TOF (fast, low precision)*
- *Storage rings (fast, many measurements at once)*
- *MR-TOFs (fast, high resolution)*
- *Penning traps ("slow", high resolution, high precision)*



e.g. Q value

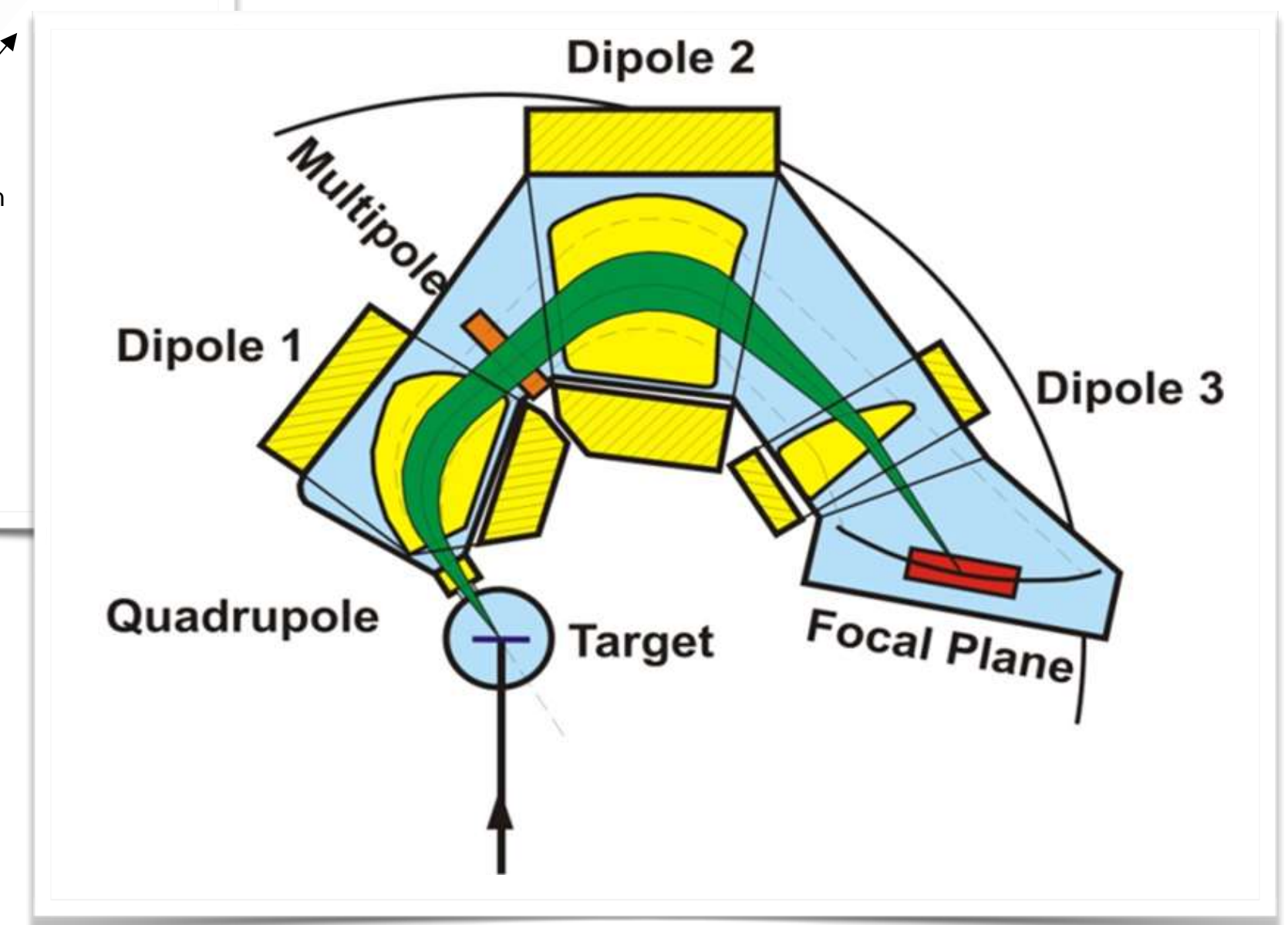
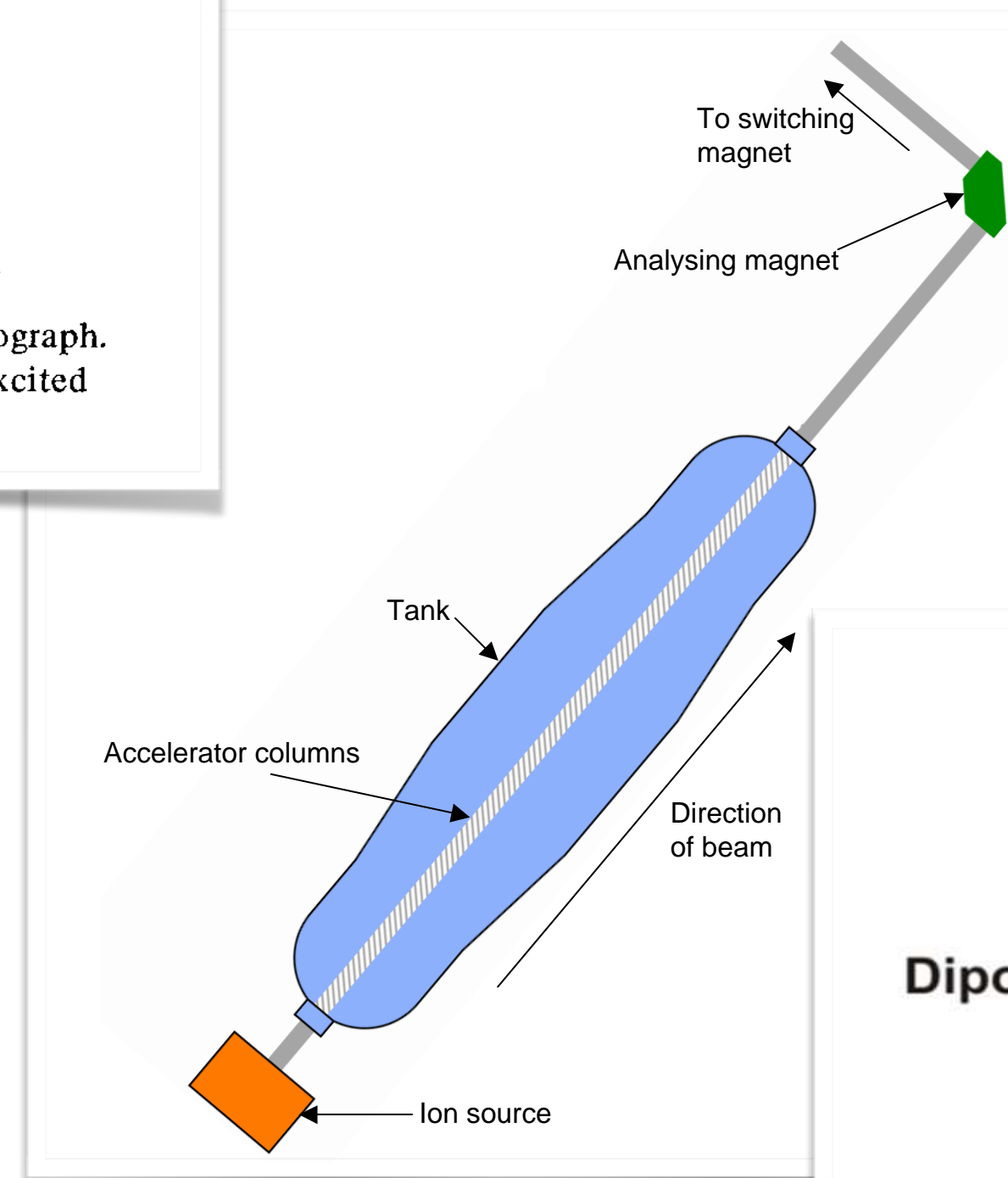
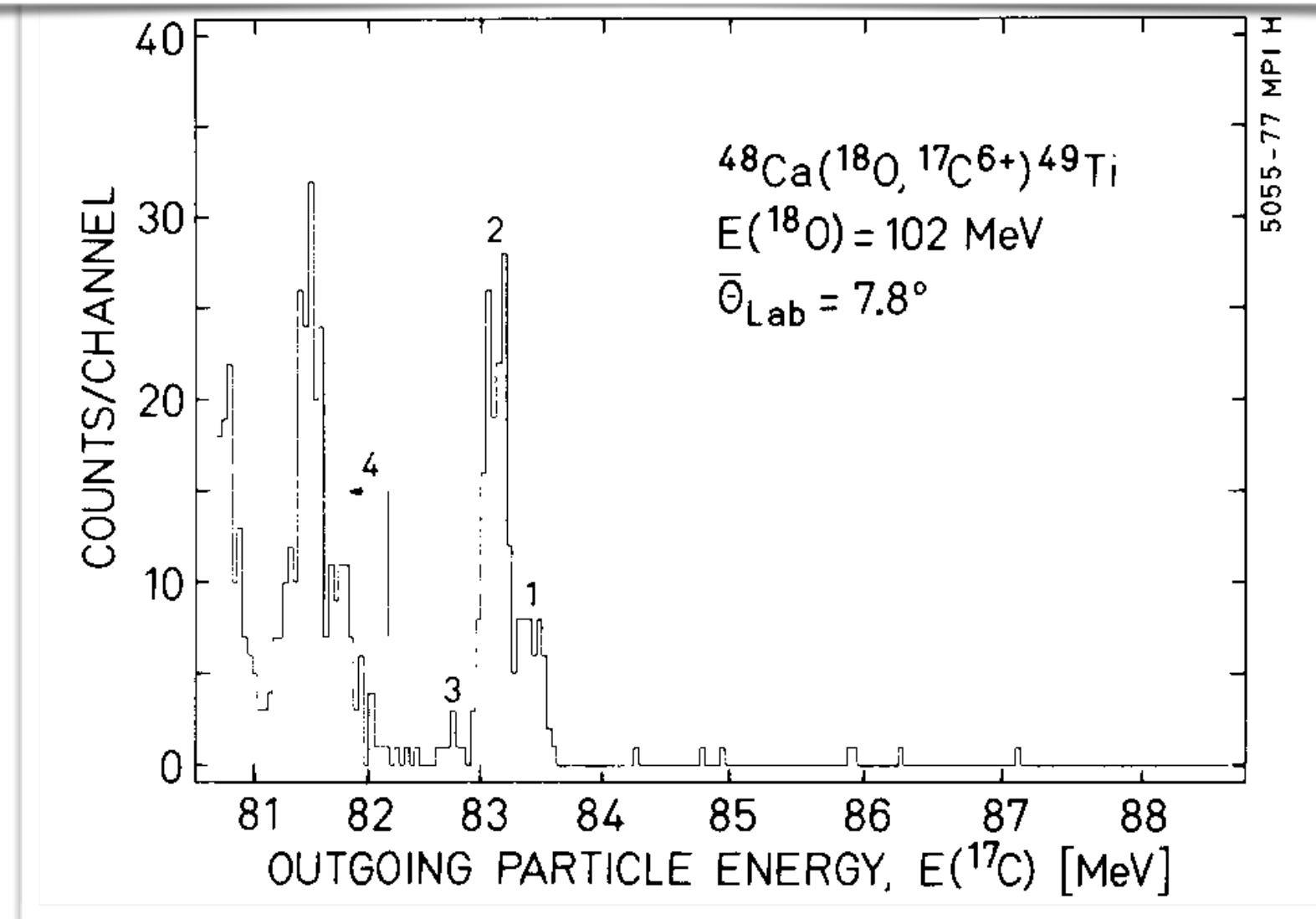
THE MASS AND LOW-LYING LEVEL STRUCTURE OF ^{17}C

J.A. NOLEN¹, T.S. BHATIA², H. HAFNER³, P. DOLL⁴,
C.A. WIEDNER and G.J. WAGNER

Max-Planck-Institut für Kernphysik, Heidelberg, Germany

Received 14 September 1977

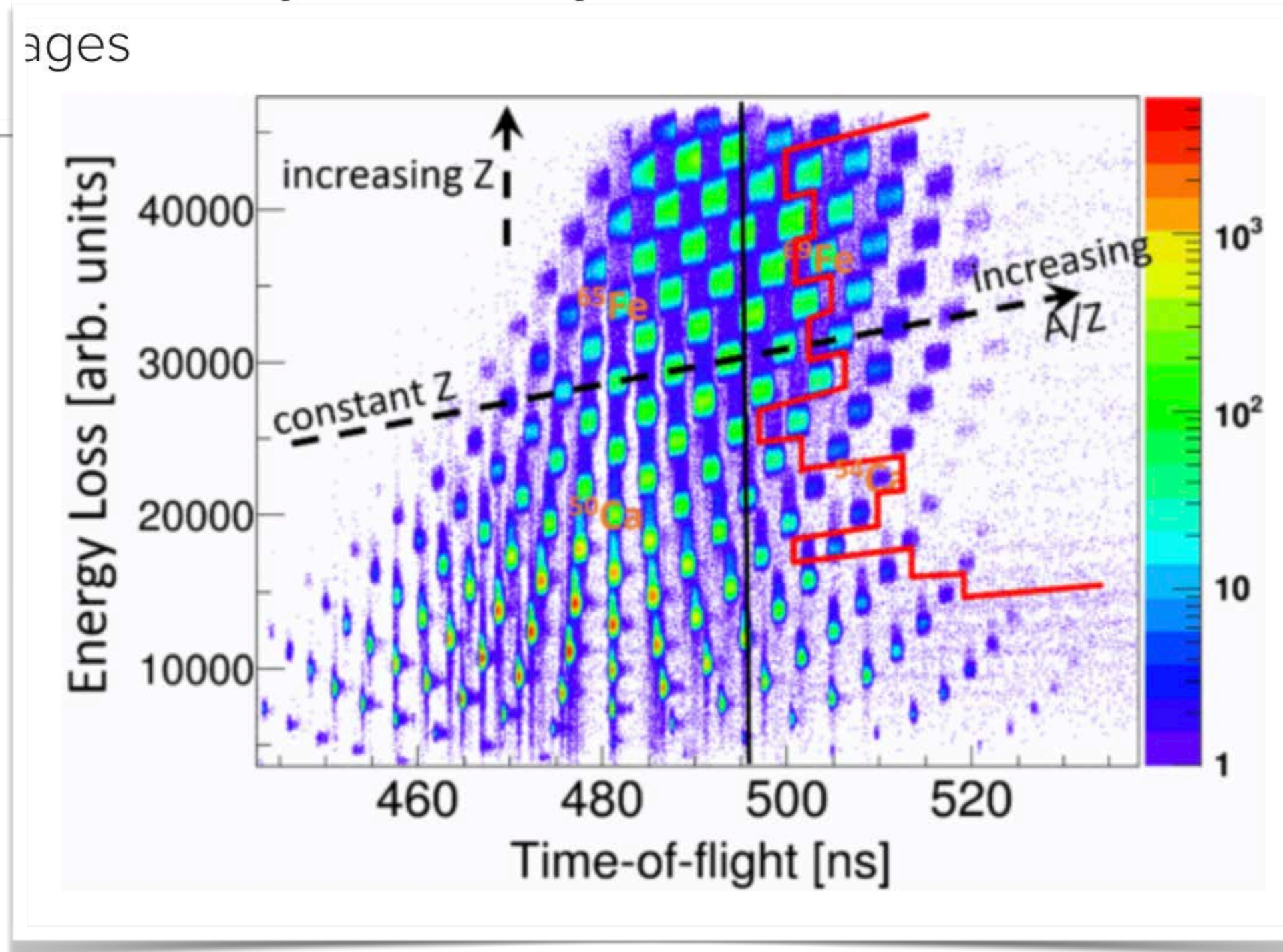
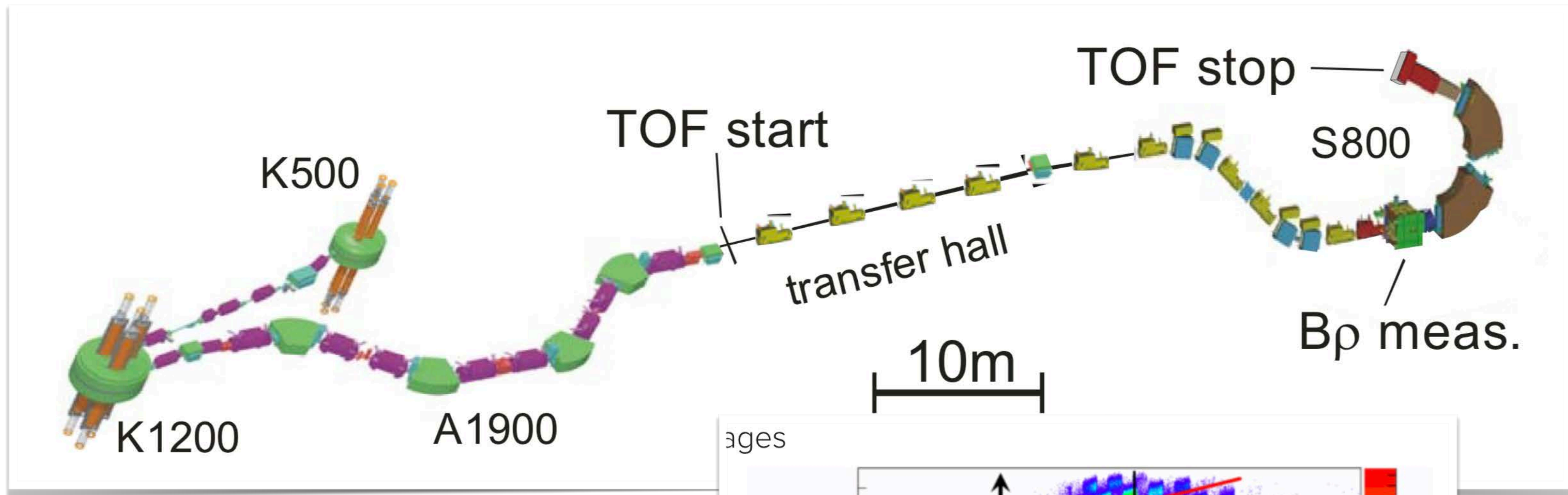
Spectra of ^{17}C ions from the $^{48}\text{Ca}(^{18}\text{O}, ^{17}\text{C})^{49}\text{Ti}$ reaction at 102 MeV were recorded with a Q3D spectrograph. The lowest state observed, assumed to be the ground state of ^{17}C , has a mass excess of 21023 ± 35 keV. An excited state of ^{17}C was also observed at an excitation energy of 292 ± 20 keV.



$$S_n = \{m(Z, N-1) + m_n - m(Z, N)\}c^2 = B(Z, N) - B(Z, N-1) = Q$$

(with knowledge of other masses and beam energies)

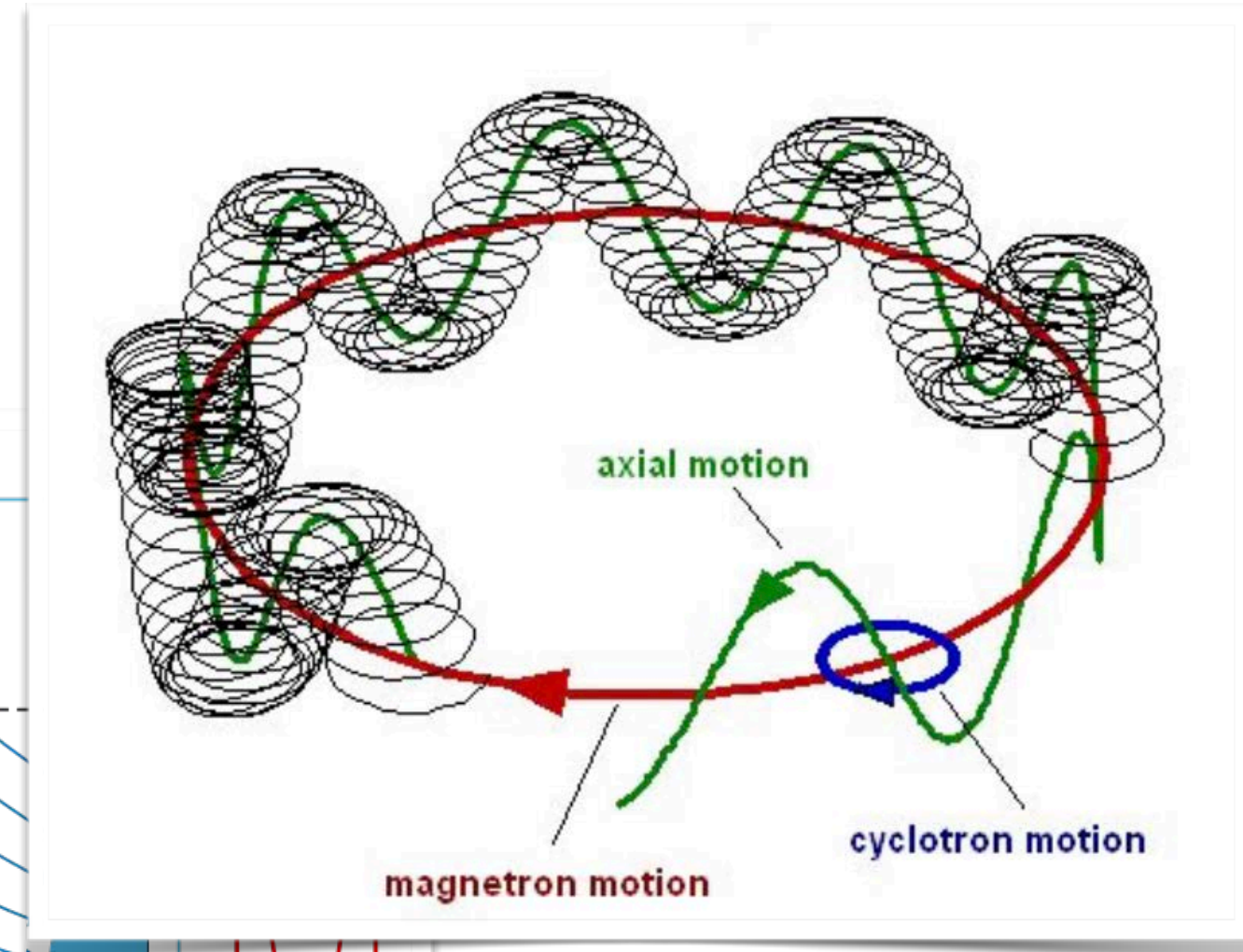
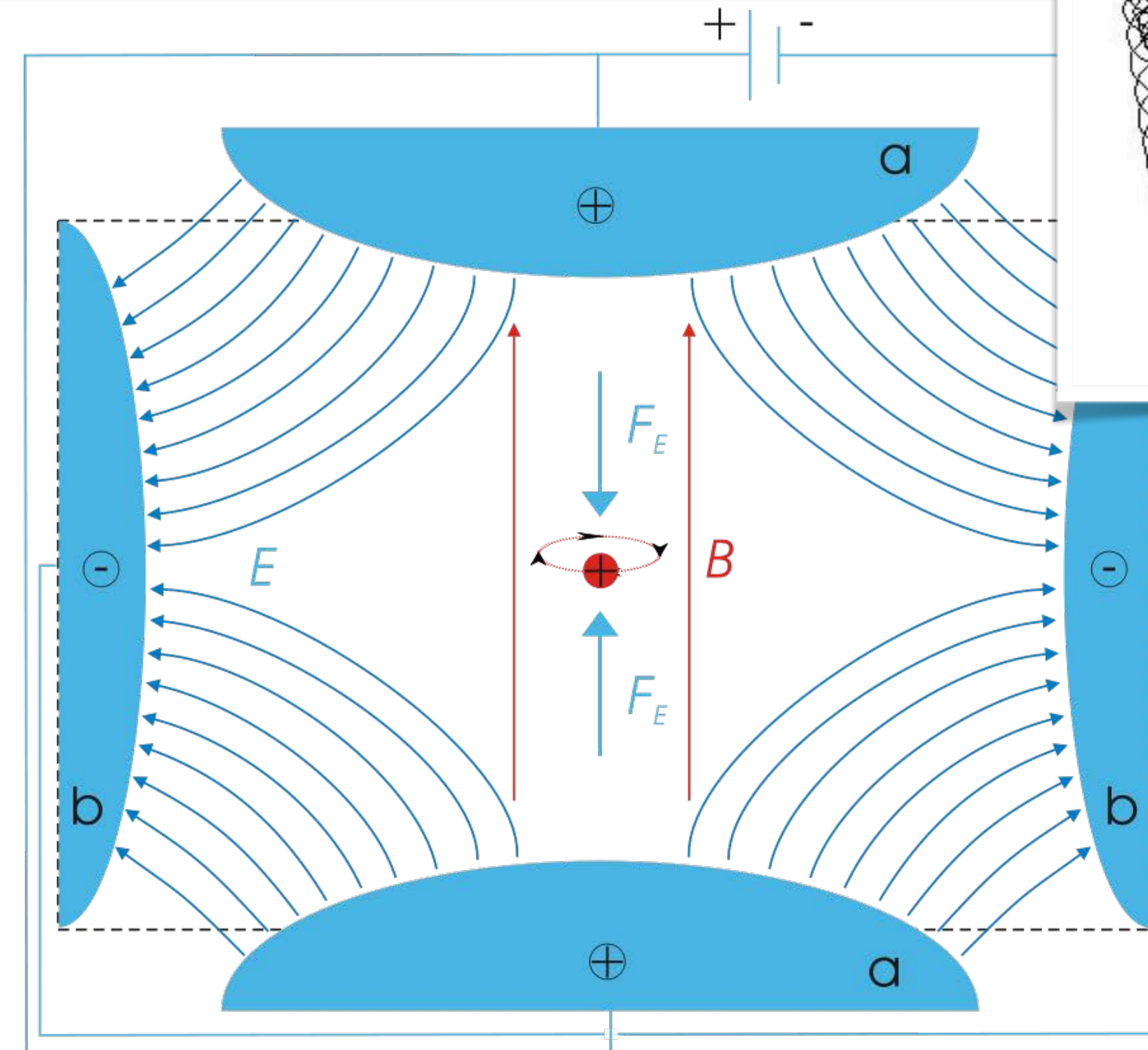
e.g. TOF (and magnetic spectrograph)



e.g. Z. Meisel, Phys. Rev. C **93**, 035805 (2016)

e.g. Penning traps

A somewhat de-facto approach for 30+ years



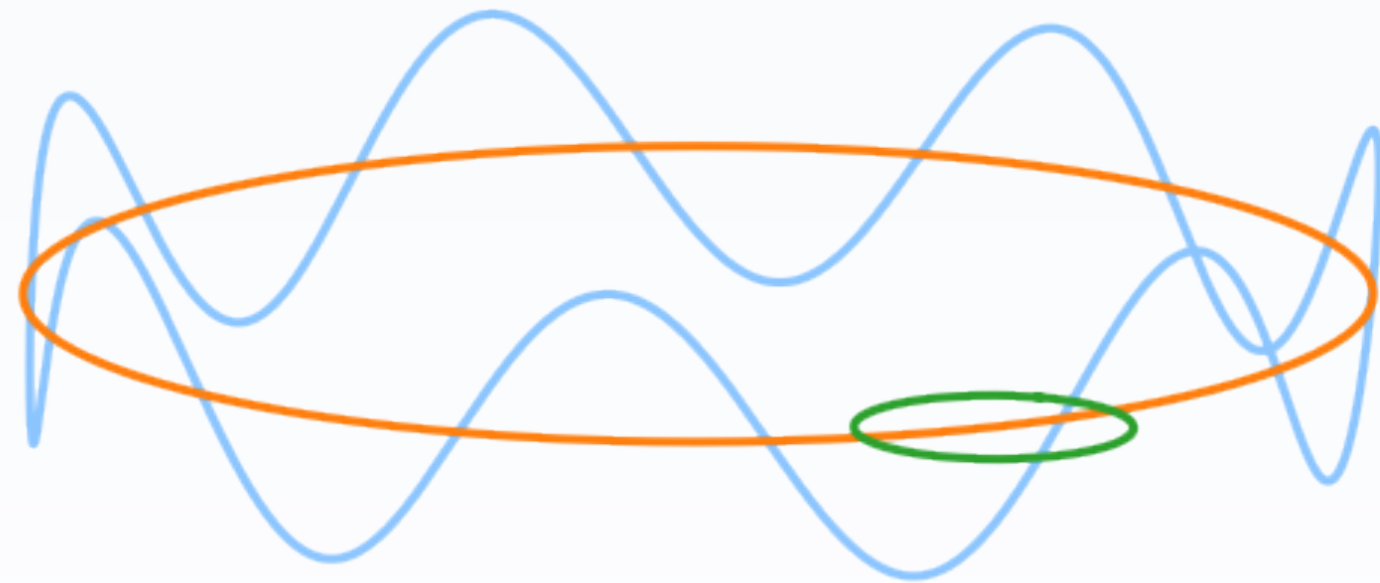
$$\omega_c = \frac{qB}{m_{ion}}$$

Ion confined in strong B field (radial confinement), with electrodes providing a potential for axial confinement and manipulation

Ion motion in Penning trap

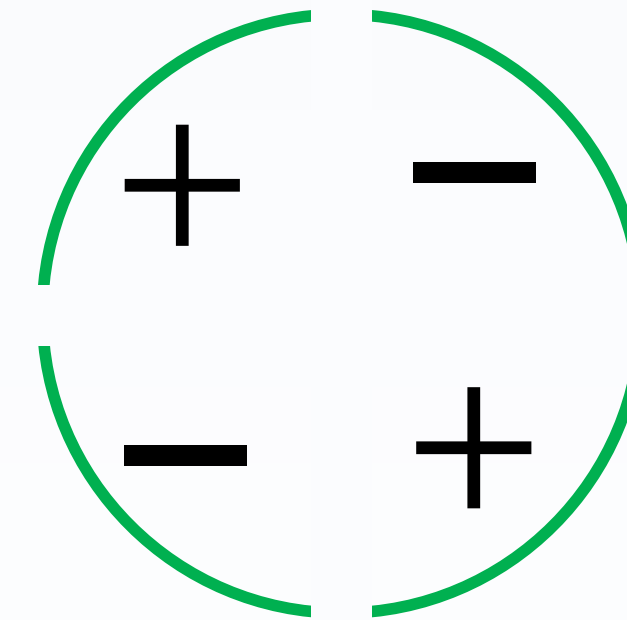
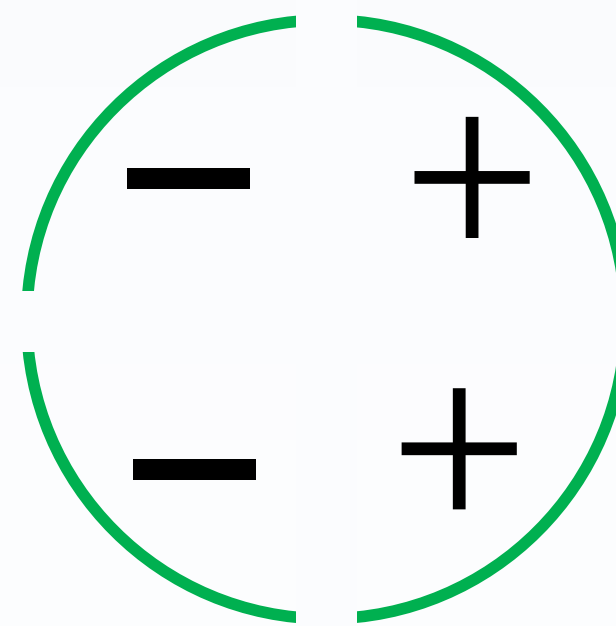
Penning trap motion

— ω_z — ω_- — ω_+



Segmented Ring electrode allows for:

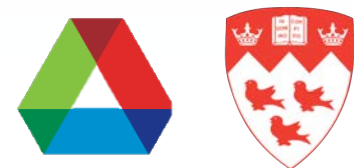
1. Driving of orbital motions through a dipole excitation
2. Interconversion of orbital motions through a quadrupole excitation at ω_c



$$\omega_c = \omega_- + \omega_+$$

$\omega_- \sim 1.6$ kHz
Mass independent

$\omega_+ \sim 1$ MHz
Mass dependent



Nuclear forum - LBNL - R. Orford

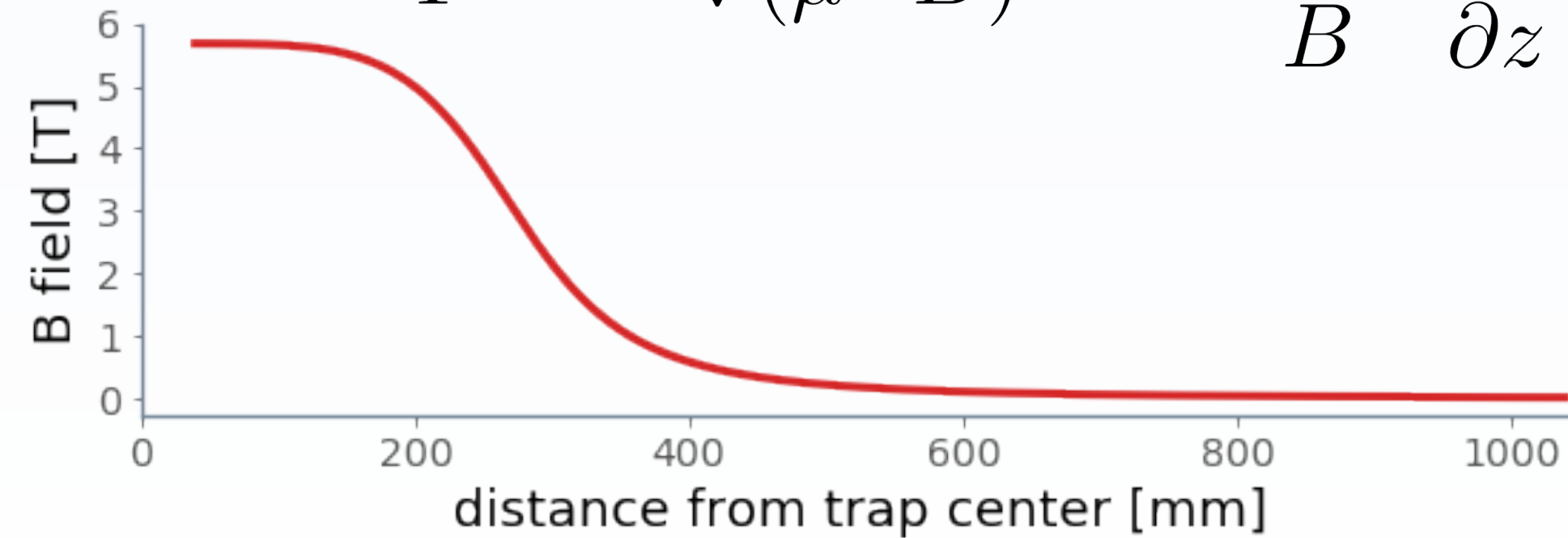
22

Penning trap mass measurements

TOF drift tube



$$\vec{F} = -\nabla(\mu \cdot \vec{B}) = -\frac{E_{rad}}{B} \frac{\partial B}{\partial z} \hat{z}$$



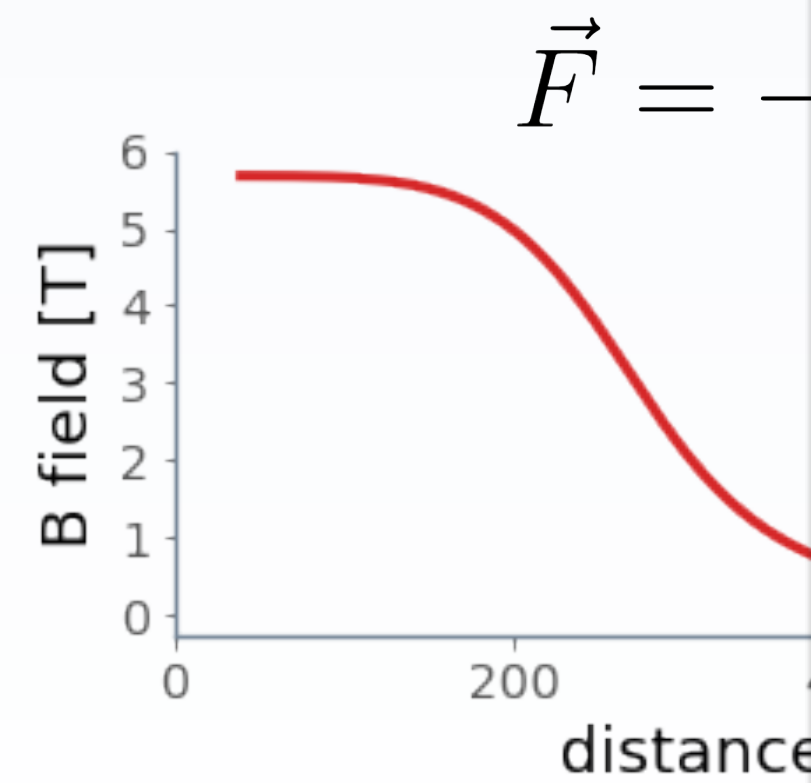
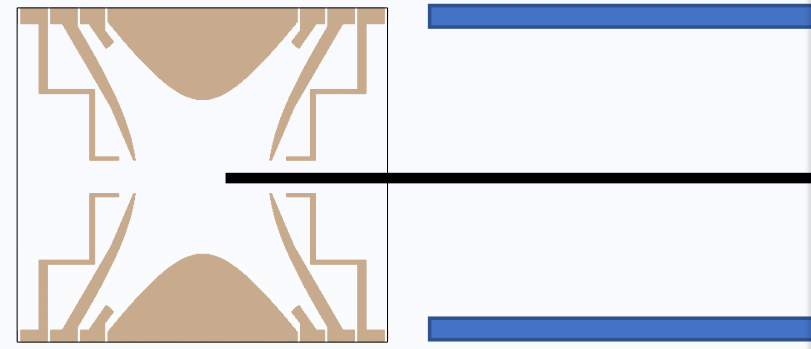
$$\omega_+ \gg \omega_-$$

Leads to Time-Of-Flight Ion-Cyclotron-Resonance (TOF-ICR) method of mass measurements

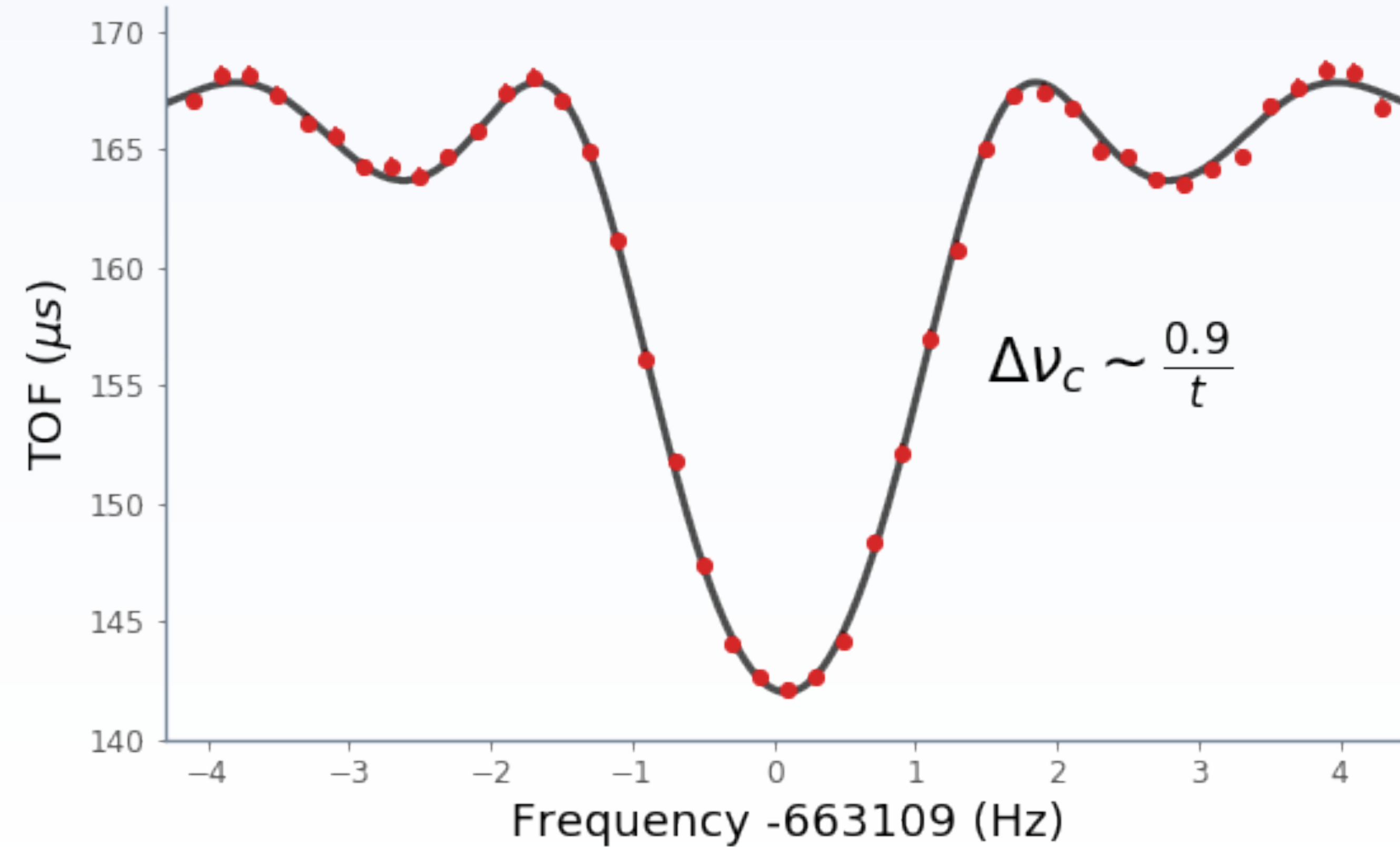


Penning trap mass measurements

TOF drift tube



TOF-ICR Method

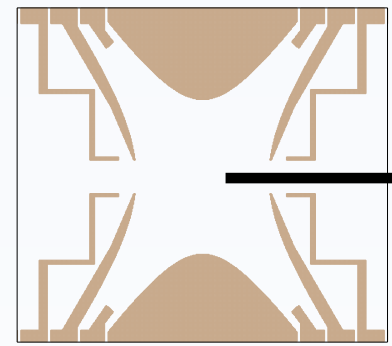


Nuclear forum - LBNL - R. Orford

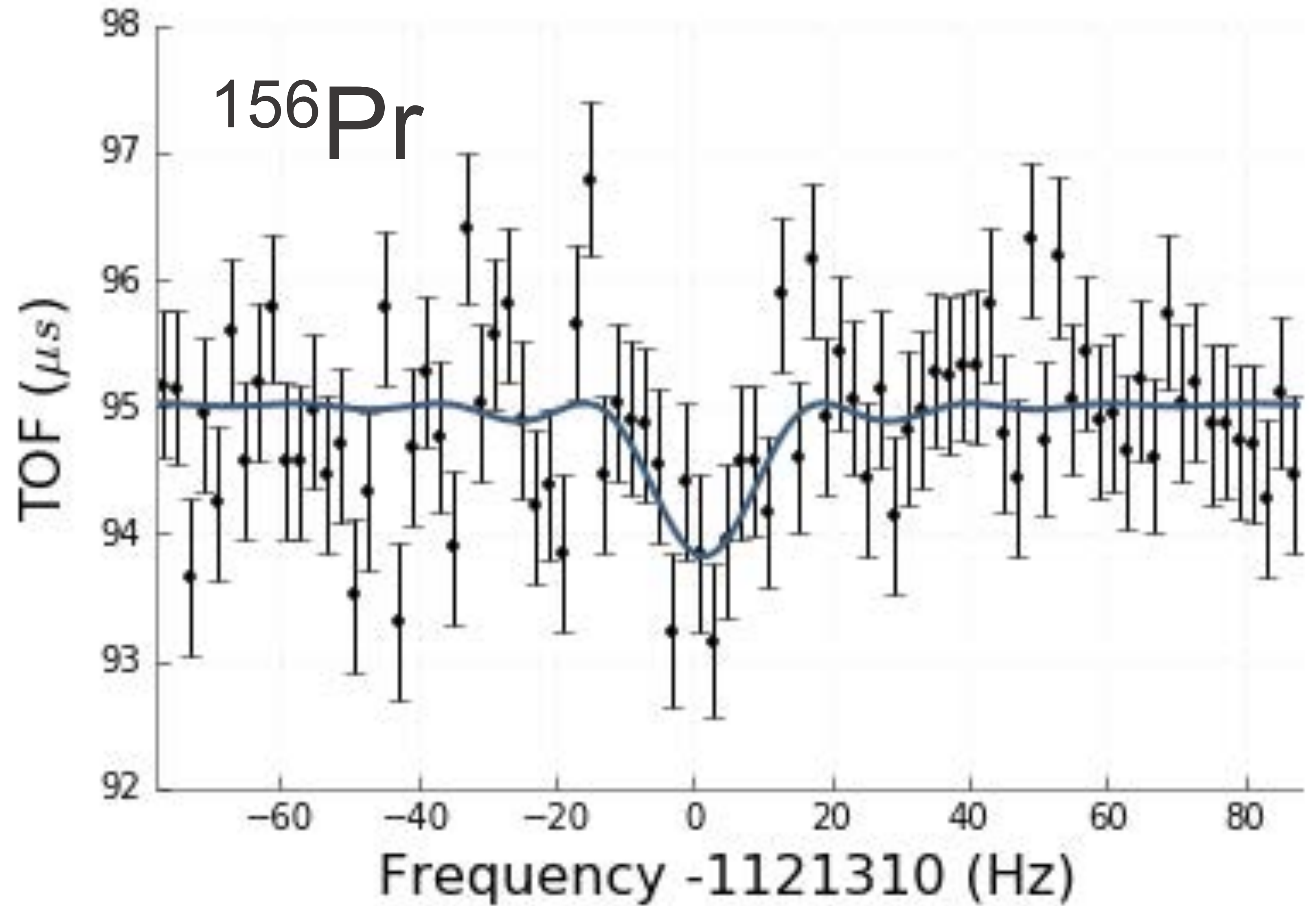
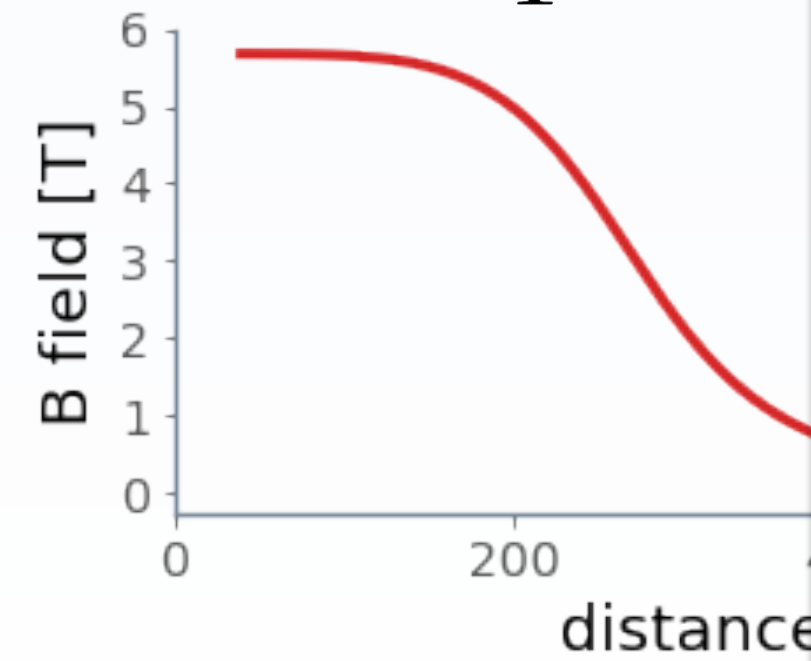
24

Penning trap ma

TOF drift



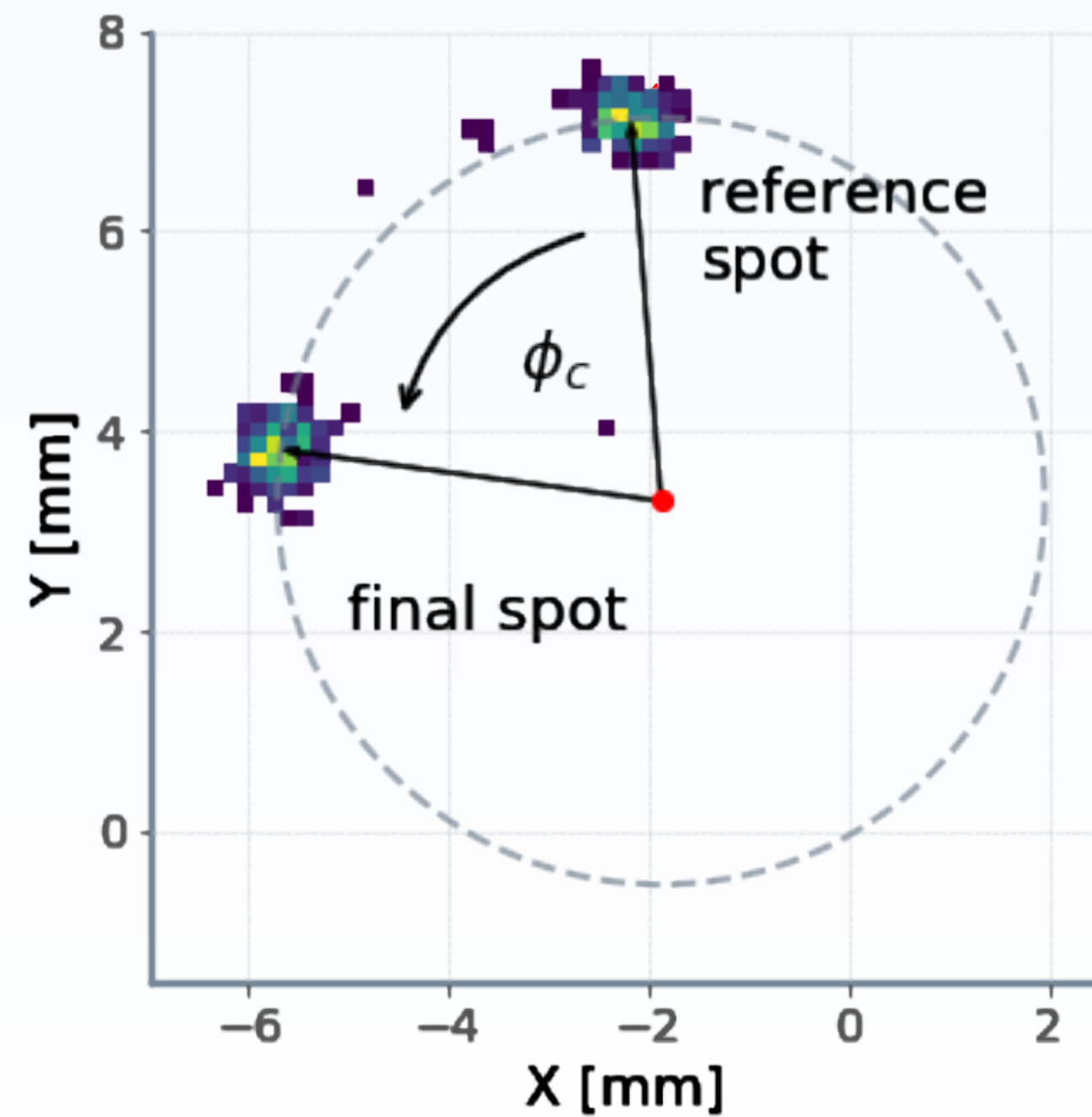
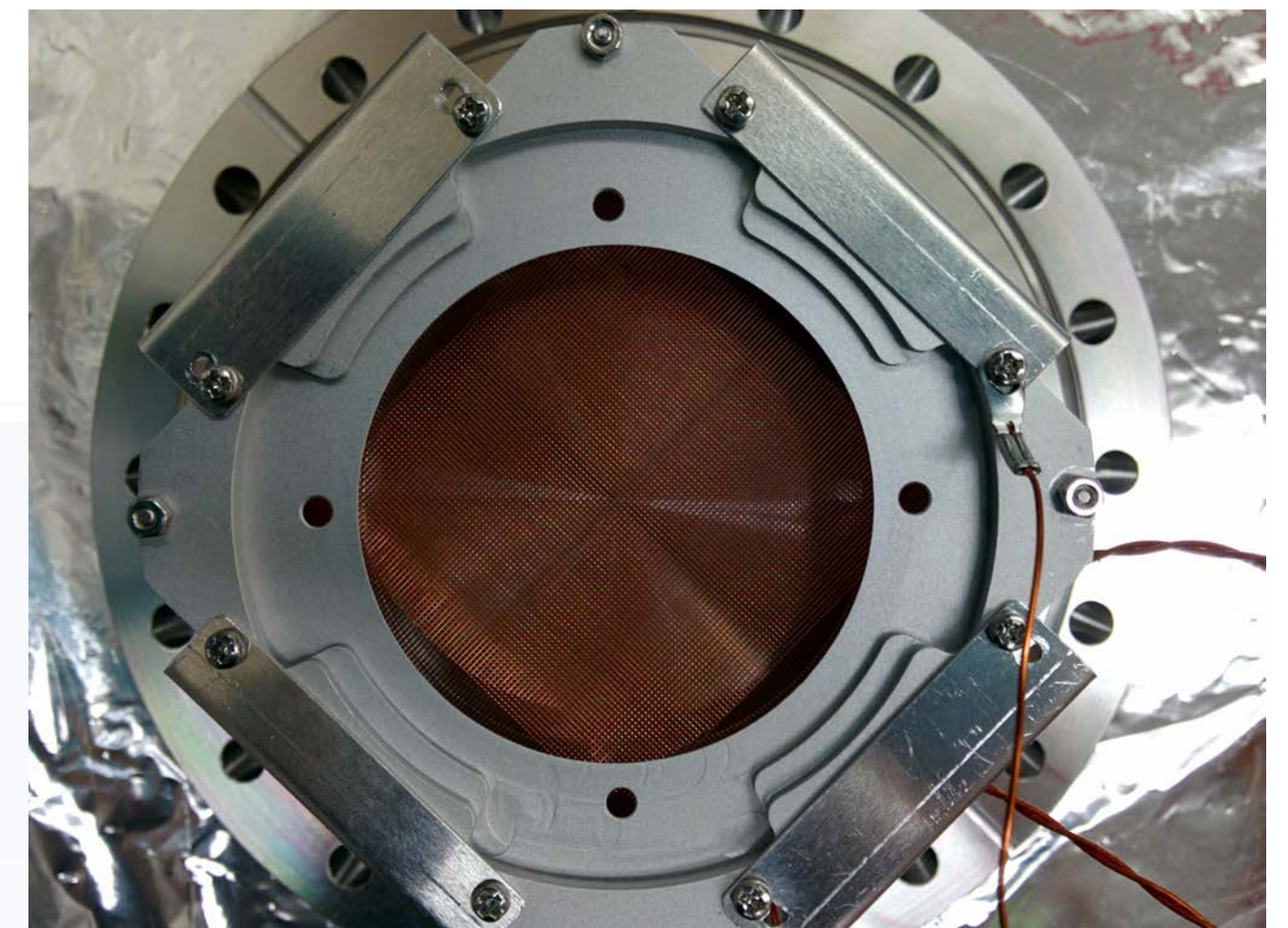
$$\vec{F} = -$$



Nuclear forum - LBNL - R. Orford

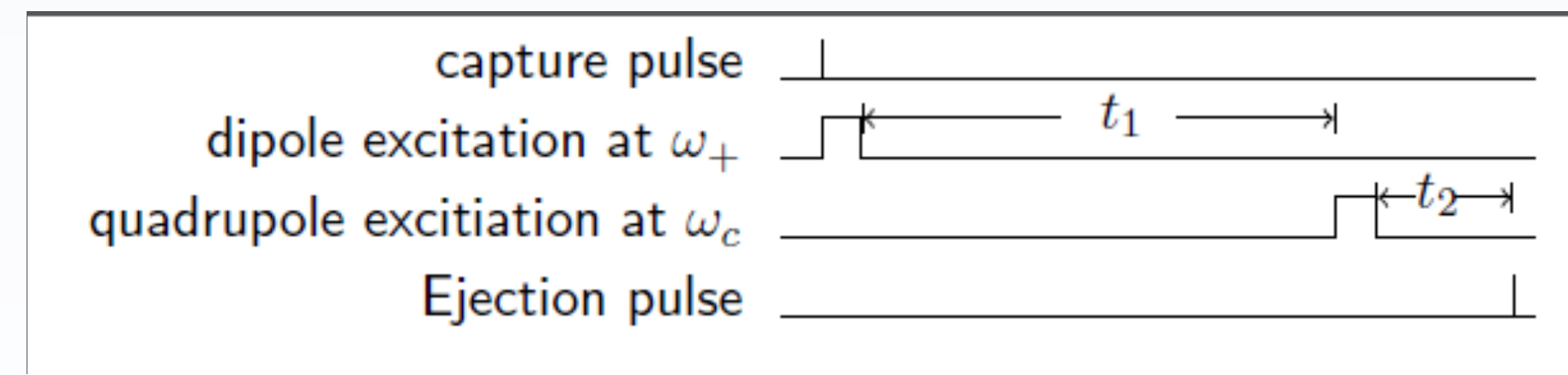
24

PI-ICR technique

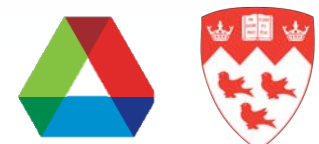


General concept: Determine the cyclotron frequency of trapped ions by measuring the phase advance of ions over a period excitation free motion

- Use a position-sensitive detector to determine ion position

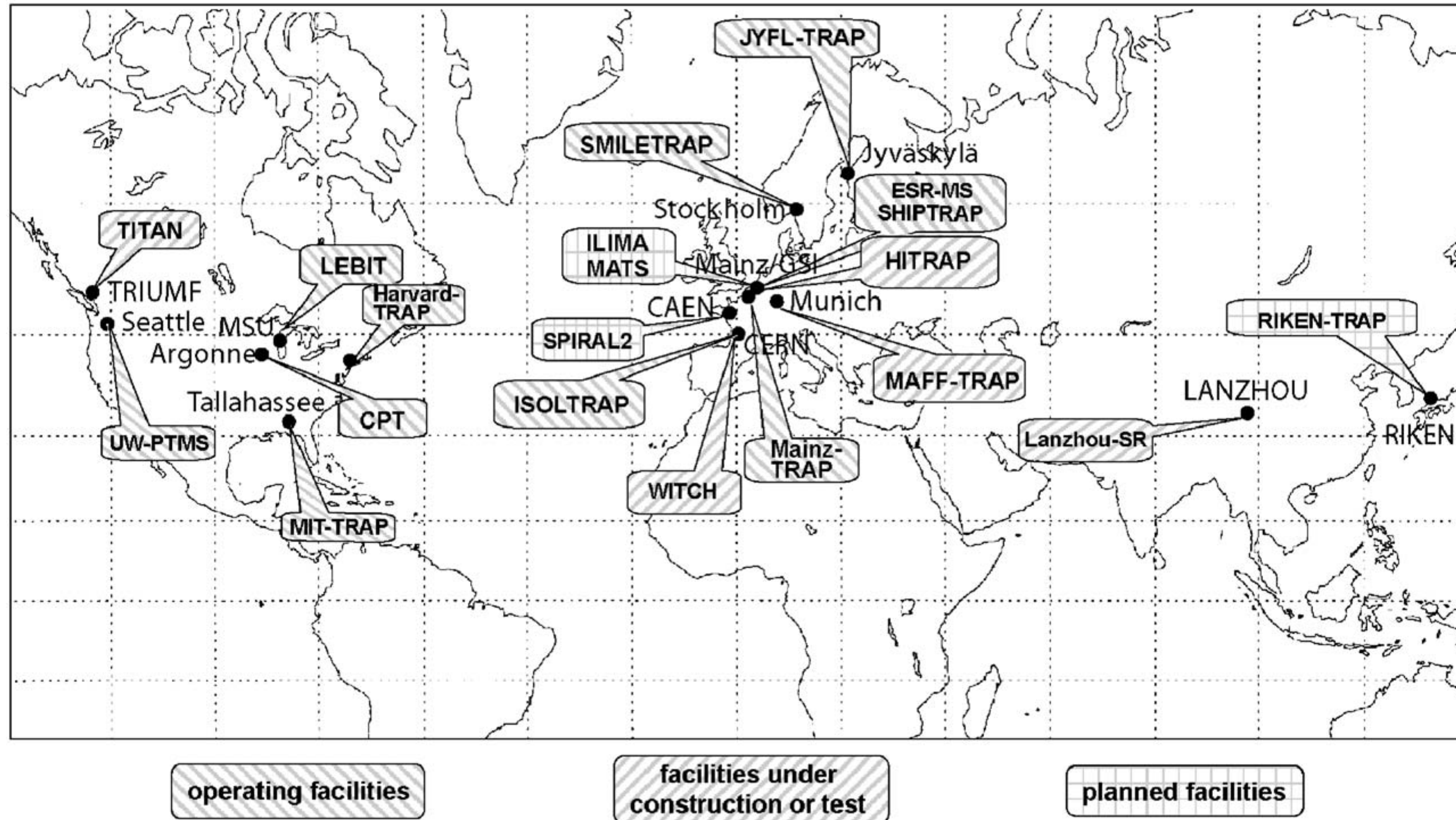


$$\omega_c = \frac{qB}{m_{ion}} = \frac{\phi_c + 2\pi N}{t_1}$$



e.g. Penning traps

WORLDWIDE EXPLOSION OF PENNING TRAP PROJECTS



Penning traps ... popular? (last 10 days)

PHYSICAL REVIEW C **100**, 014304 (2019)

Mass measurements of neutron-rich isotopes near $N = 20$ by in-trap β decay with the ISOLTRAP spectrometer

P. Ascher,^{1,*} N. Althubiti,^{2,3} D. Atanasov,^{4,†} K. Blaum,⁴ R. B. Cakirli,⁵ S. Grévy,¹ F. Herfurth,⁶ S. K. V. Manea,^{8,†} D. Neidherr,⁶ M. Rosenbusch,⁹ L. Schweikhard,¹⁰ A. Welker,¹¹ F. Wienholtz,¹⁰ R. N. Woerner,¹¹

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⁵University of Istanbul, Department of Physics 34134, Istanbul, Turkey

⁶GSI Helmholtzzentrum für Schwerionenforschung GmbH, Darmstadt, Germany

⁷CSNSM-IN2P3-CNRS, Université de Paris Sud, Orsay, France

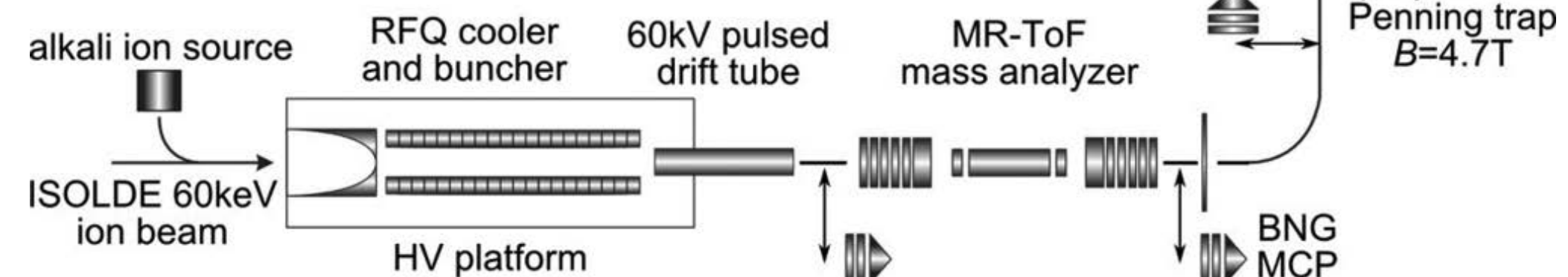
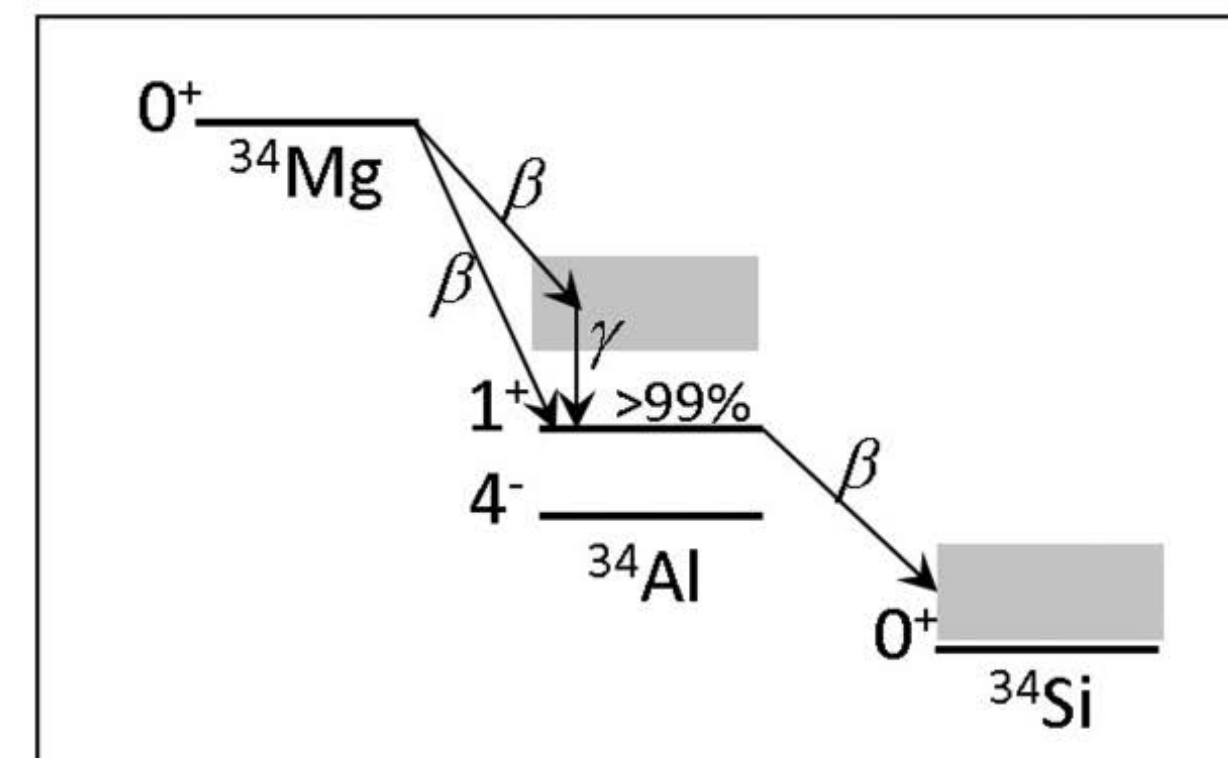
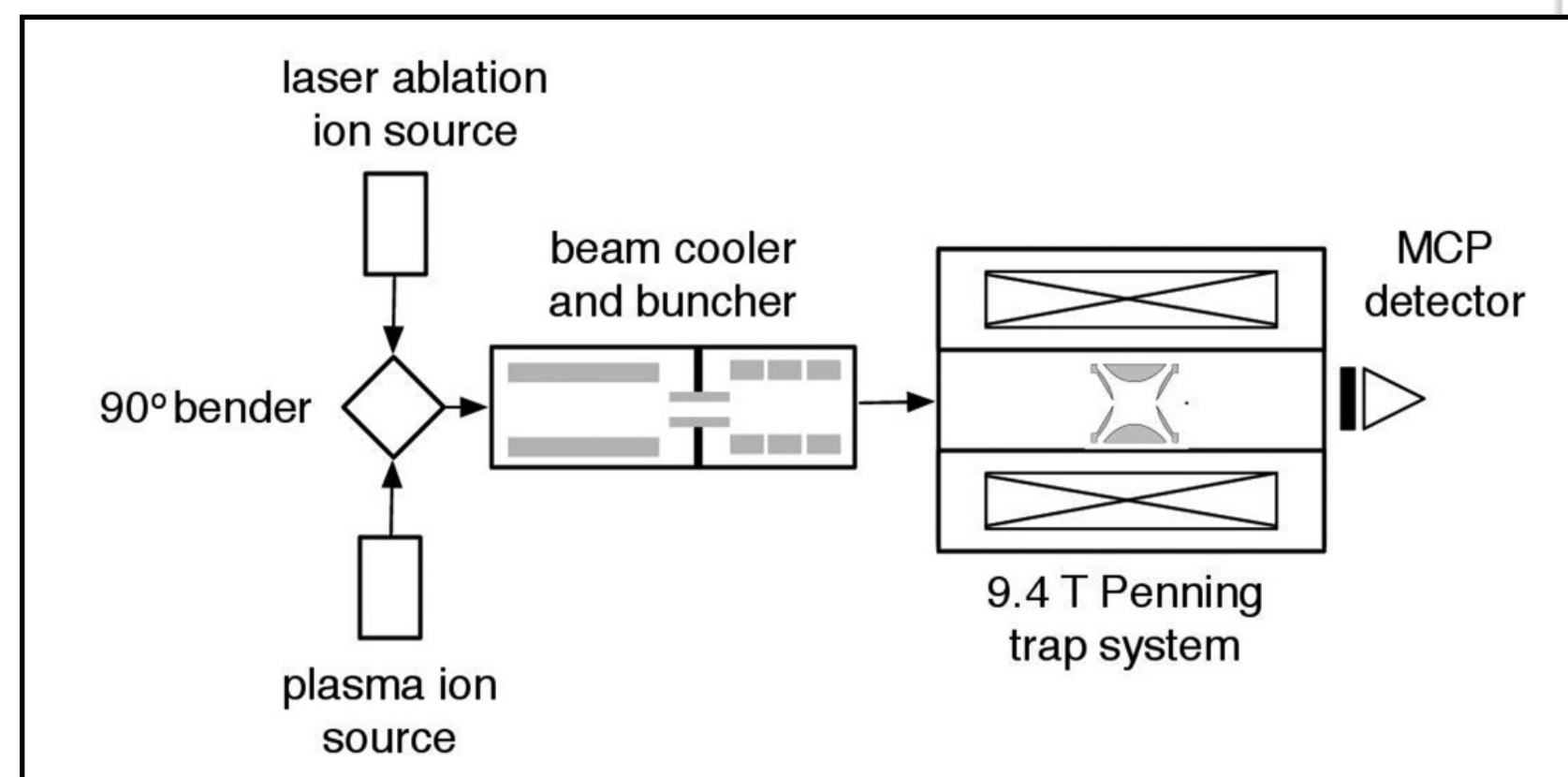
⁸Experimental Physics Department, CERN, Geneva, Switzerland

⁹RIKEN Nishina Center, Wako, Saitama 351-0198, Japan

¹⁰Institut für Physik, Universität Greifswald, 17487 Greifswald, Germany

¹¹Institut für Kern- und Teilchenphysik, Technische Universität Dresden, Germany

(Received 28 February 2019; revised manuscript received 21 May 2019; published 8 July 2019)



PHYSICAL REVIEW C **100**, 014308 (2019)

Direct determination of the ^{138}La β -decay Q value using Penning trap mass spectrometry

R. Sandler,^{1,2,*} G. Bollen,^{2,3,4} J. Dissanayake,¹ M. Eibach,^{2,5} K. Gulyuz,¹ A. Hamaker,^{2,4} C. Izzo,^{2,4} X. Mougeot,⁶ D. Puentes,^{2,4} F. G. A. Quarati,^{7,8} M. Redshaw,^{2,1} R. Ringle,² and I. Yandow^{2,4}

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³Facility for Rare Isotope Beams, East Lansing, Michigan 48824, USA

⁴Department of Physics and Astronomy, Michigan State University, East Lansing, Michigan 48824, USA

⁵Institut für Physik, Universität Greifswald, 17487 Greifswald, Germany

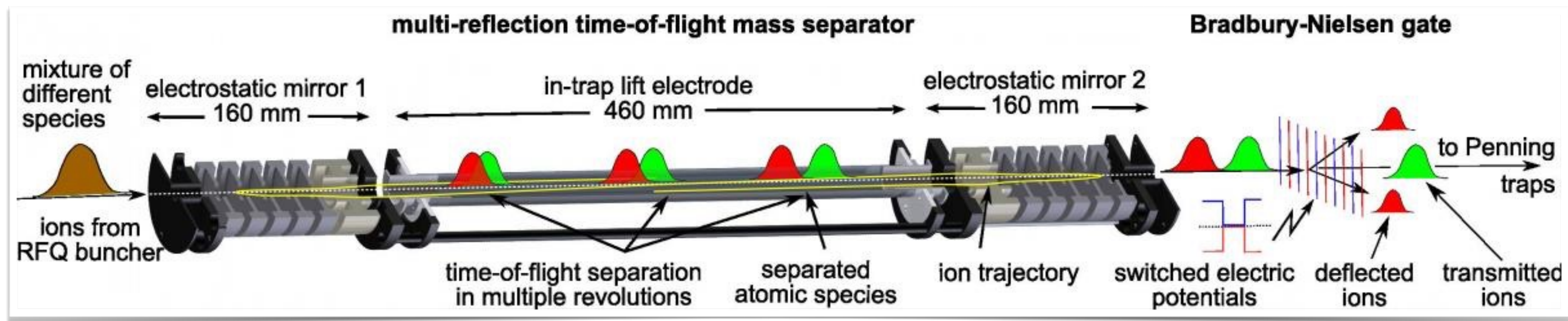
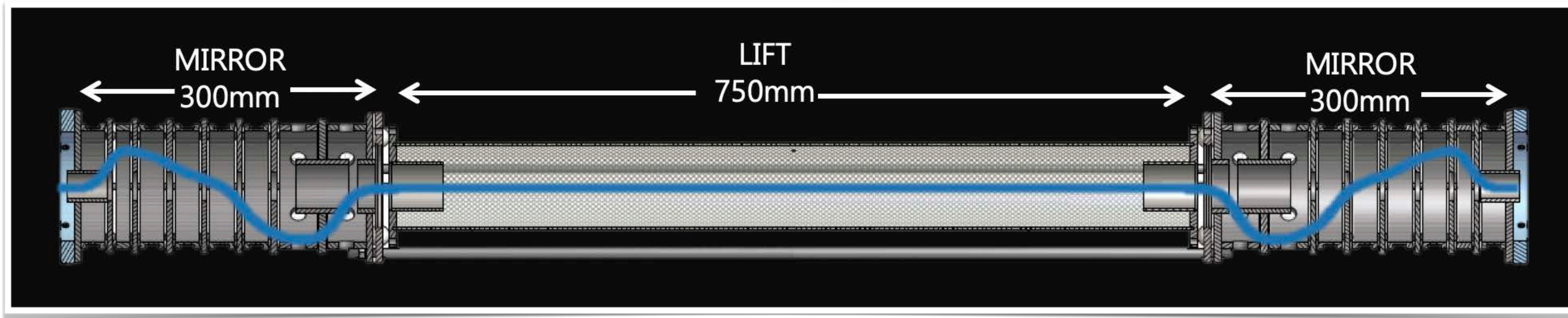
⁶CEA, LIST, Laboratoire National Henri Becquerel (LNE-LNHB), Bât. 602 PC111, CEA-Saclay 91191 Gif-sur-Yvette Cedex, France

⁷AS, RST, LM, Delft University of Technology, Mekelweg 15, 2629JB Delft, The Netherlands

⁸Gonitec BV, Johannes Bildersstraat 60, 259EJ Den Haag, The Netherlands

(Received 29 April 2019; published 11 July 2019)

MR-TOF

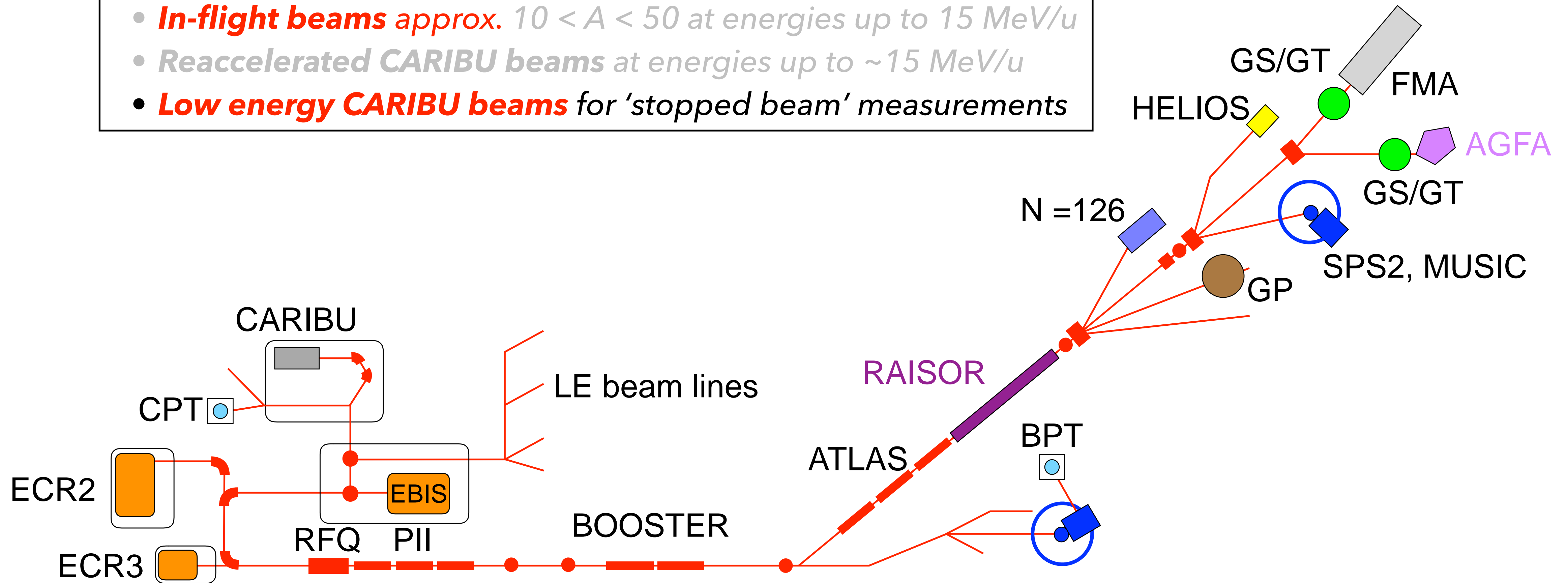


Ions cycle back and forth. A time separation occurs such that $t \sim \sqrt{m/q}$. Fast. High resolution.

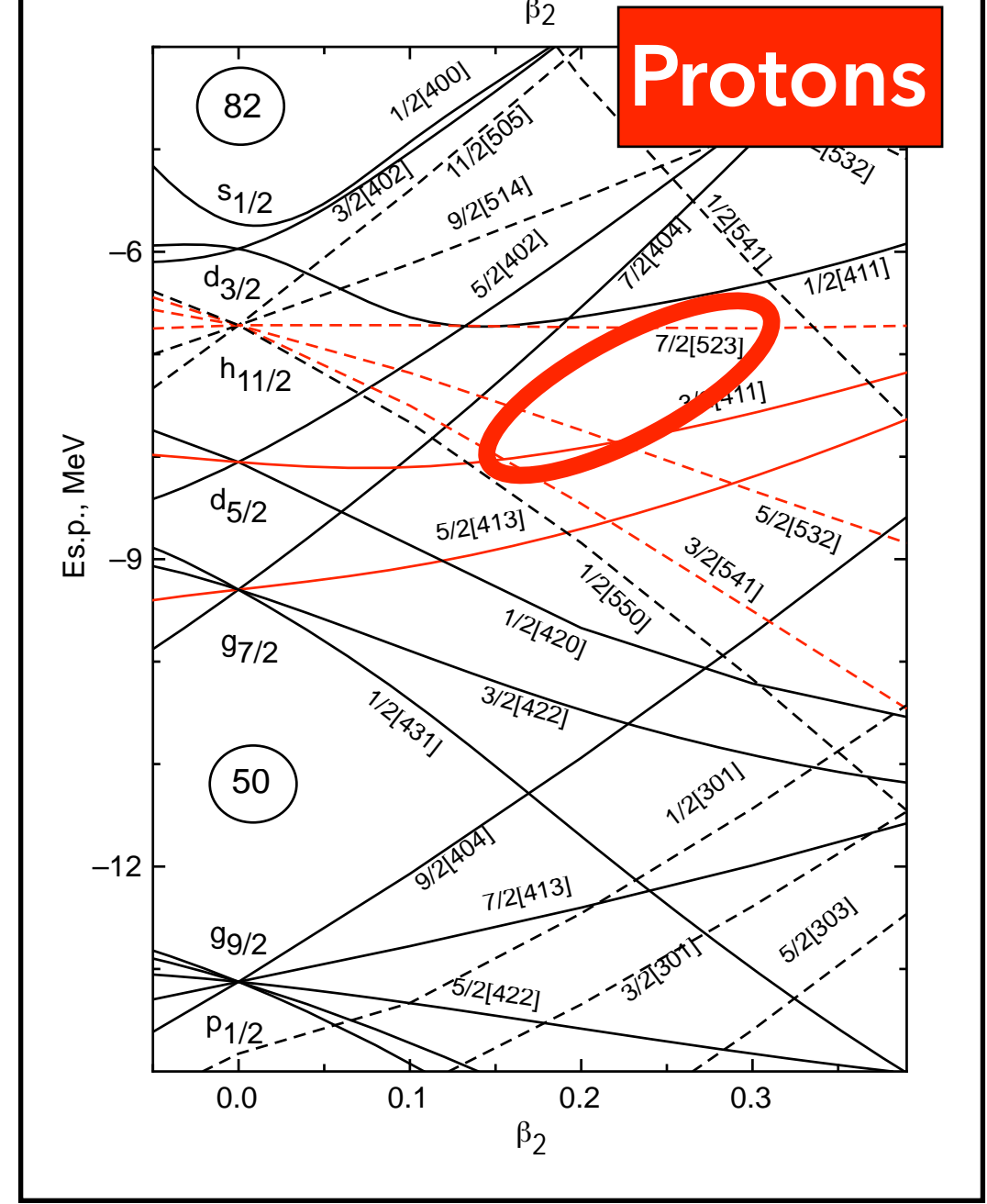
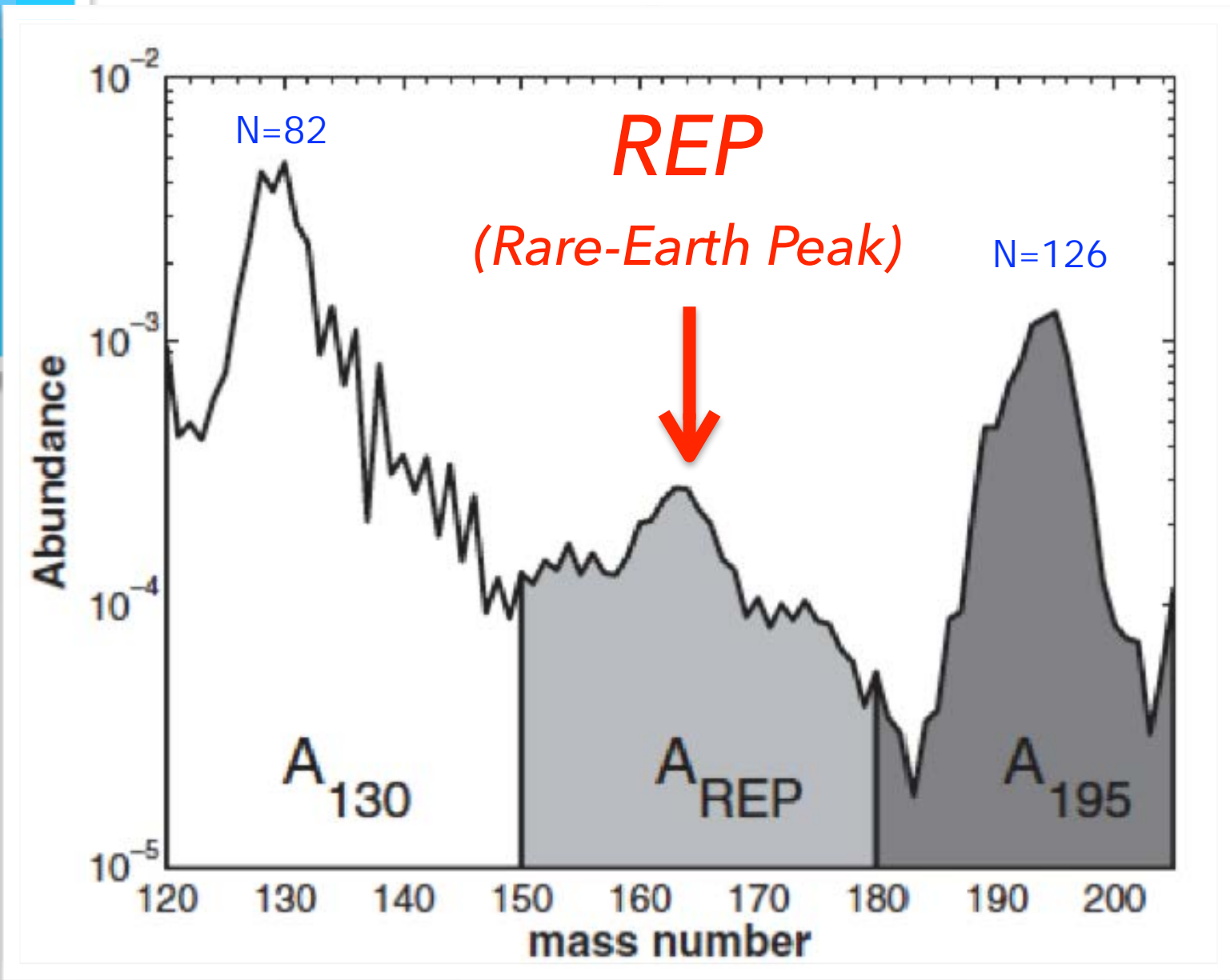
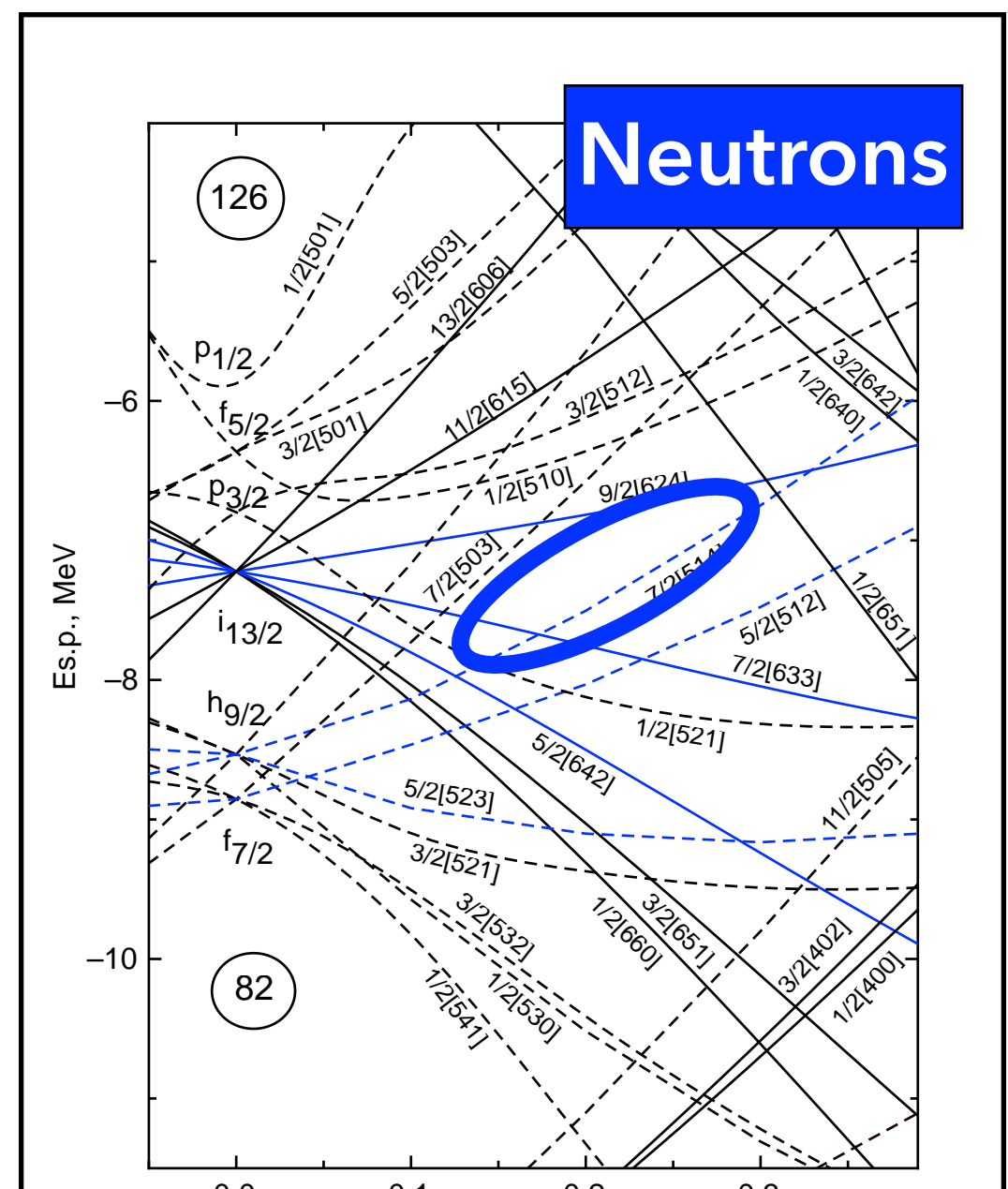
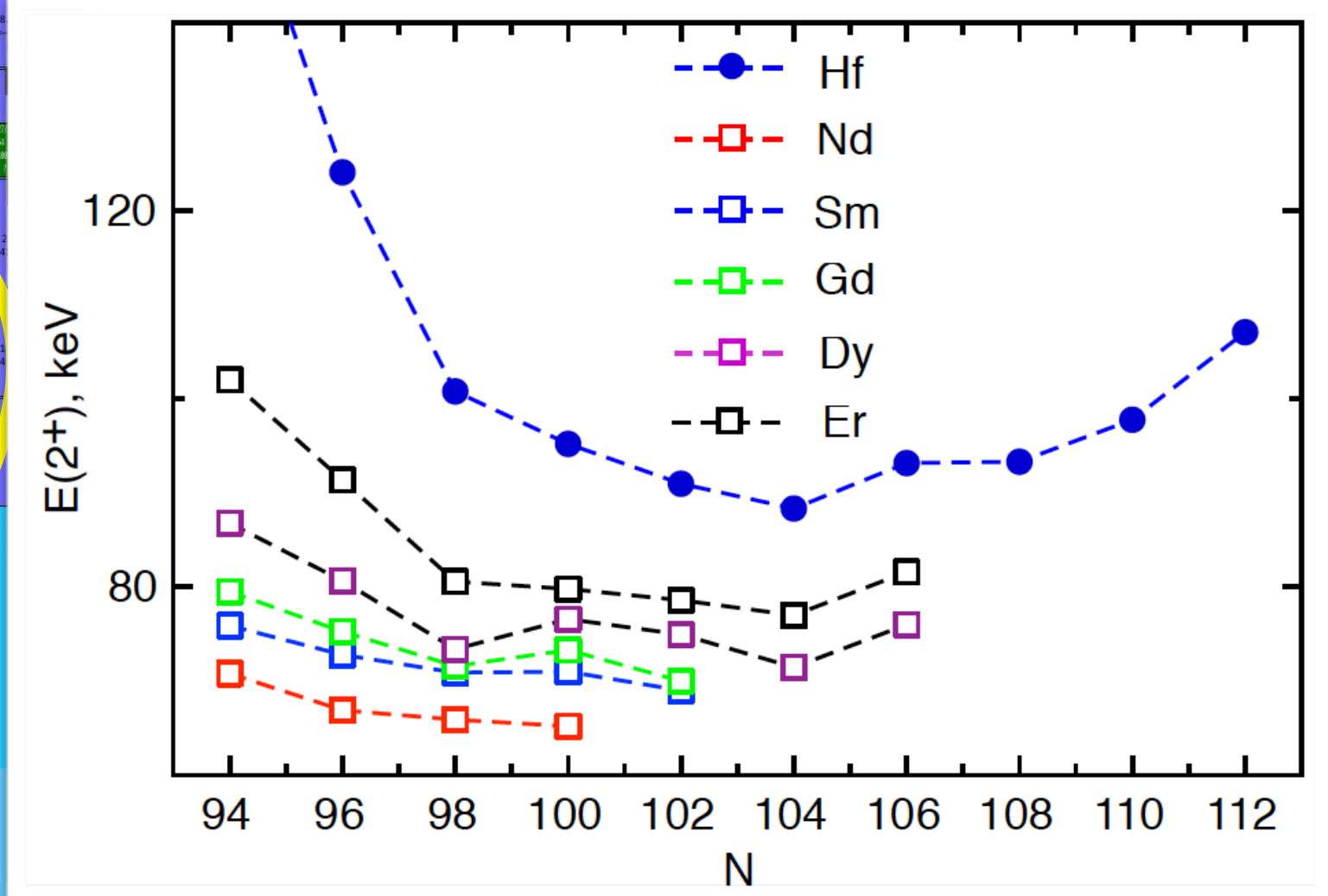
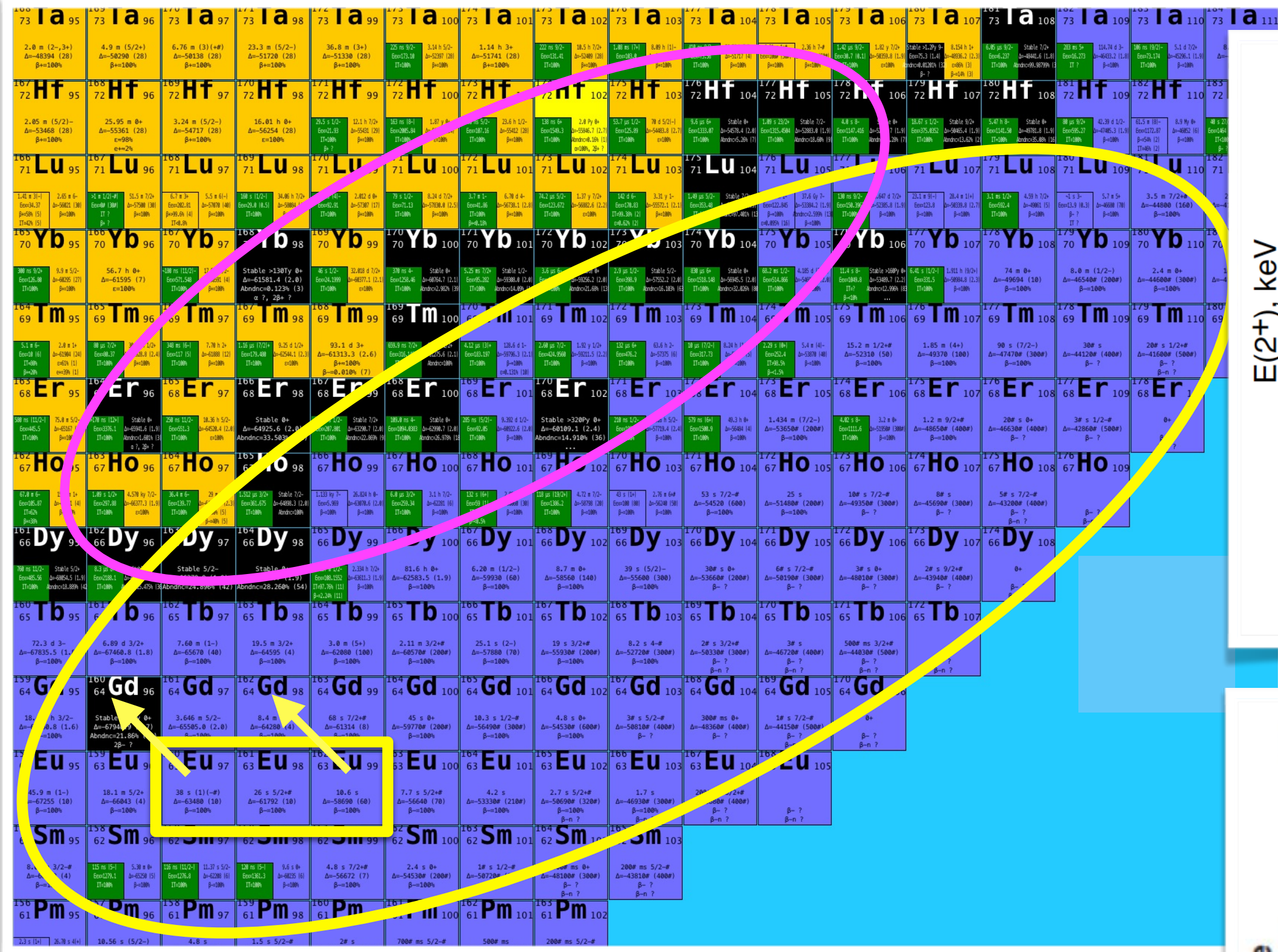
Images and suggestions from Rodney Orford (LBNL) and Jason Clark (ANL)

ATLAS

- *Stable beams at high intensity and energies up to 18 MeV/u*
- **In-flight beams approx.** $10 < A < 50$ at energies up to 15 MeV/u
- *Reaccelerated CARIBU beams at energies up to ~15 MeV/u*
- **Low energy CARIBU beams** for 'stopped beam' measurements

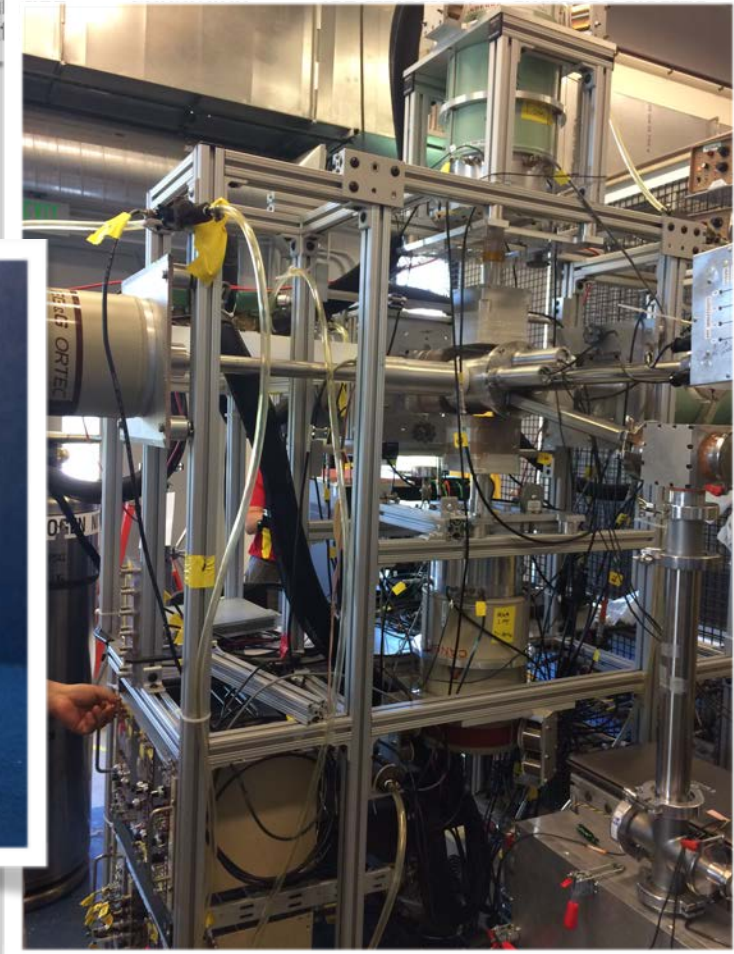
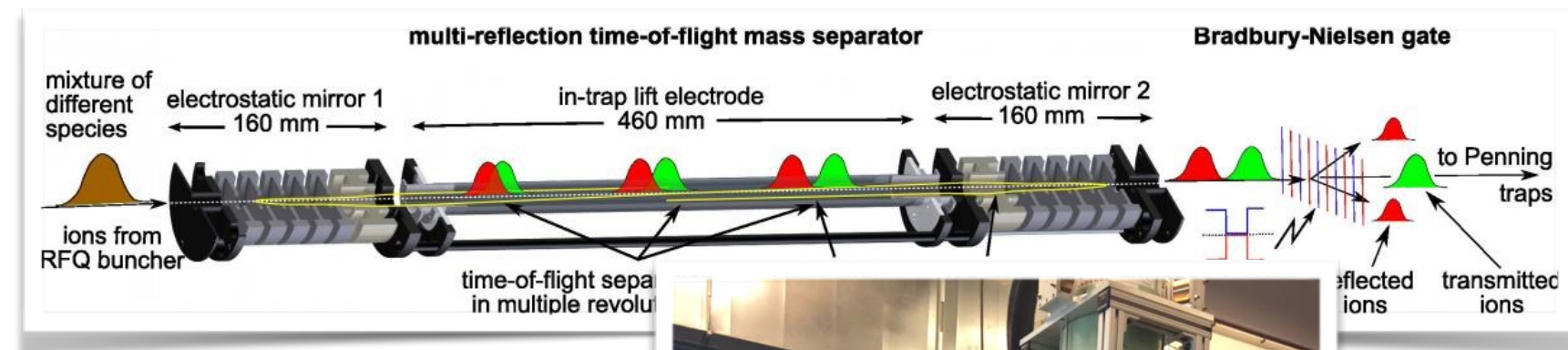


Deformed, neutron-rich nuclei

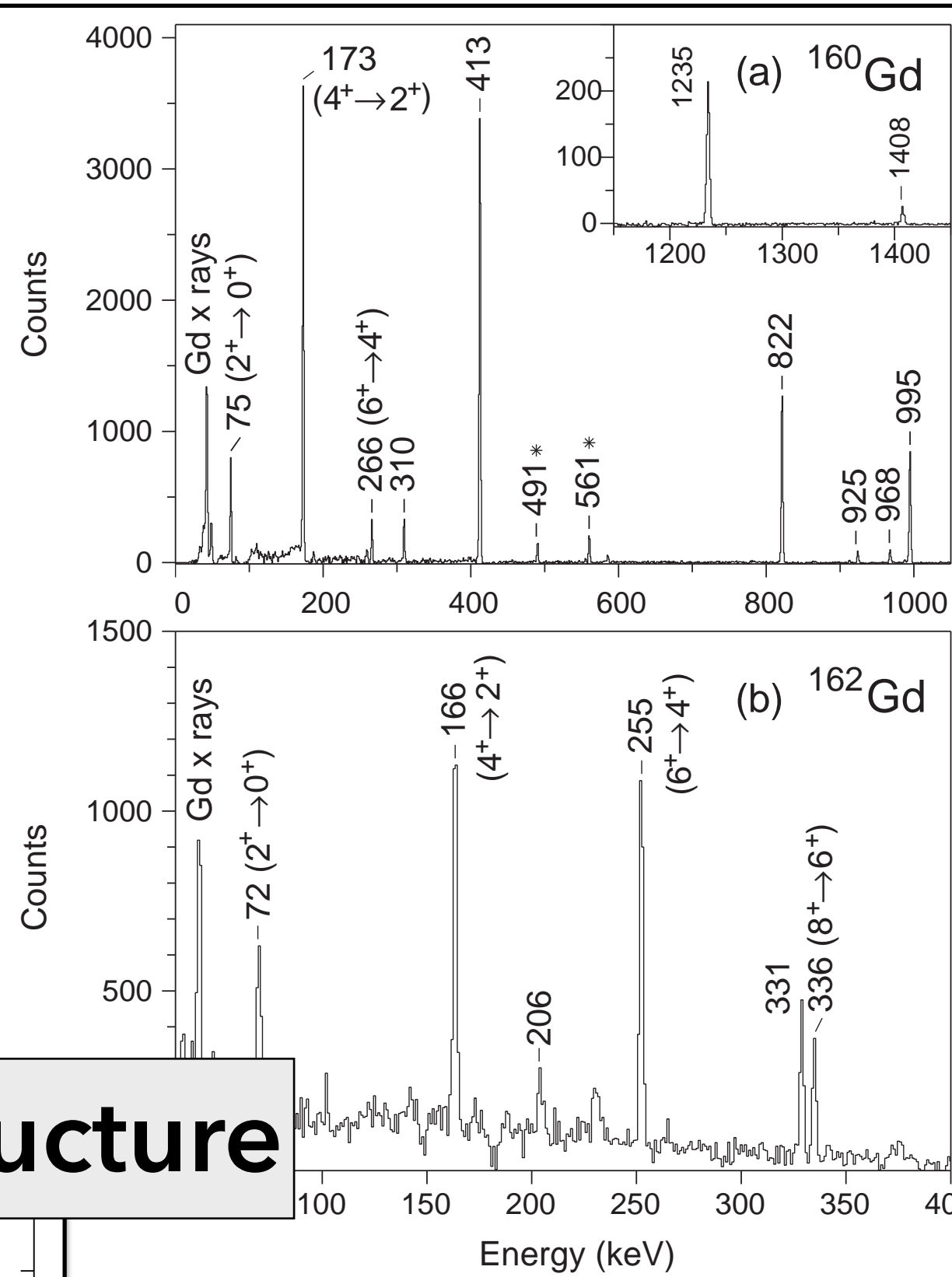
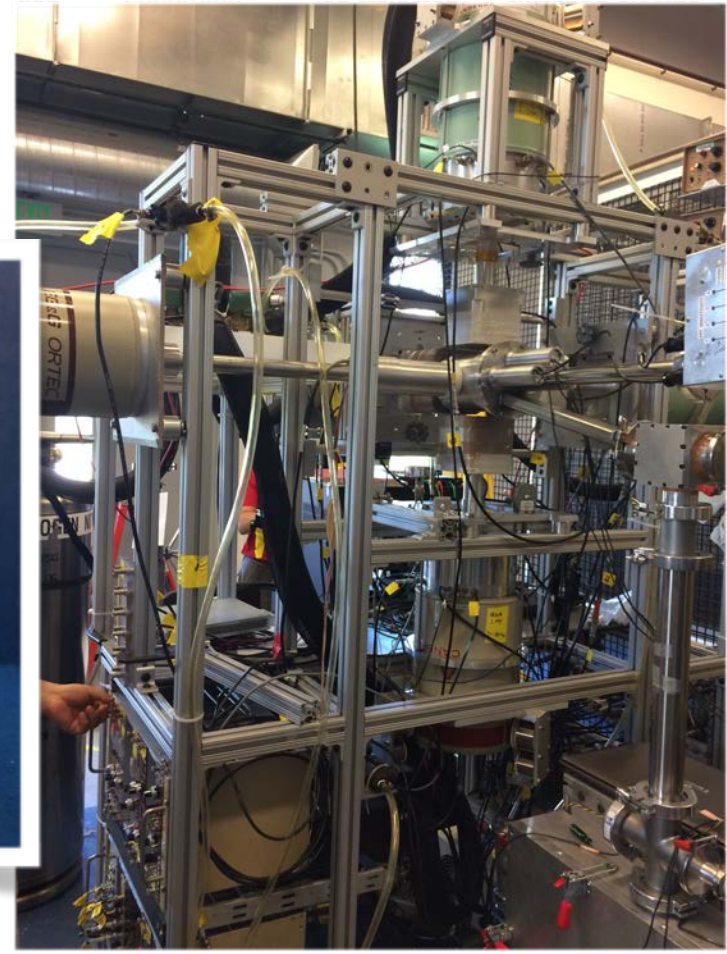
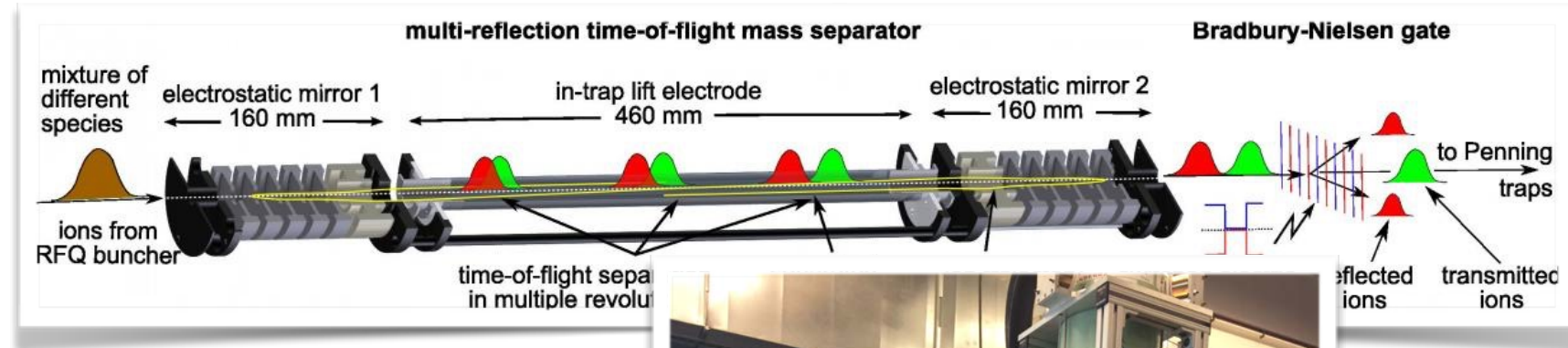


CARIBU allows access to terra incognita: What stabilizes these deformed shapes? What role does the structure of these nuclei have on the r -process abundance?

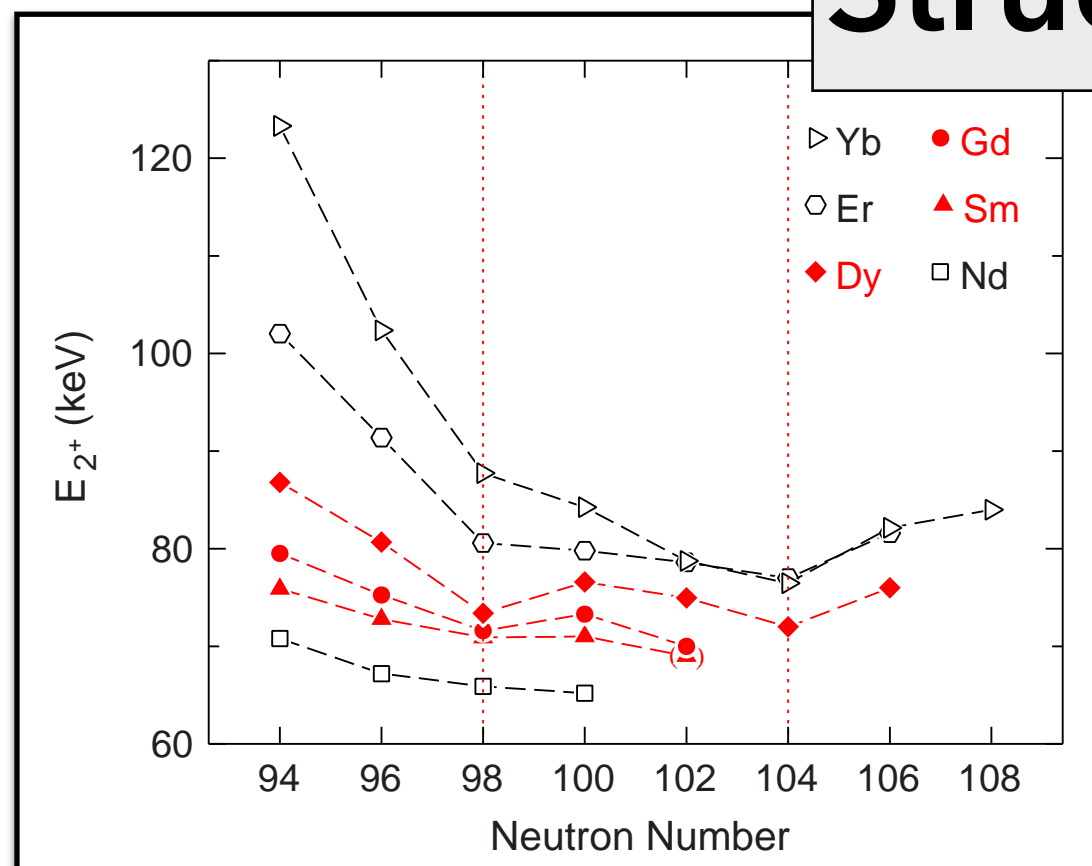
Detailed spectroscopy



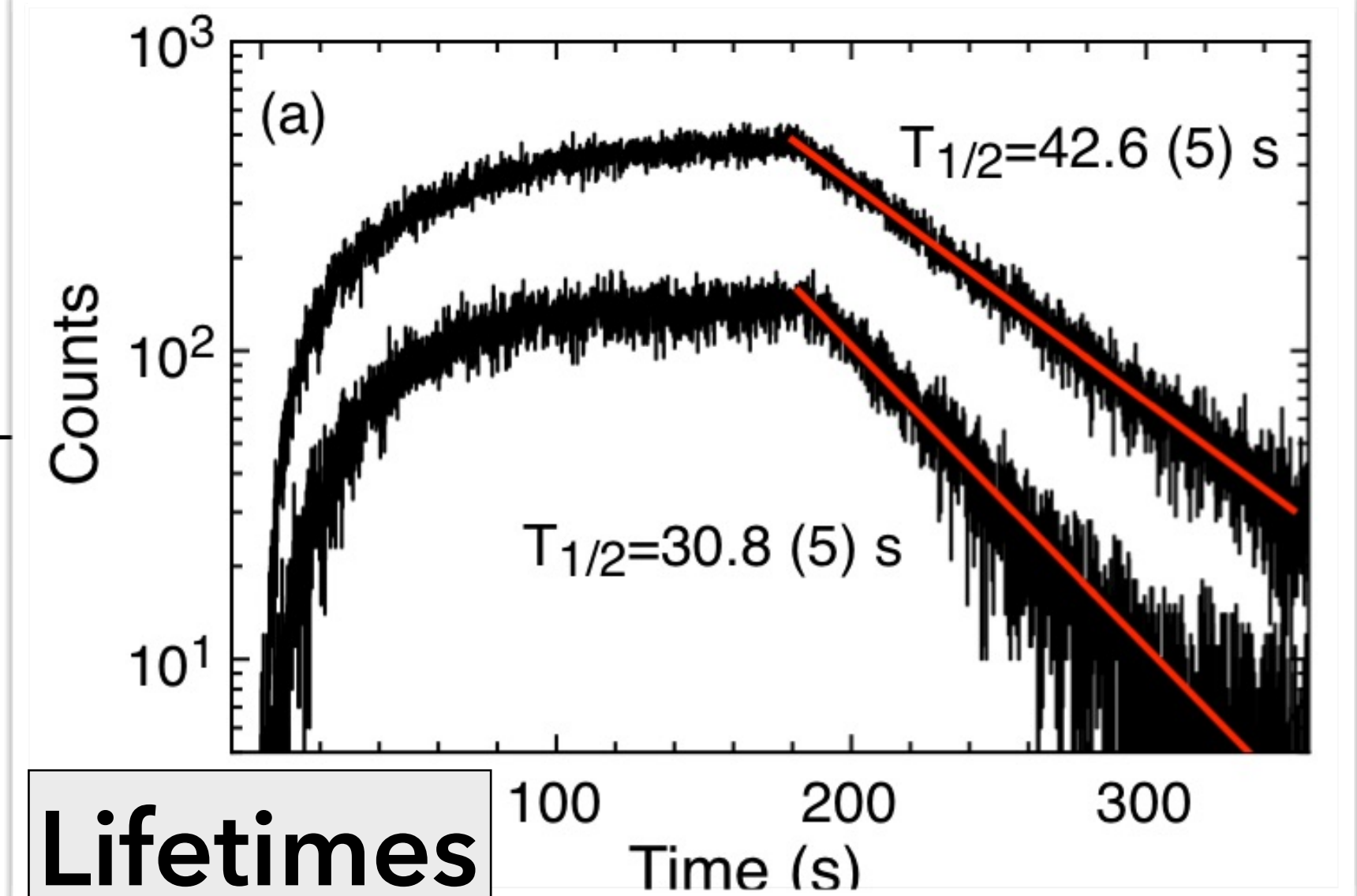
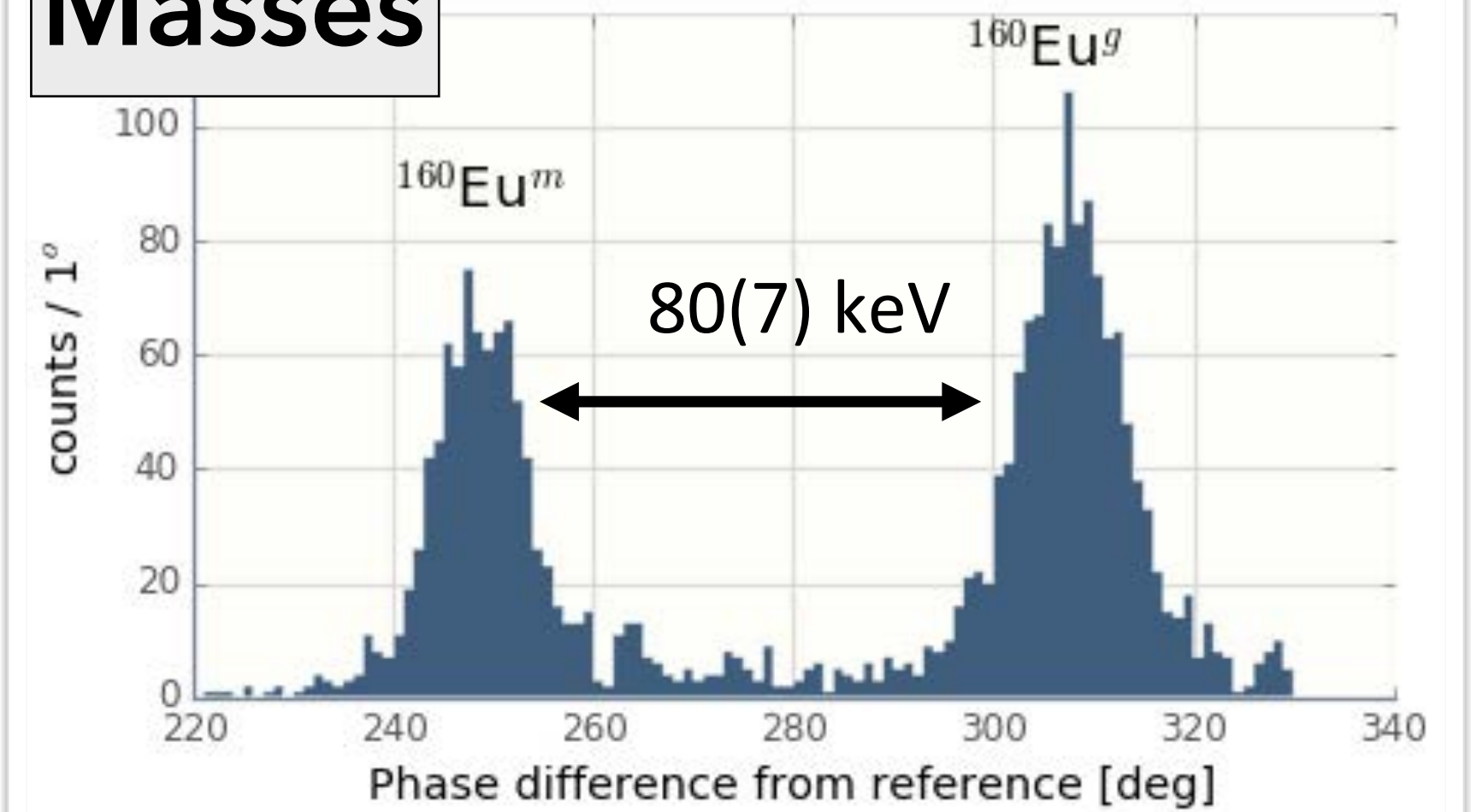
Detailed spectroscopy



Structure

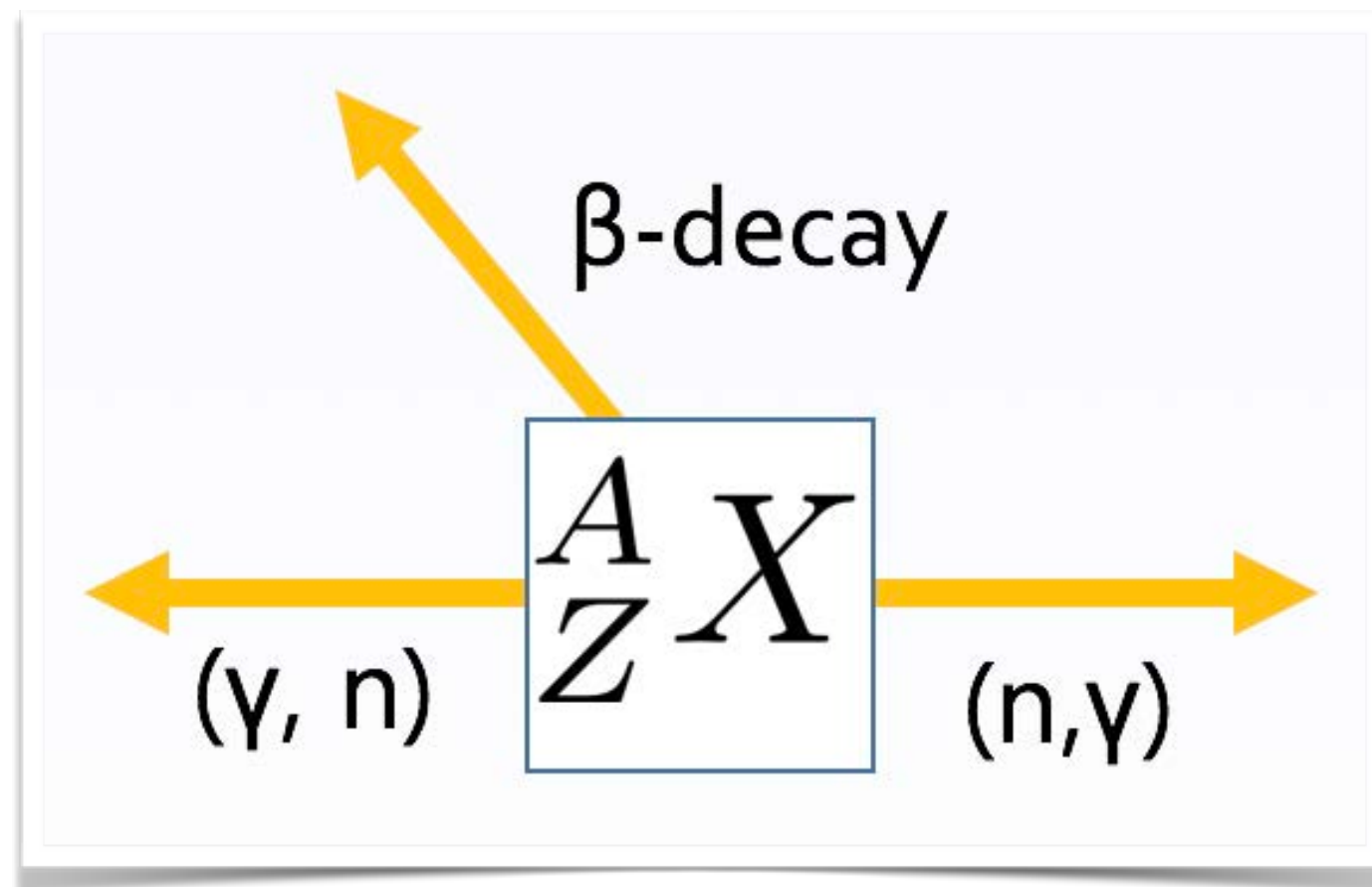


Masses



Lifetimes

Understanding REP formation

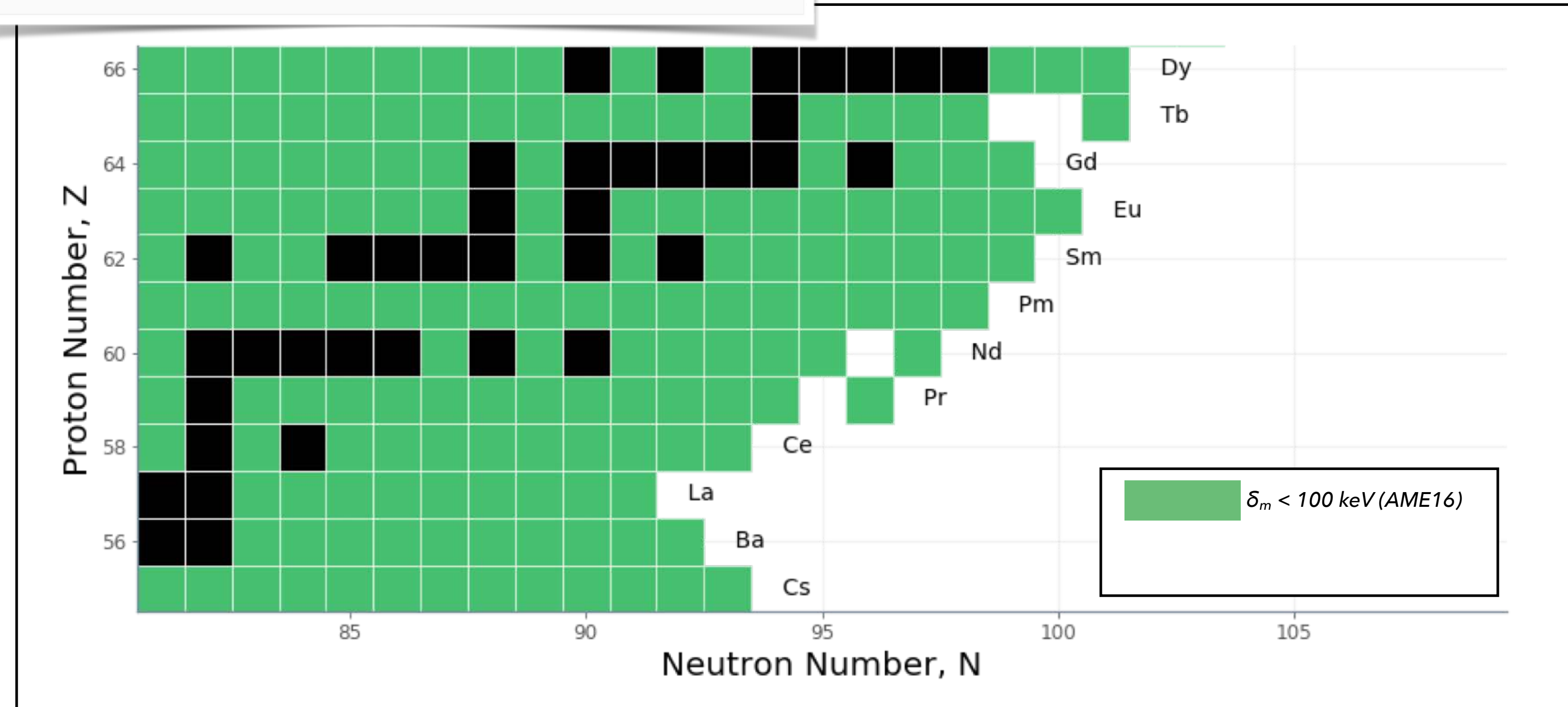
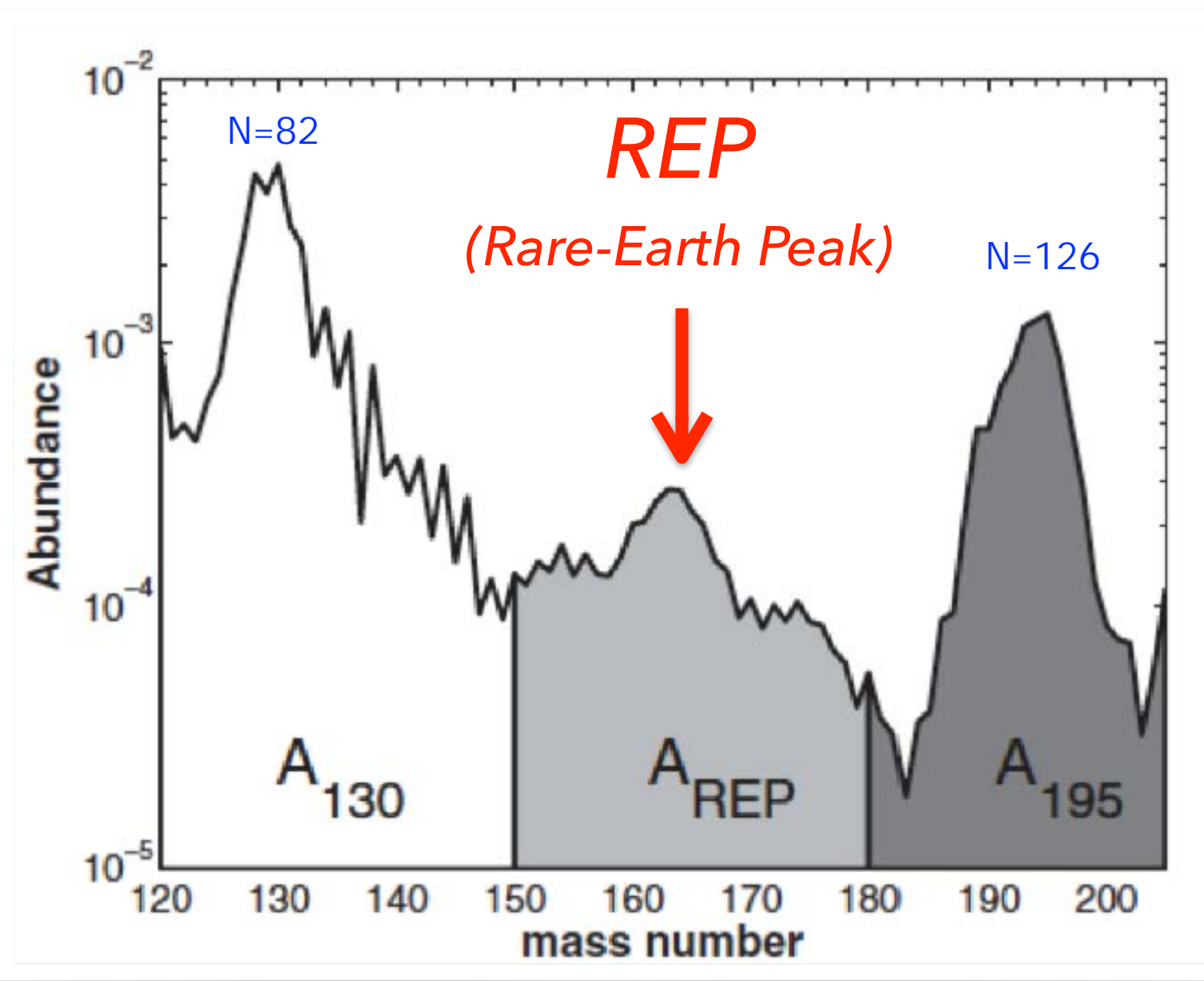


General classification:

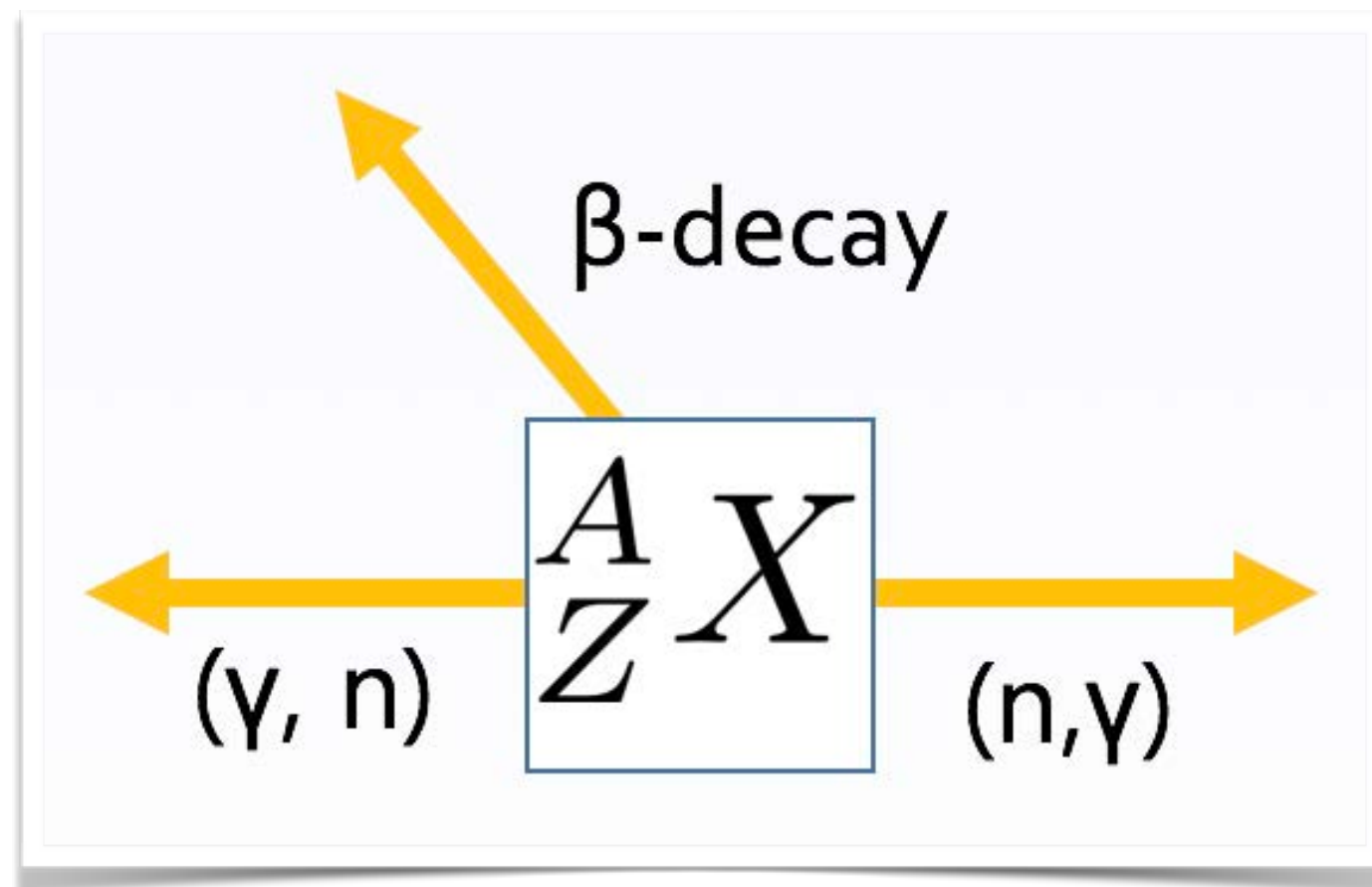
Hot: $(n, \gamma) \rightleftharpoons (\gamma, n)$

Cold: no equilibrium, competition between decay and neutron-capture

Masses measurements necessary to gain insights into what **environment produces the observed abundance peaks**



Understanding REP formation

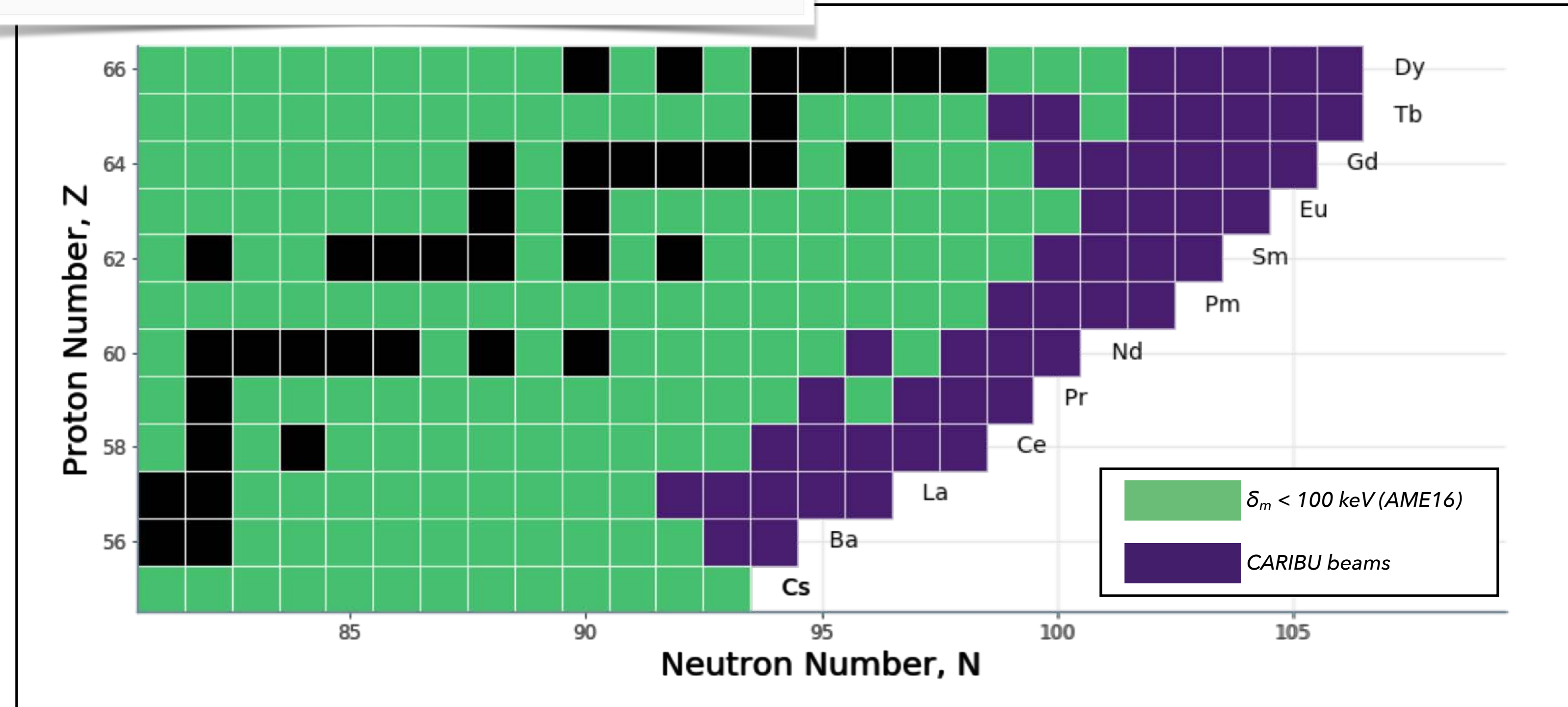
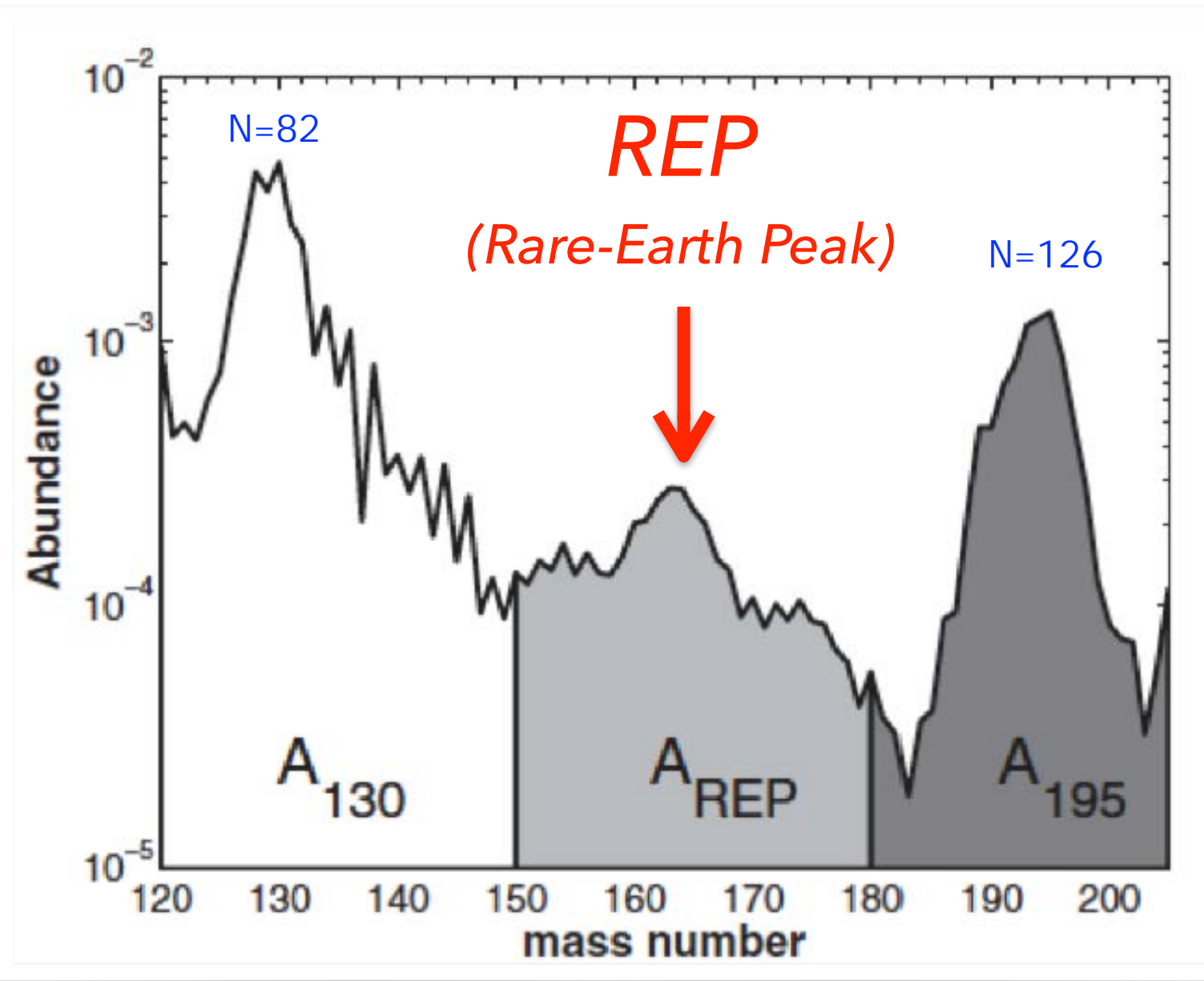


General classification:

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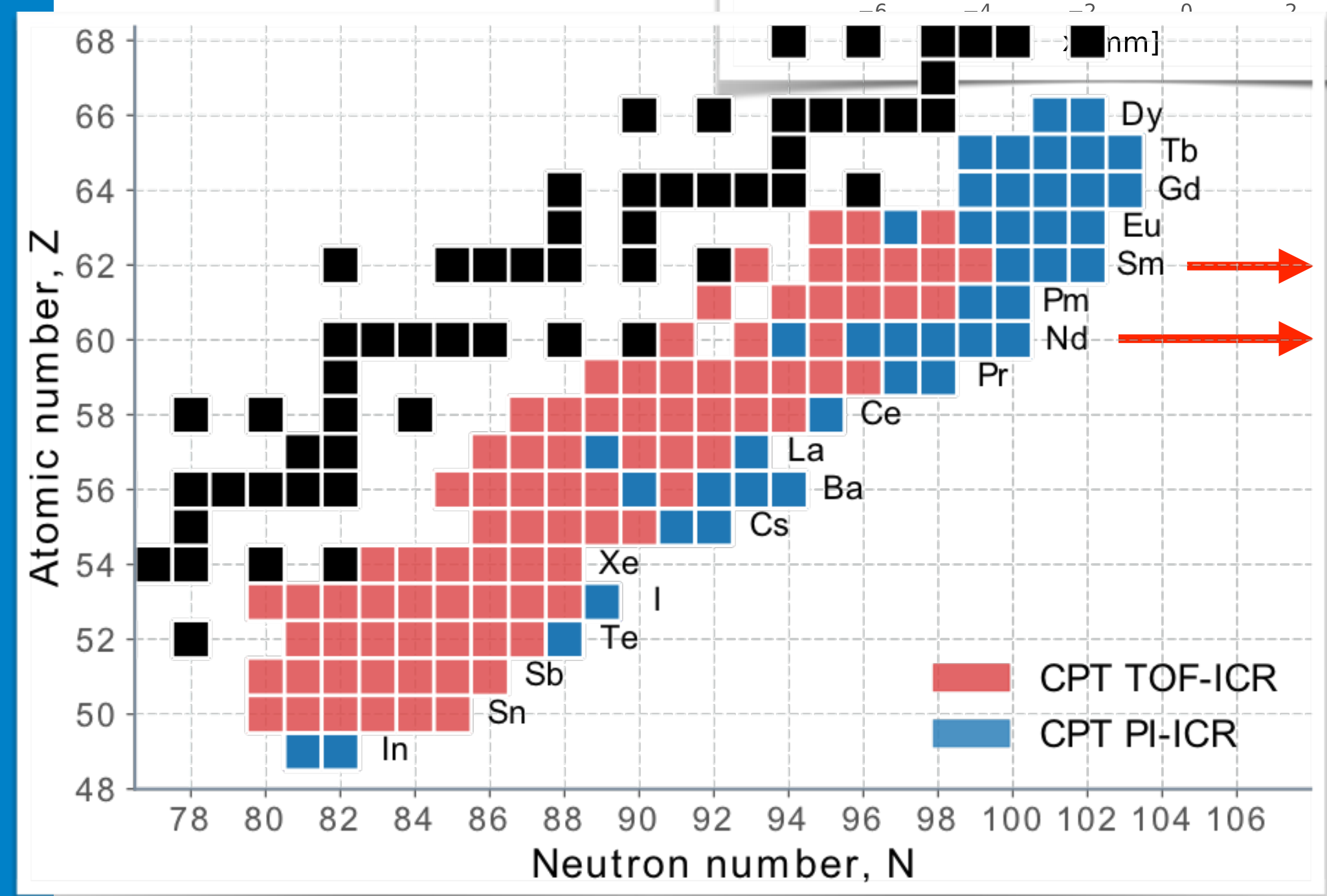
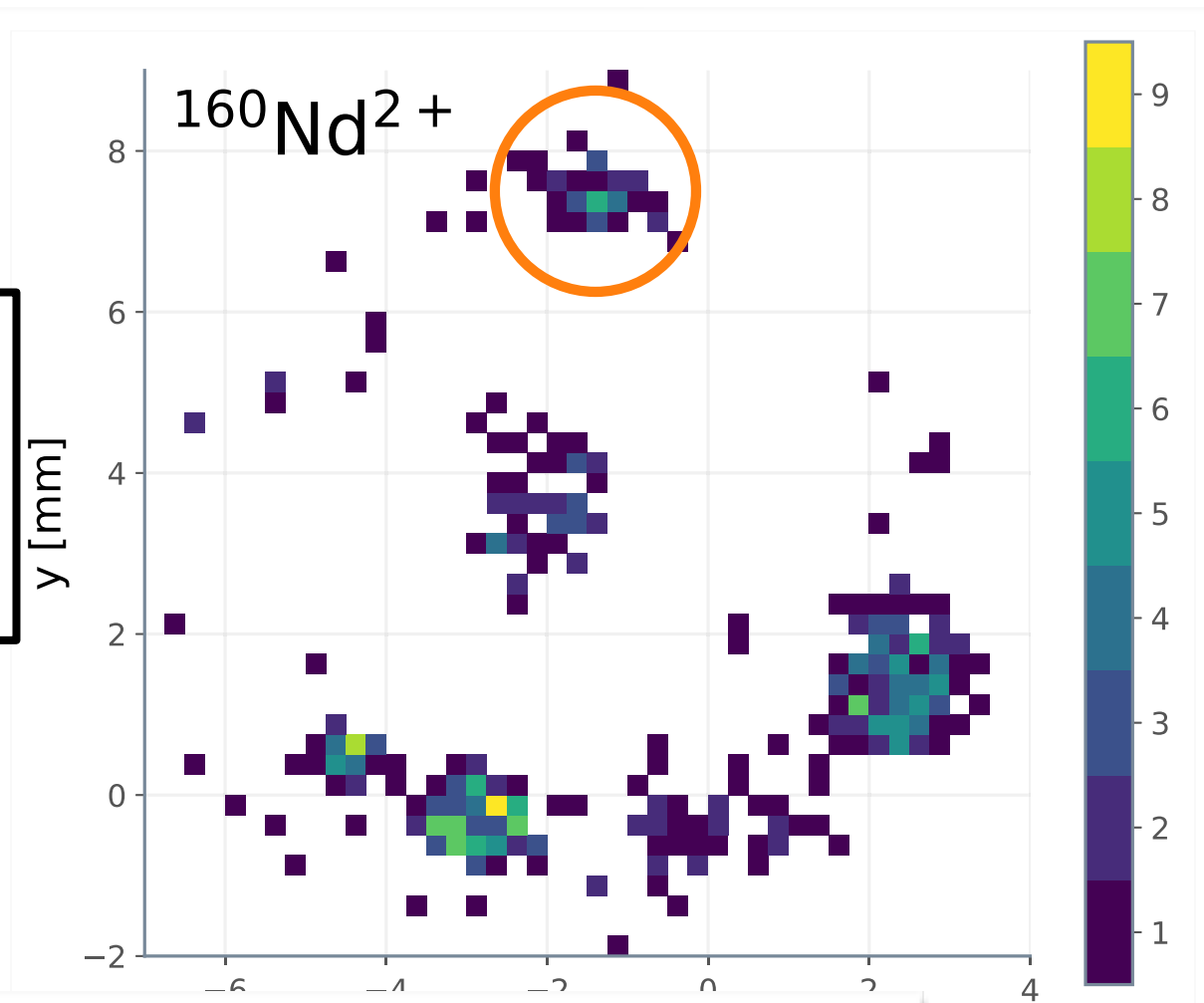
Cold: no equilibrium, competition between decay and neutron-capture

Masses measurements necessary to gain insights into what **environment produces the observed abundance peaks**

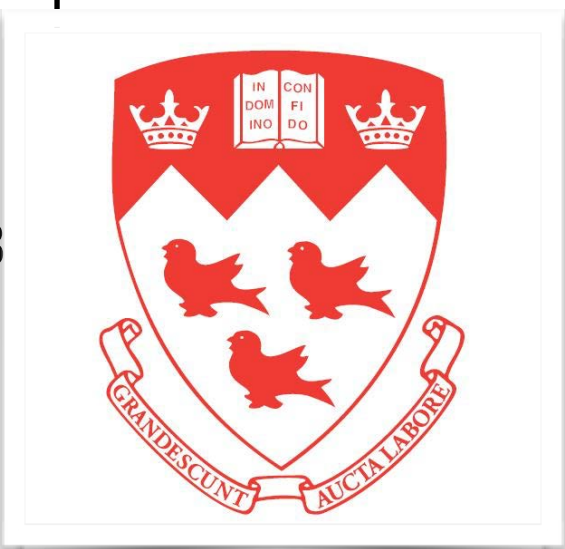
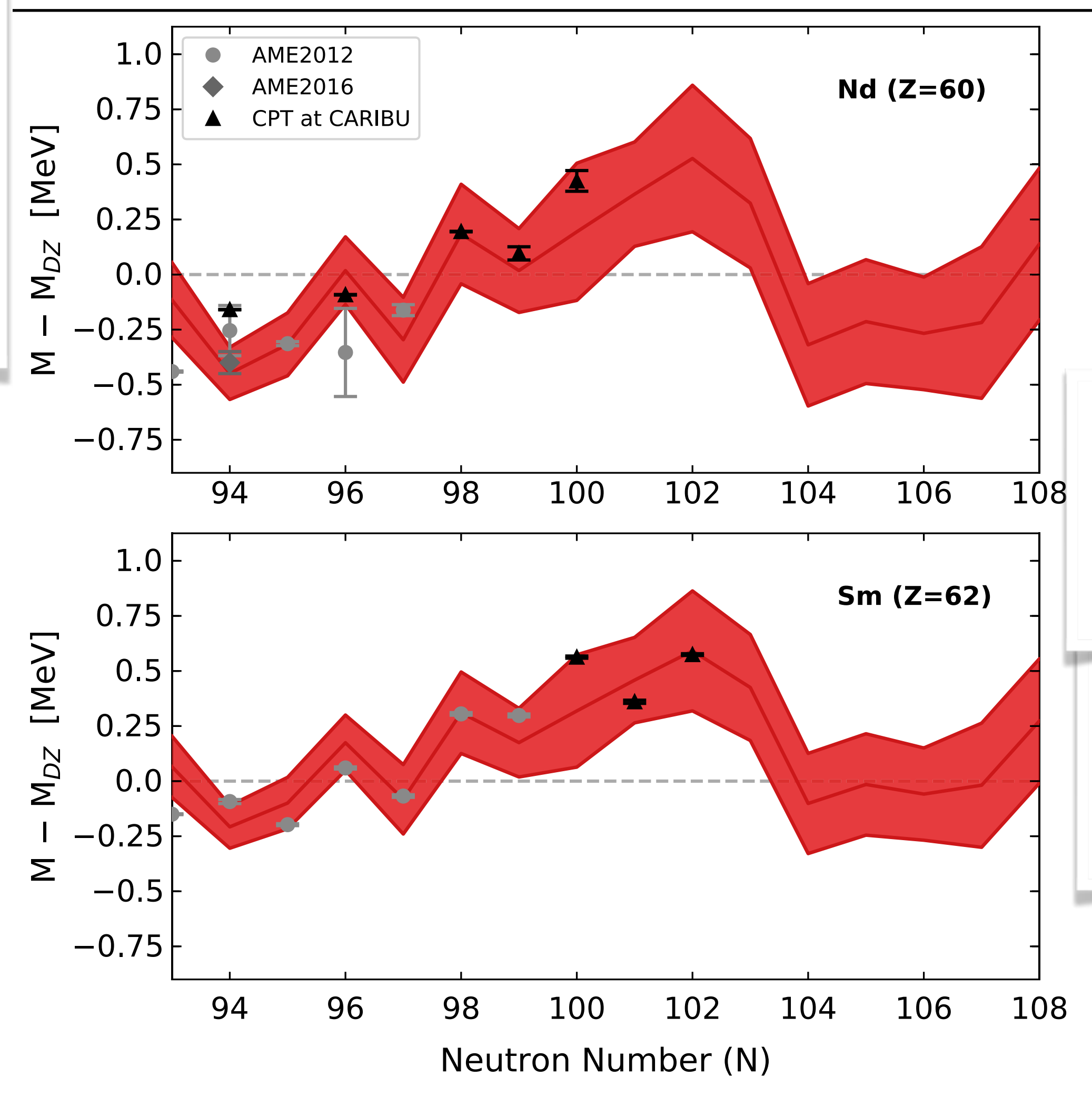


Masses, mass models

TOF: Time Of Flight
 PI: Phase Imaging
 ICR: Ion-Cyclotron Resonance



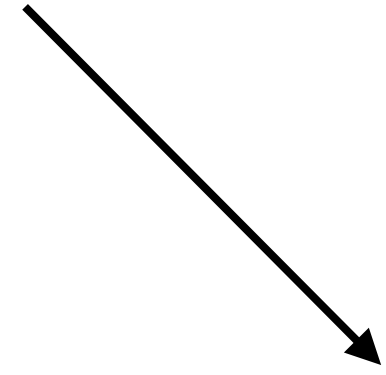
Monte Carlo Mass Corrections (MCMC): "reverse engineering" the masses



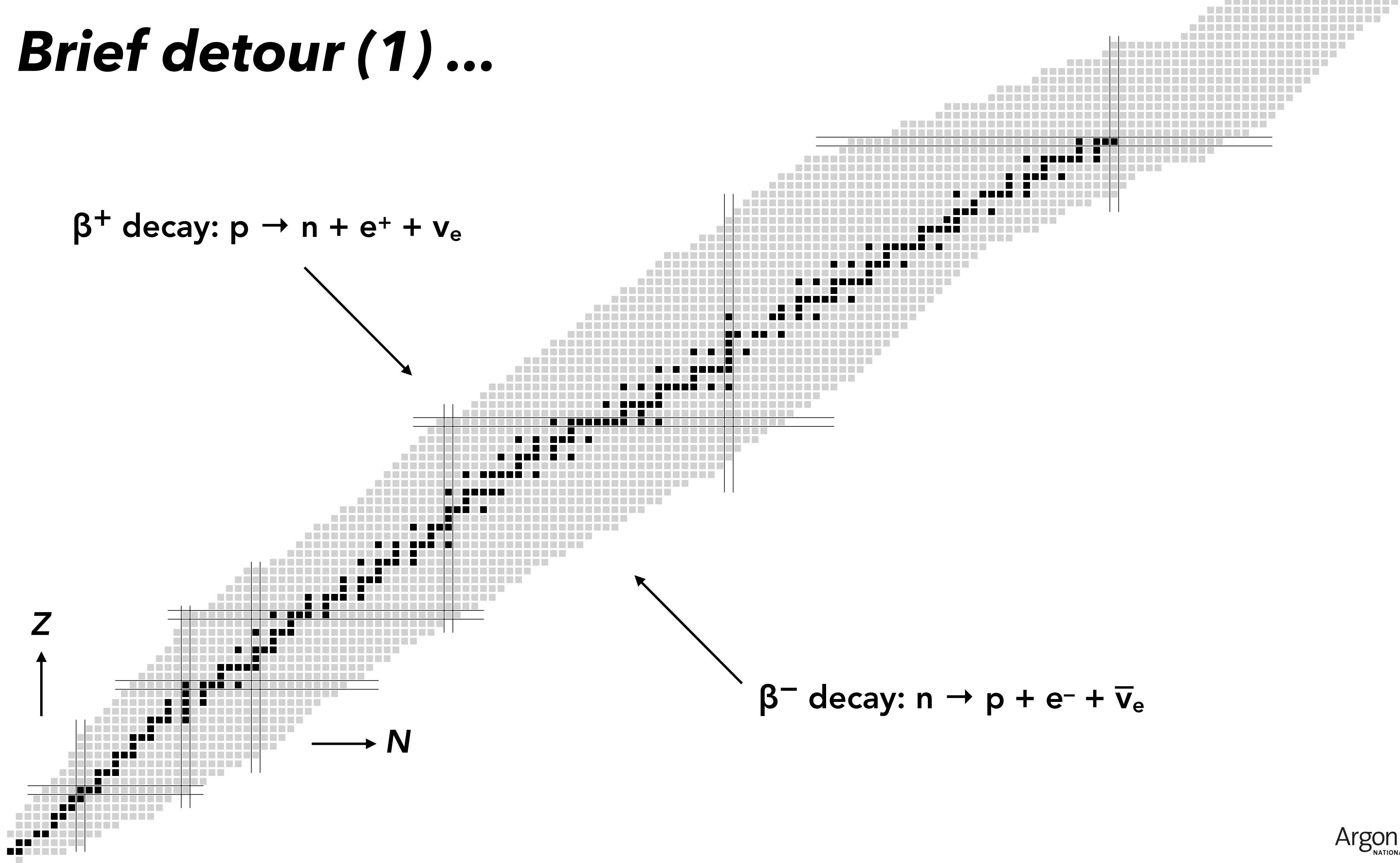
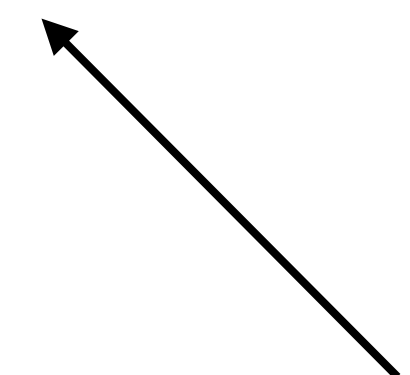
R. Orford et al., Phys. Rev. Lett. **120**, 262702 (2018)

Brief detour (1) ...

β^+ decay: $p \rightarrow n + e^+ + \nu_e$

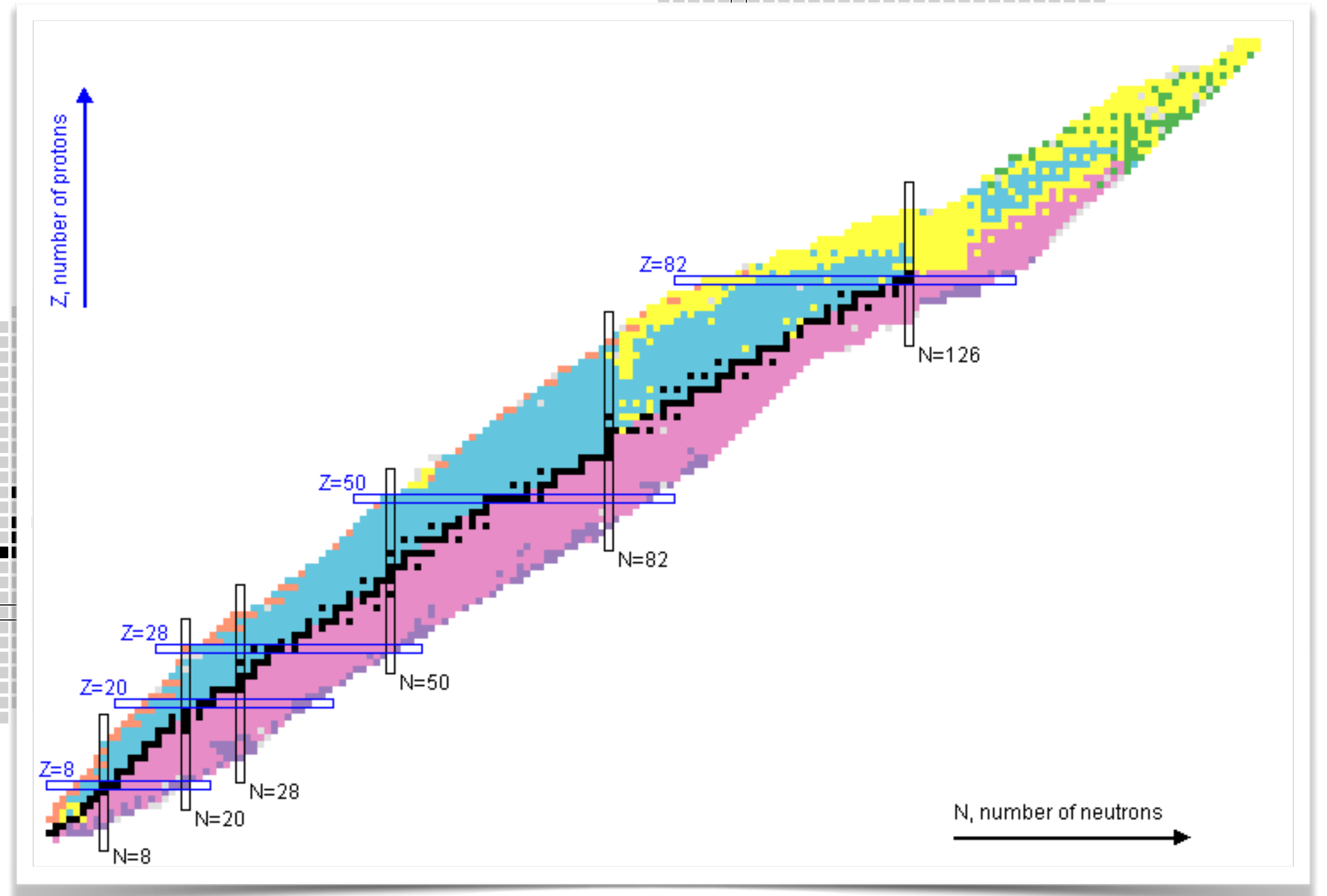
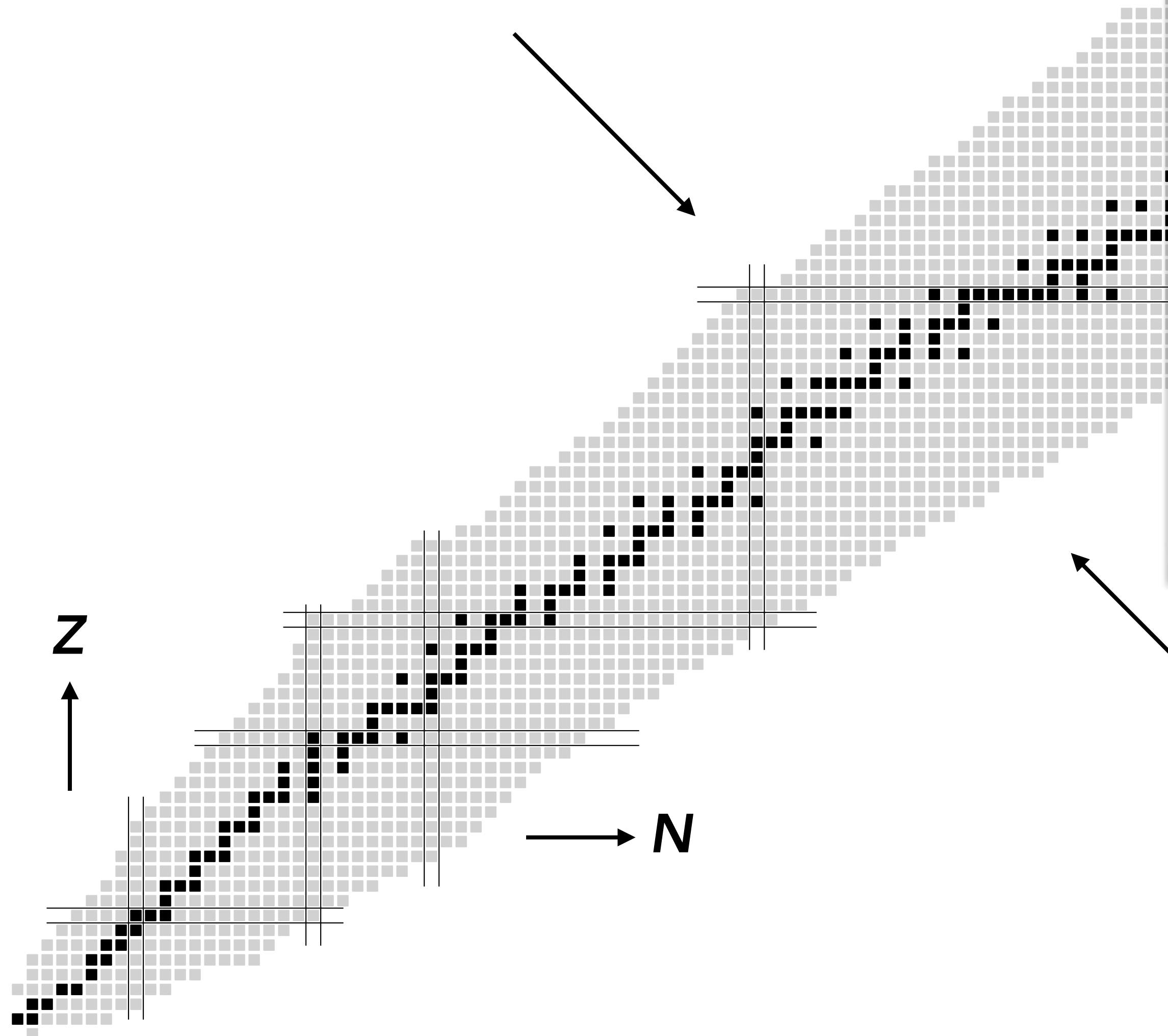
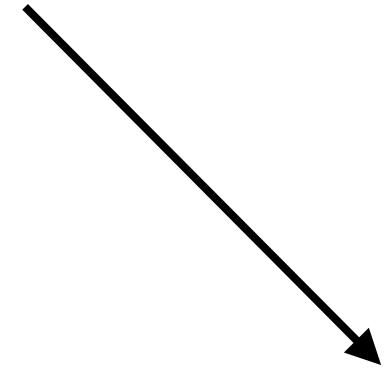


β^- decay: $n \rightarrow p + e^- + \bar{\nu}_e$



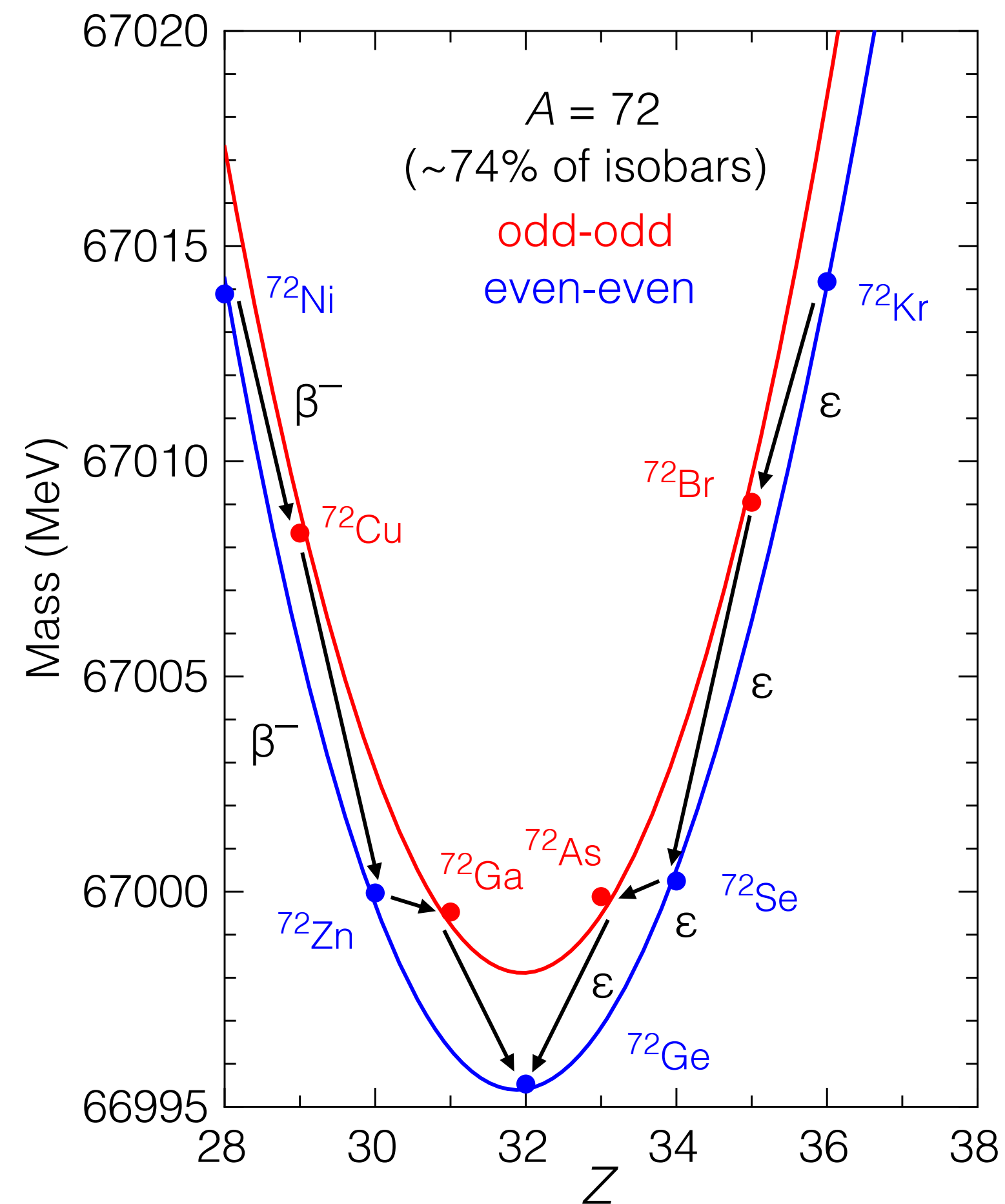
Brief detour (1) ...

β^+ decay: $p \rightarrow n + e^+ + \nu_e$



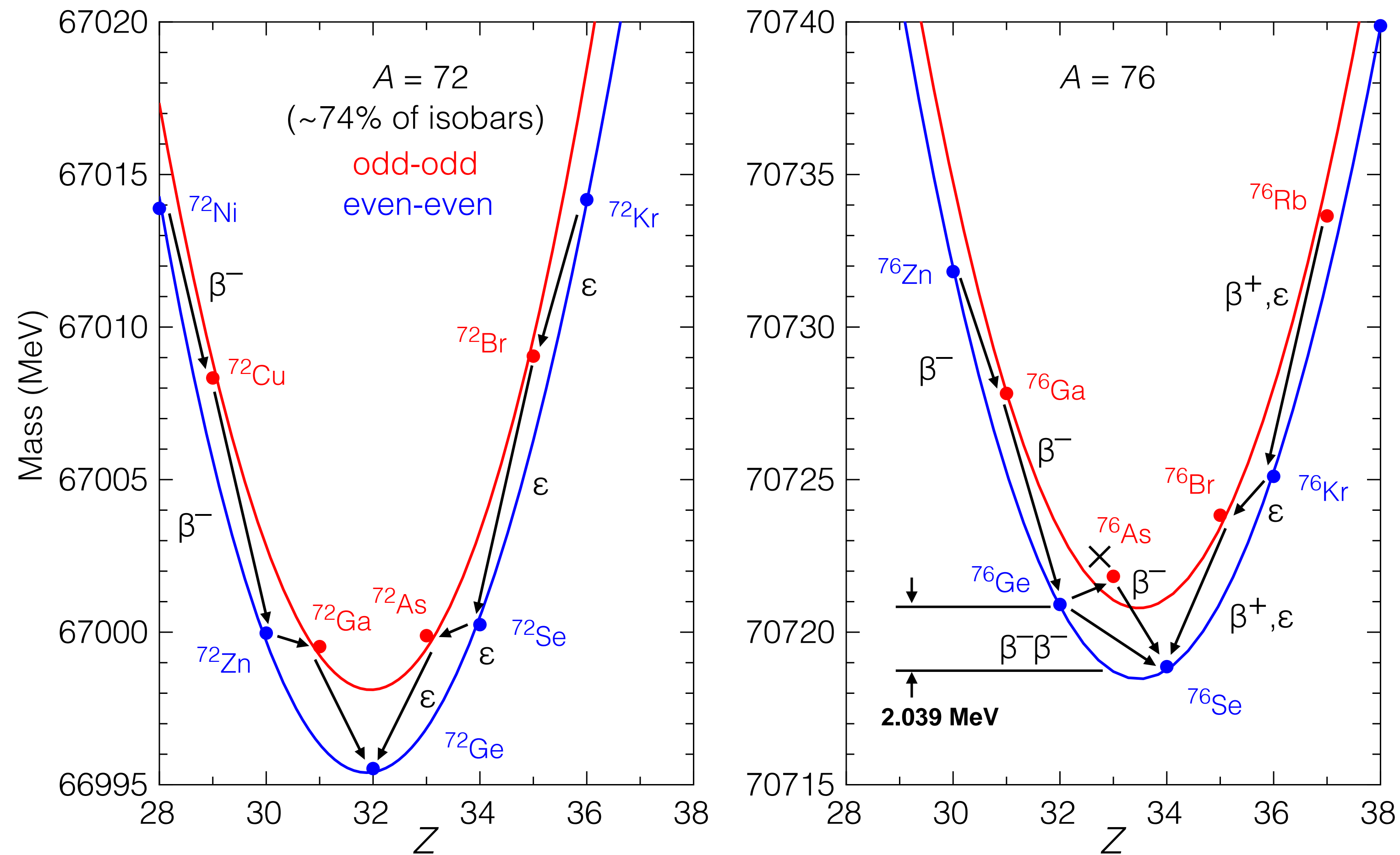
β^- decay: $n \rightarrow p + e^- + \bar{\nu}_e$

Beta decay, double beta decay ...



Pairing in nuclei results in a displacement of even-even and odd-odd mass parabolas for given isobars. Data from AME 2012.

Beta decay, double beta decay ...

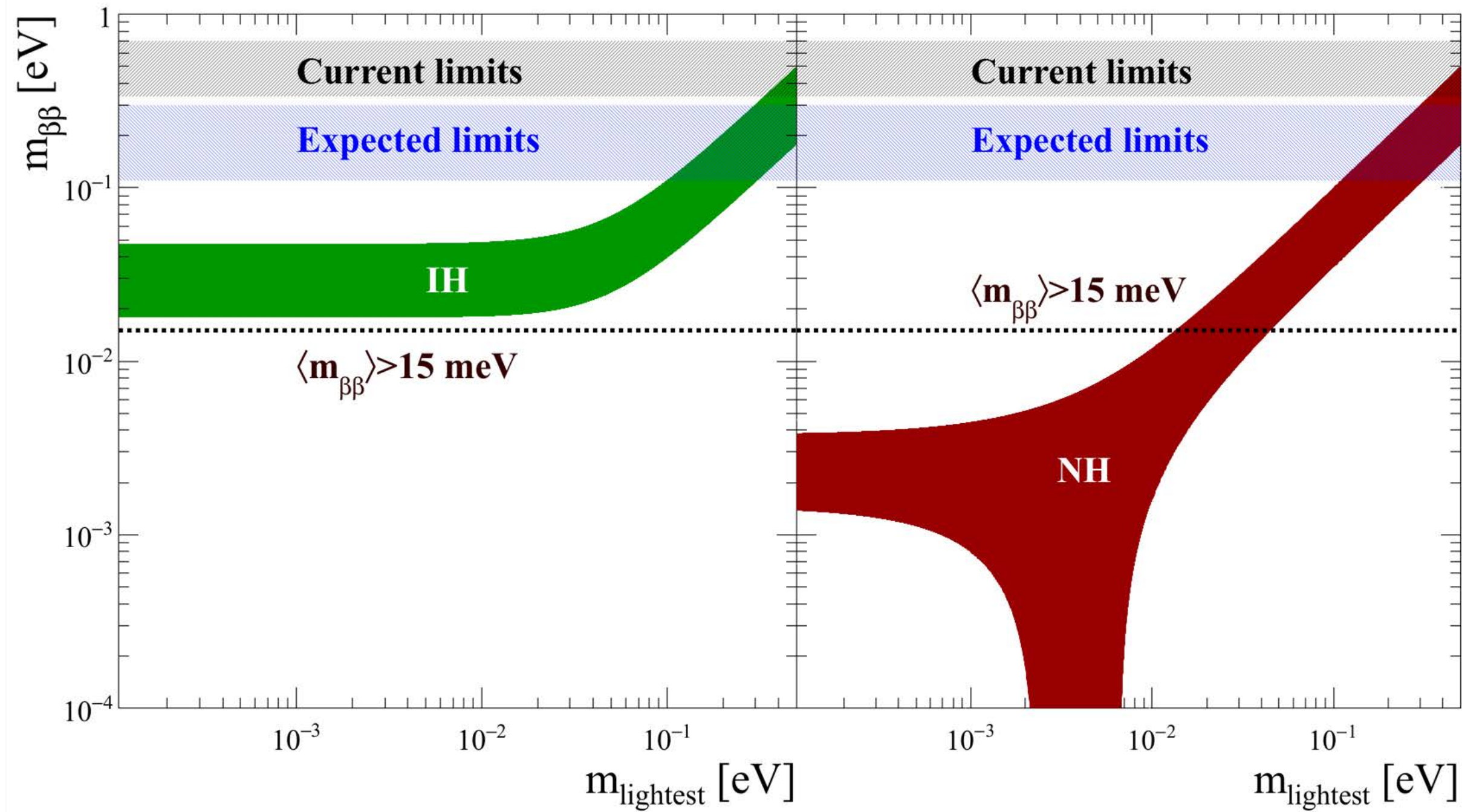


Pairing in nuclei results in a displacement of even-even and odd-odd mass parabolas for given isobars. Data from AME 2012. **Precise masses \Rightarrow precise Q value.**

REACHING FOR THE HORIZON

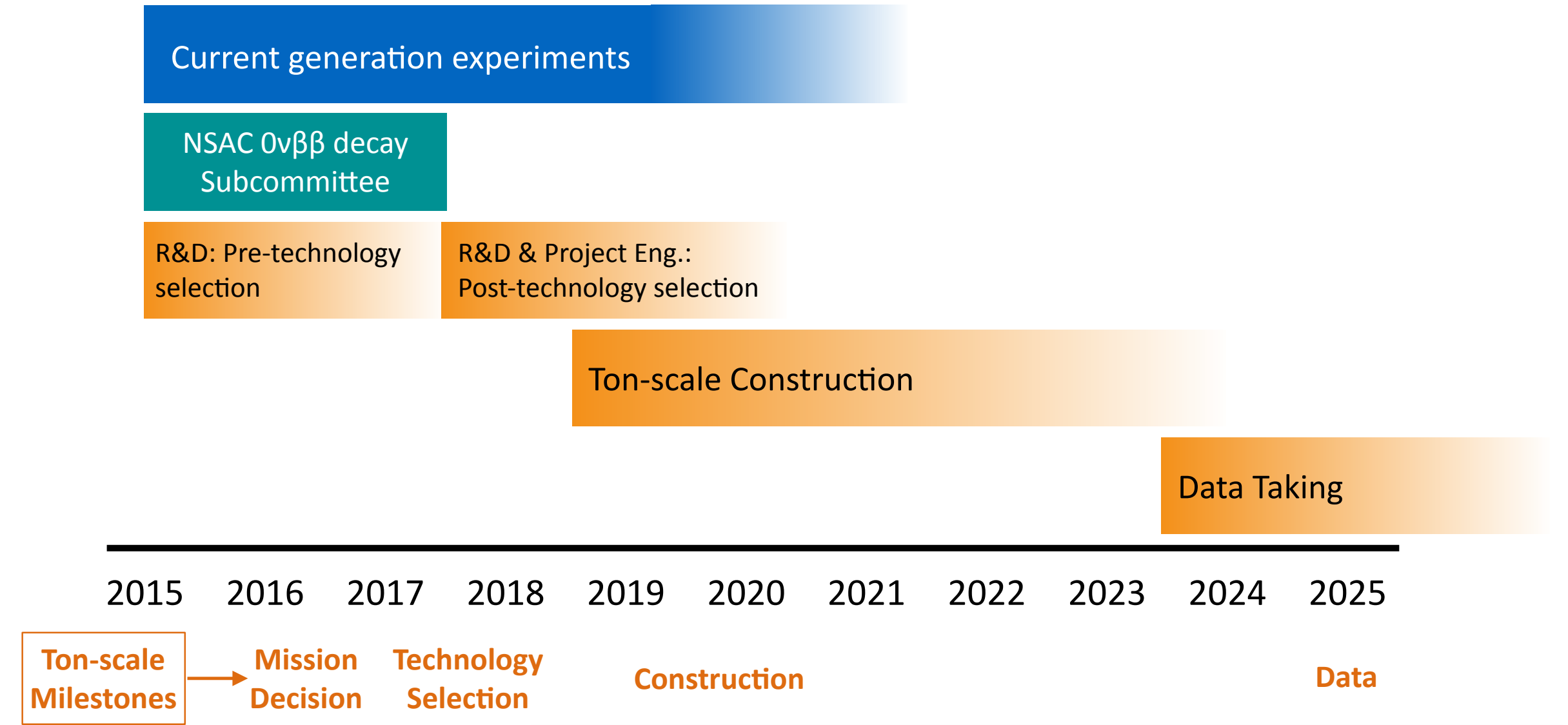
The 2015
LONG RANGE PLAN
for NUCLEAR SCIENCE





Ton-scale Neutrinoless Double Beta Decay ($0\nu\beta\beta$) - A Notional Timeline

Search for Lepton Number Violation

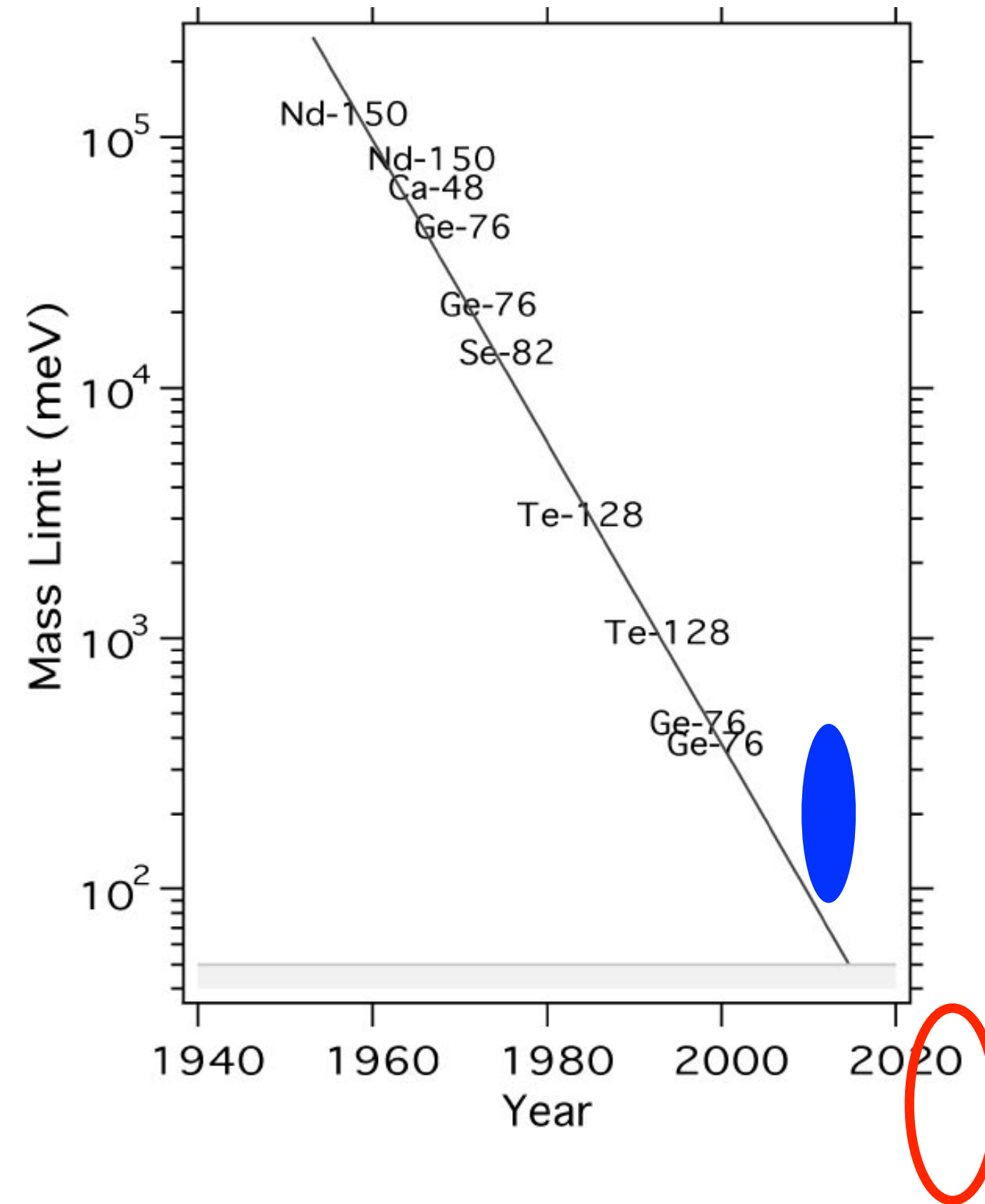
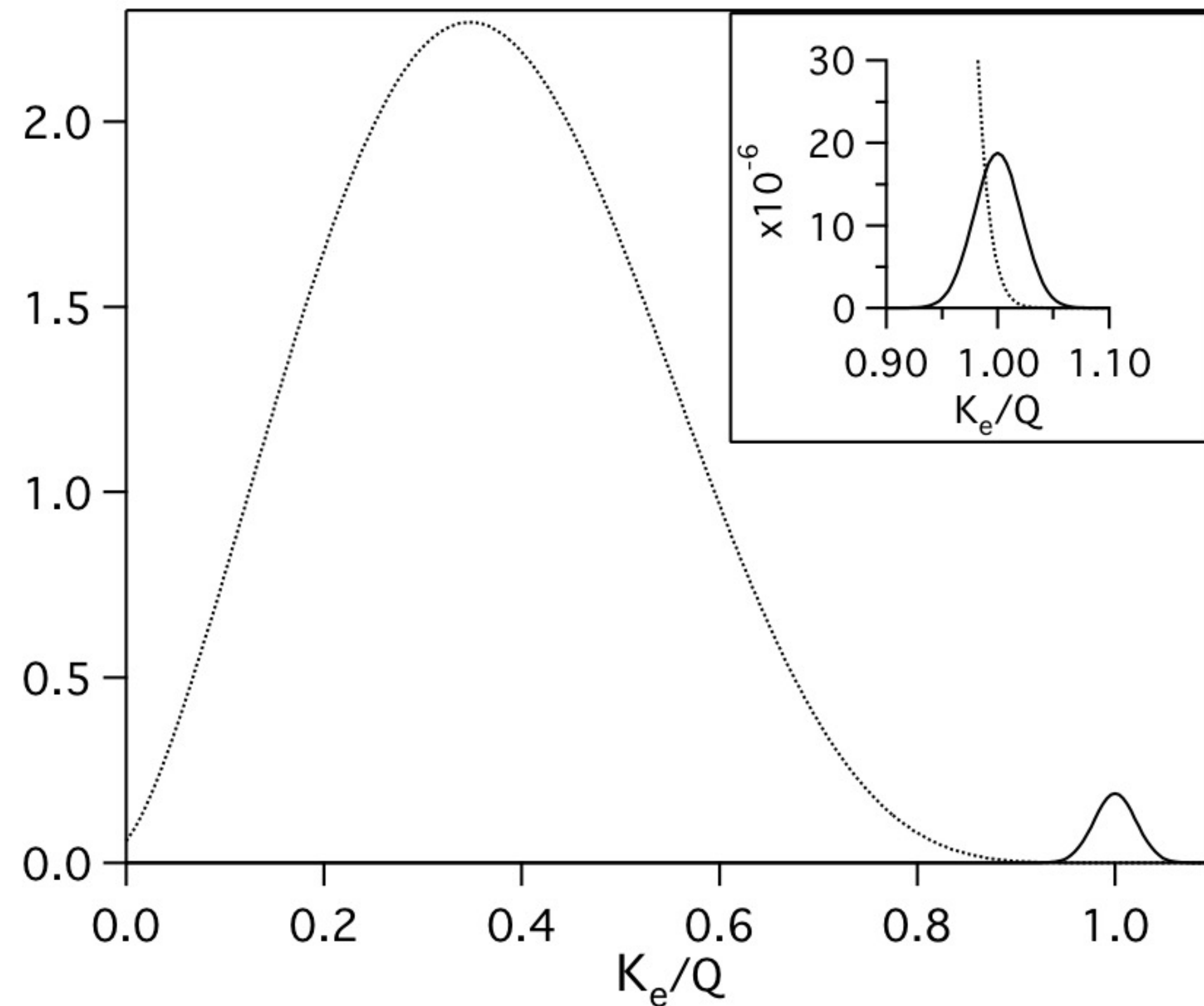


“The **second recommendation** specifically targets the development and deployment of a ton-scale neutrino-less double beta decay experiment. **Demonstration experiments at the scale of 100 kg are currently underway** to identify the requirements and candidate technologies for a larger, next-generation experiment, which is needed to be sensitive to postulated new physics. **An ongoing NSAC subcommittee is helping to guide the process of the down-select, from several current options to one U.S.-led ton-scale experiment.**”

“Construction of this flagship experiment is expected to require five years, with capital investment peaking at about \$50M/year during this period.”

“Since neutrinoless double beta decay measurements use the atomic nucleus as a laboratory, **nuclear theory is critical in connecting experimental results to the underlying lepton-number violating interactions and parameters through nuclear matrix elements**, which account for the strong interactions of neutrons and protons. Currently, **there exists about a factor of two uncertainty in the relevant matrix elements**, but by the time a ton-scale experiment is ready to take data, we expect reduced uncertainties as a result of the application to this problem of improved methods to solve the nuclear many-body physics.”

(Neutrino-less) double beta decay



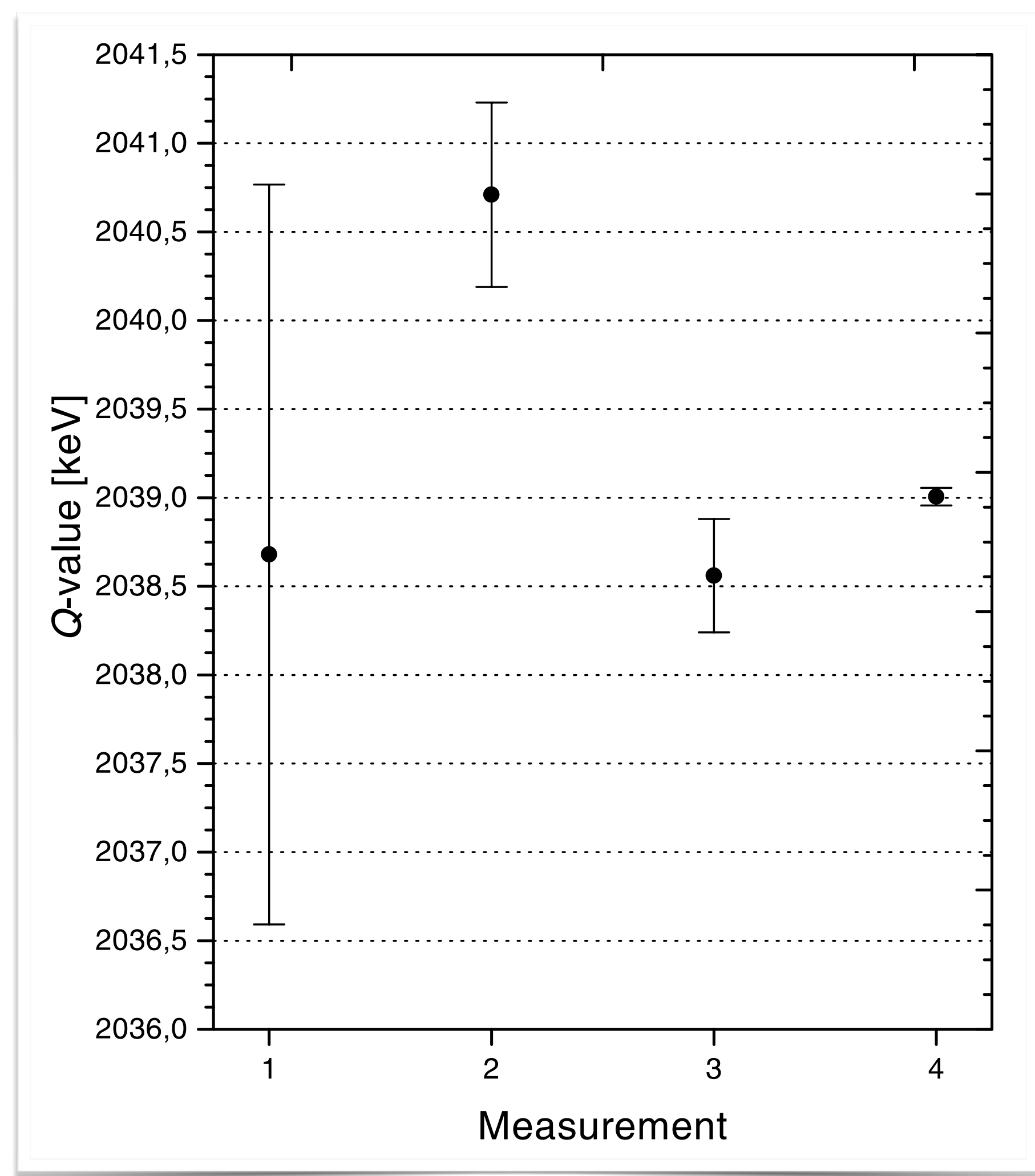
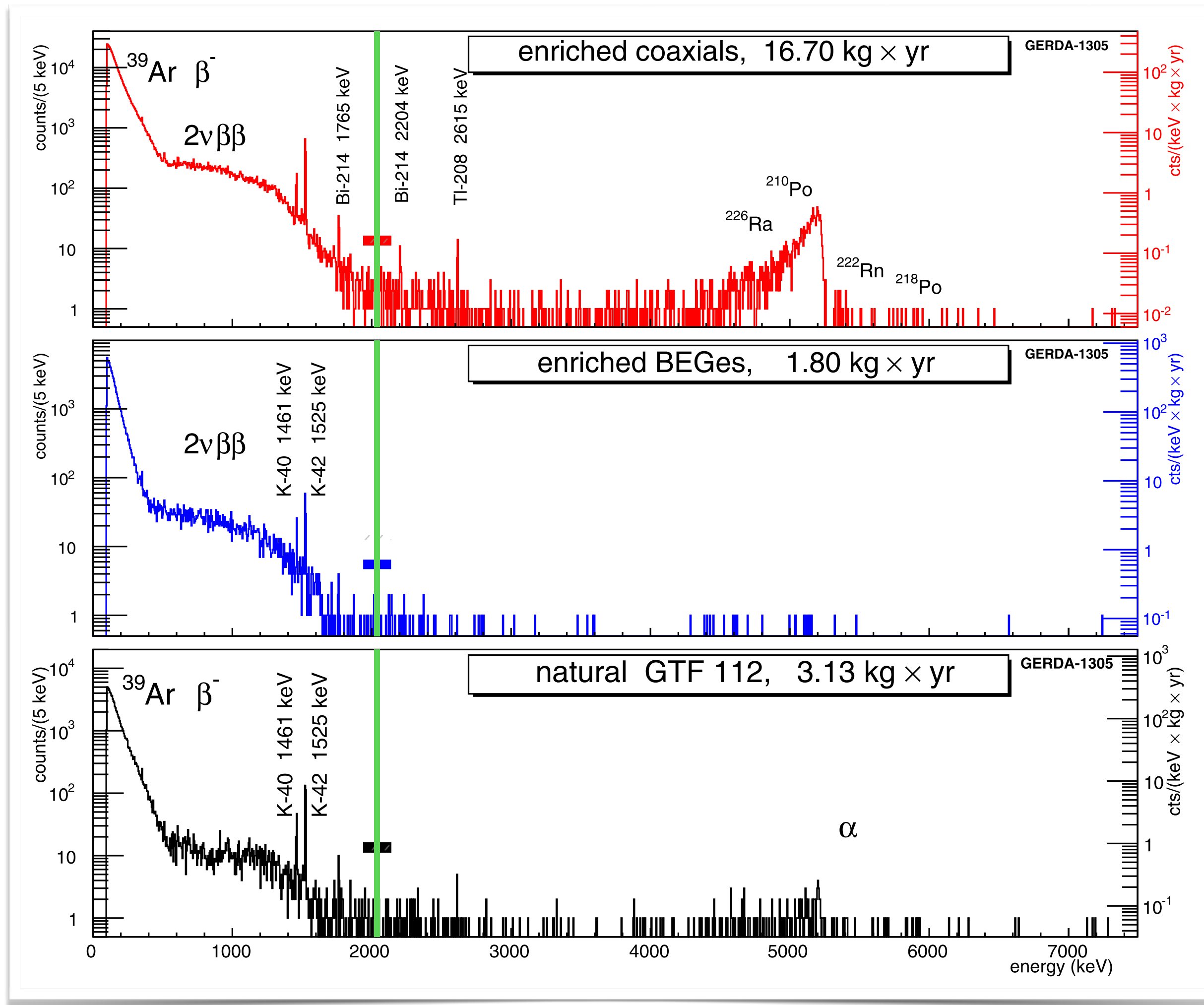
EXO-200: **200-400 meV**
 KamLAND-Zen: **61-165 meV**
 (~1 year ago)
 GERDA: **200-400 meV**.

Next Gen. > 15 meV

The search is on ... what does the future hold?

$$[T_{1/2}^{0\nu}]^{-1} = (\text{Phase Space Factor}) \times |Nuclear Matrix Element|^2 \times |\langle m_{\beta\beta} \rangle|^2$$

Knowing the Q value is essential



Q5!

$$[T_{1/2}^{0\nu}]^{-1} = (\text{Phase Space Factor}) \times |\text{Nuclear Matrix Element}|^2 \times |\langle m_{\beta\beta} \rangle|^2$$

Sizes

Neutron and weak-charge distributions of the ^{48}Ca nucleus

G. Hagen^{1,2*}, A. Ekström^{1,2}, C. Forssén^{1,2,3}, G. R. Jansen^{1,2}, W. Nazarewicz^{1,4,5}, T. Papenbrock^{1,2}, K. A. Wendt^{1,2}, S. Bacca^{6,7}, N. Barnea⁸, B. Carlsson³, C. Drischler^{9,10}, K. Hebeler^{9,10}, M. Hjorth-Jensen^{4,11}, M. Miorelli^{6,12}, G. Orlandini^{13,14}, A. Schwenk^{9,10} and J. Simonis^{9,10}

What is the size of the atomic nucleus? This deceptively simple question is difficult to answer. Although the electric charge distributions in atomic nuclei were measured accurately already half a century ago, our knowledge of the distribution of neutrons is still deficient. In addition to constraining the size of atomic nuclei, the neutron distribution also impacts the number of nuclei that can exist and the size of neutron stars. We present an *ab initio* calculation of the neutron distribution of the neutron-rich nucleus ^{48}Ca . We show that the neutron skin (difference between the radii of the neutron and proton distributions) is significantly smaller than previously thought. We also make predictions for the electric dipole polarizability and the weak form factor; both quantities that are at present targeted by precision measurements. Based on *ab initio* results for ^{48}Ca , we provide a constraint on the size of a neutron star.

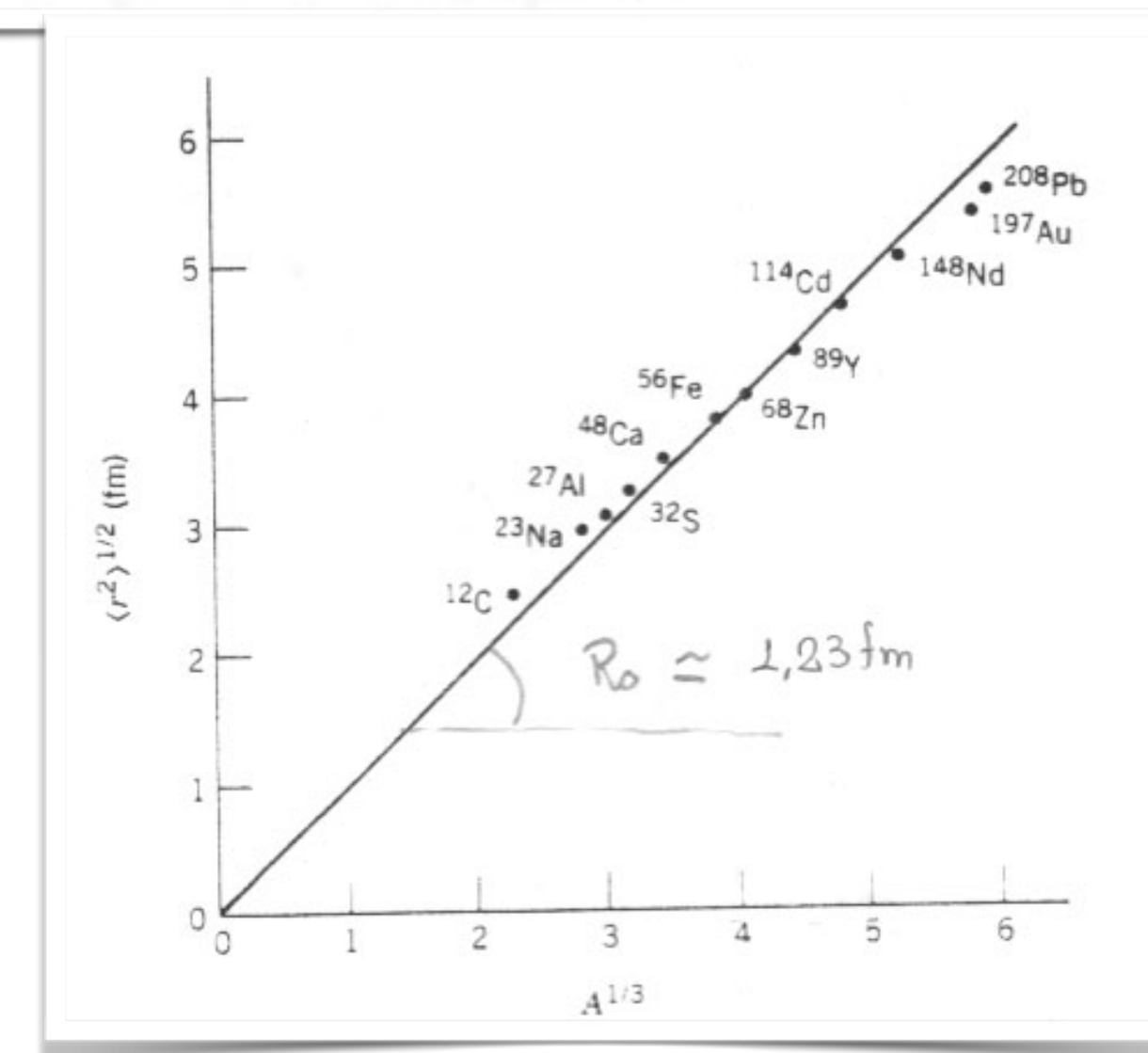
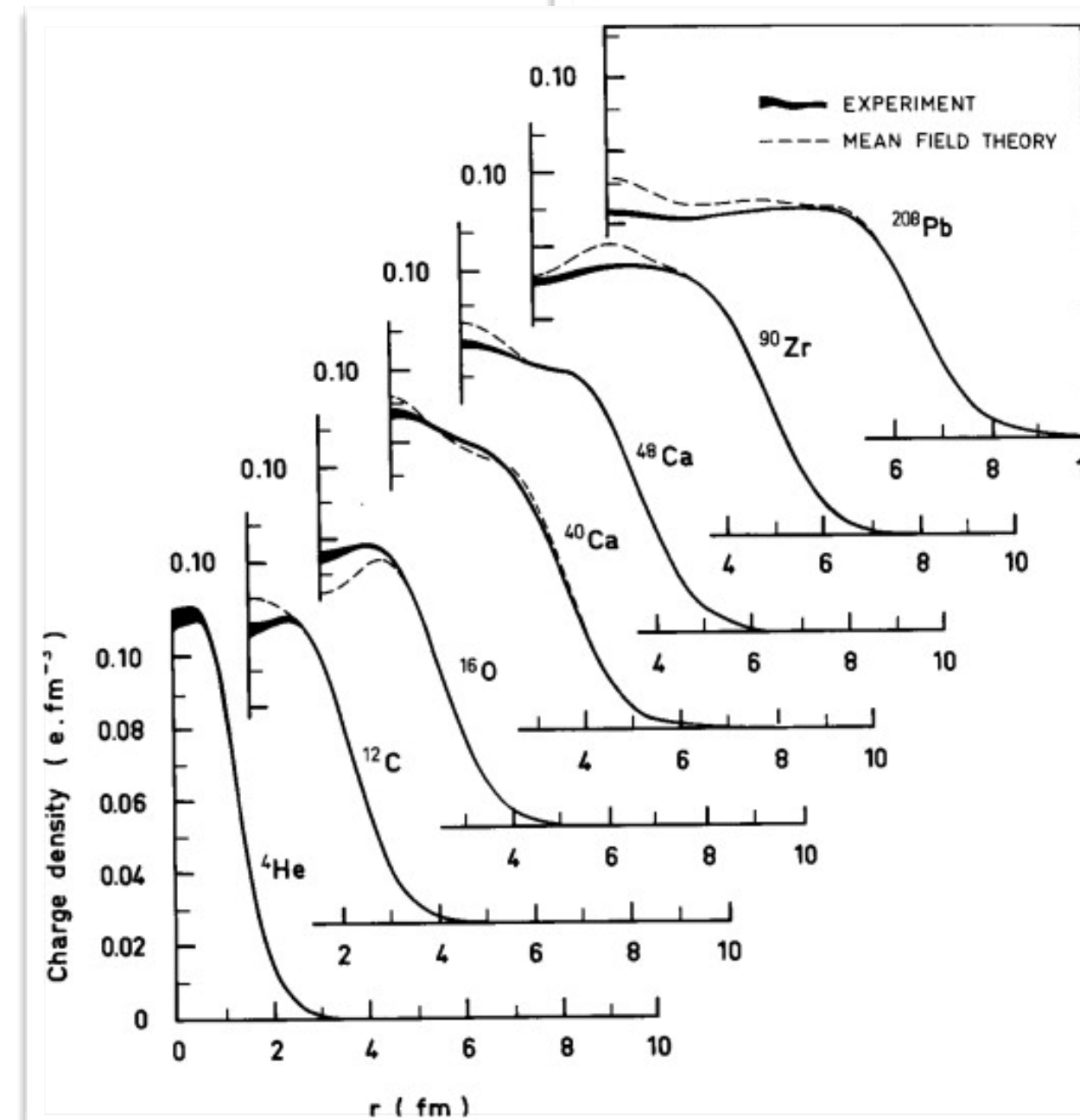
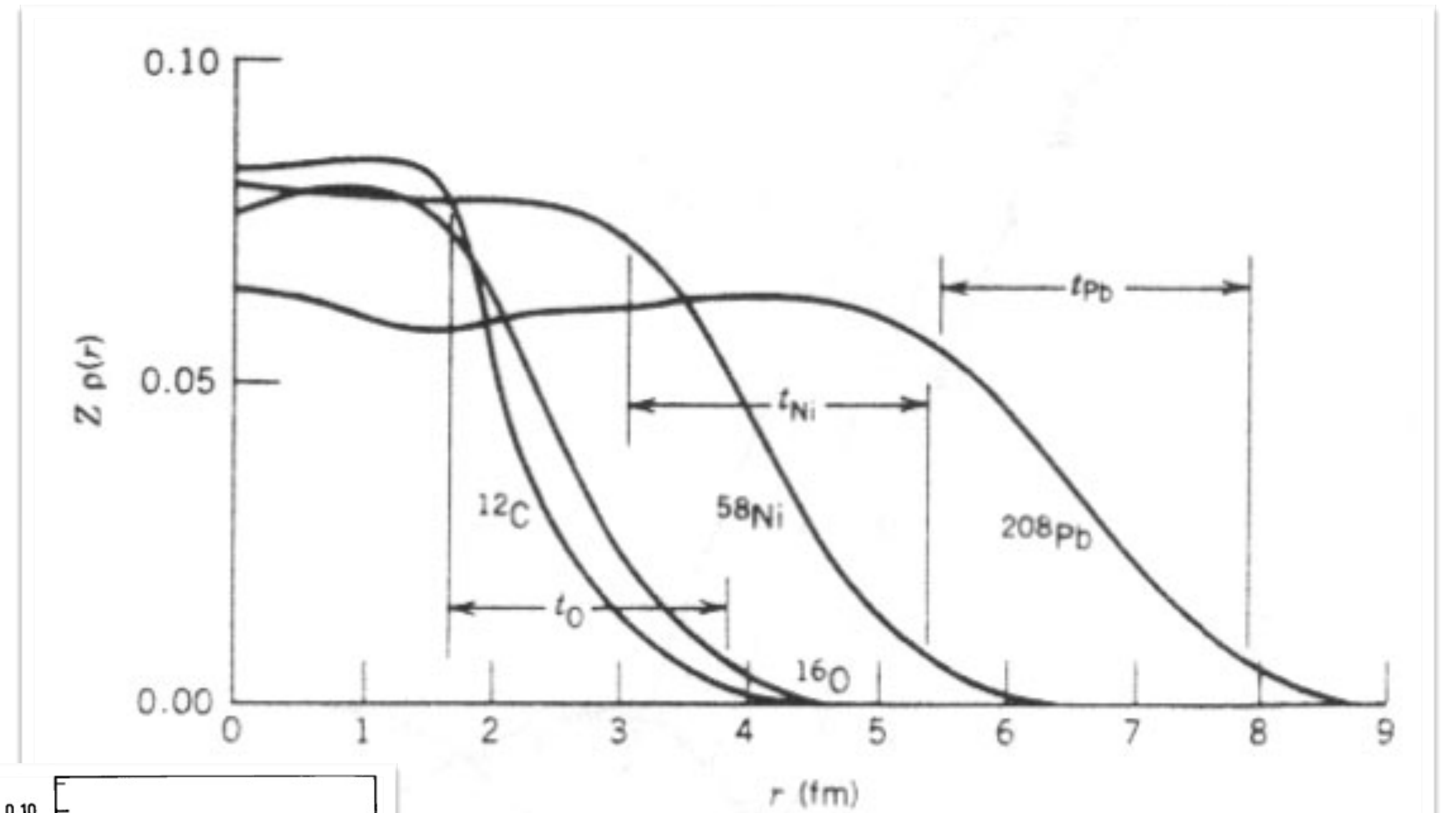
Sizes

- Sizes of nuclei along with their mass (binding) is a fundamental property of the nucleus
- Radius links to size of the nuclear potential
- Matter and charge radii are similar for most nuclei, but dramatic differences seen in exotic systems
- Neutron skins
- Matter radii, neutron structure, can modify charge radius (or center-of-mass motion)?
- Shapes of nuclei can result in changes in charge radii
- a series of examples ... via dramatic examples

Proton, neutron, matter

$$r_c^2 = \frac{1}{Z} \int \rho_c(r) r^2 d^3 r$$

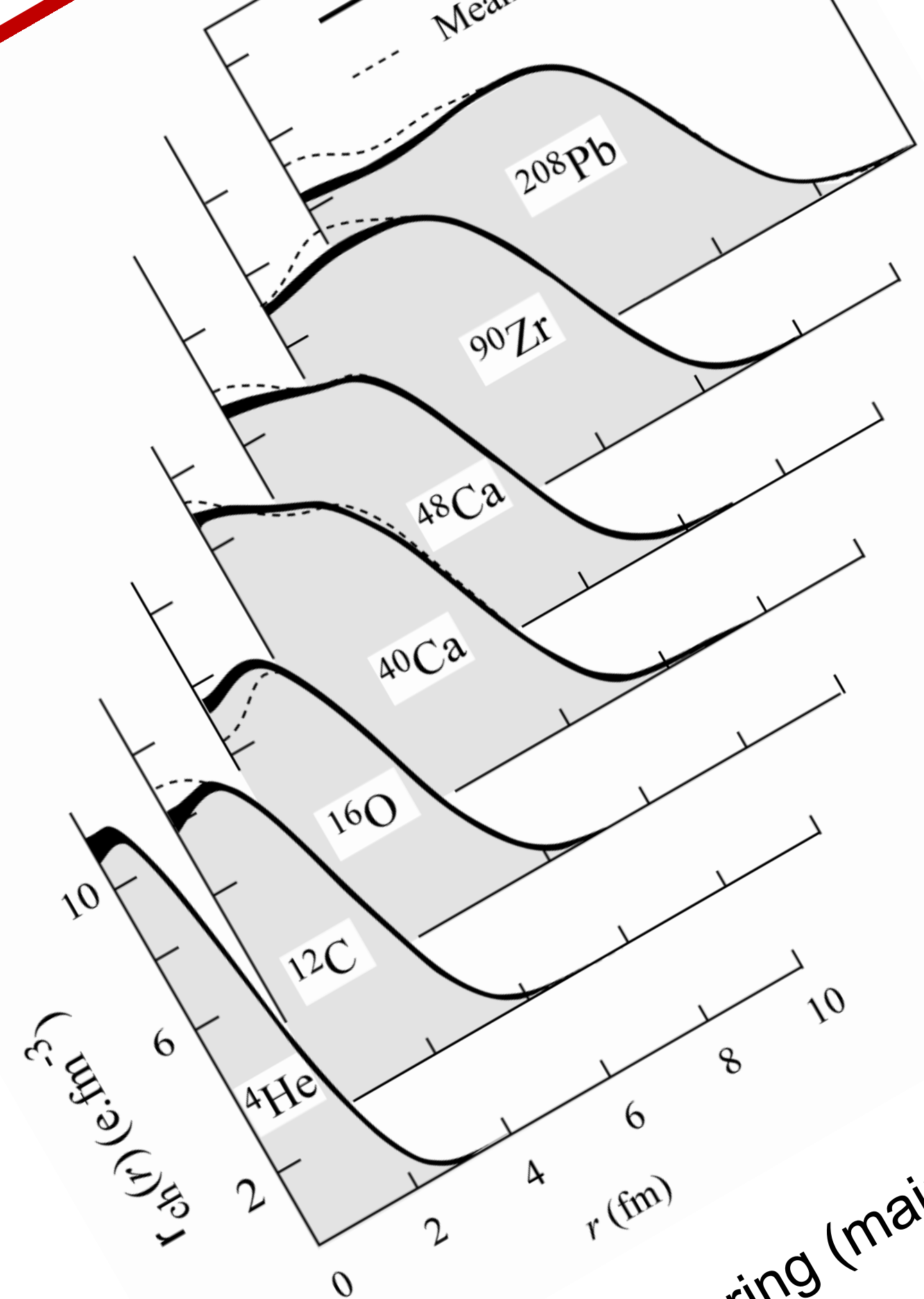
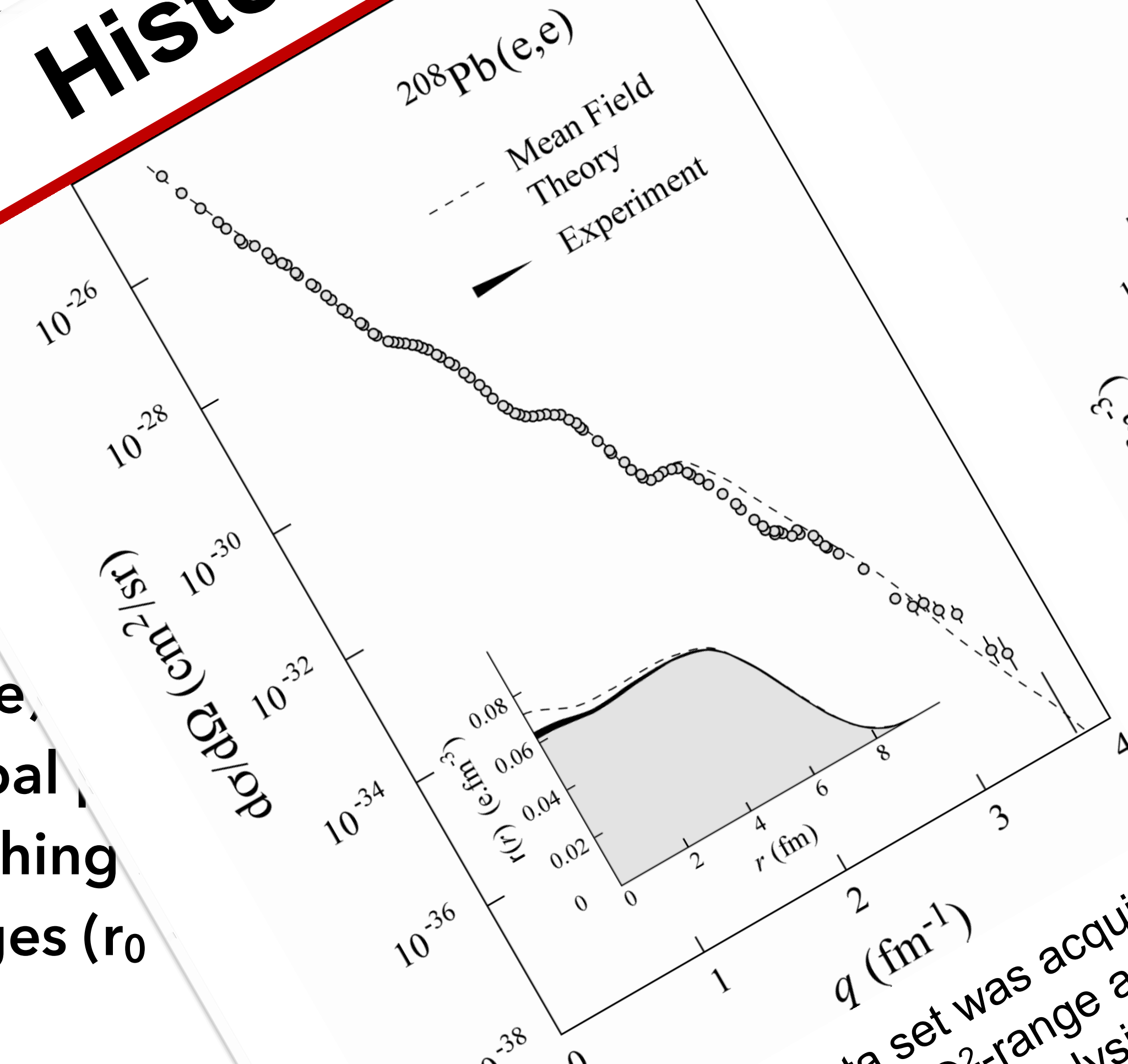
- Can be probed via elastic scattering, isotope shifts (precise)
- Matter distributions (radii) through interaction cross sections, or hadronic scattering reactions (less precise)
- A global picture of something like $r = r_0 A^{1/3}$ emerges ($r_0 \sim 1.15\text{-}3 \text{ fm}$)



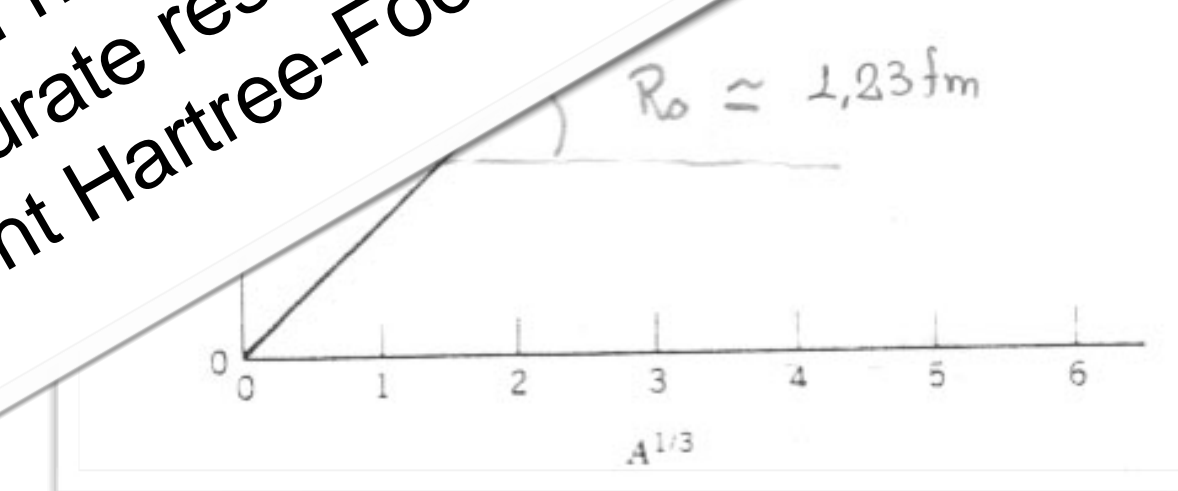
Proton, neutron, ... History: Charge Distribution

$$r_c^2 = \frac{1}{Z} \int \rho_c(r) r^2 dr$$

- the ...
- sect ...
- scatte ...
- precise ...
- A global ...
- something ...
- emerges (r_0 ...)



large data set was acquired on elastic electron scattering (mainly from ... for large Q²-range and for variety of nuclei ... "ident" analysis provided accurate results on charge distribution well ... field Density-Dependent Hartree-Fock calculations



Jefferson Lab

Krane's text book ... Rolf's slides from this school

FORM FACTORS OF NUCLEI AT LOW ENERGY

Elastic scattering $\rightarrow W^2 = M^2 \rightarrow Q^2 = 2M(E-E') \rightarrow Q^2$ and ν are correlated

$$d\sigma/d\Omega \text{ (and not } d\sigma/d\Omega dE) = \sigma_M F_0^2(q)$$

- For a point charge with charge Z one has $F_0(q) = Z$.
- For a charge with a finite size $F_0(q)$ will be smaller than Z , because different parts of $\rho(r)$ will give destructive contributions in the integral that constitutes $F_0(q)$.
- Often one includes the factor Z in σ_M and not in F_0 , such that $F_0(0) = 1$.

$$F(q) = \frac{4\pi}{Zq} \int \rho(r) \sin(qr) r dr$$

Scatter from uniform sphere with radius R at low q : $\sin(qr) = qr - (1/6)(qr)^3$

1st term disappears (charge normalization)

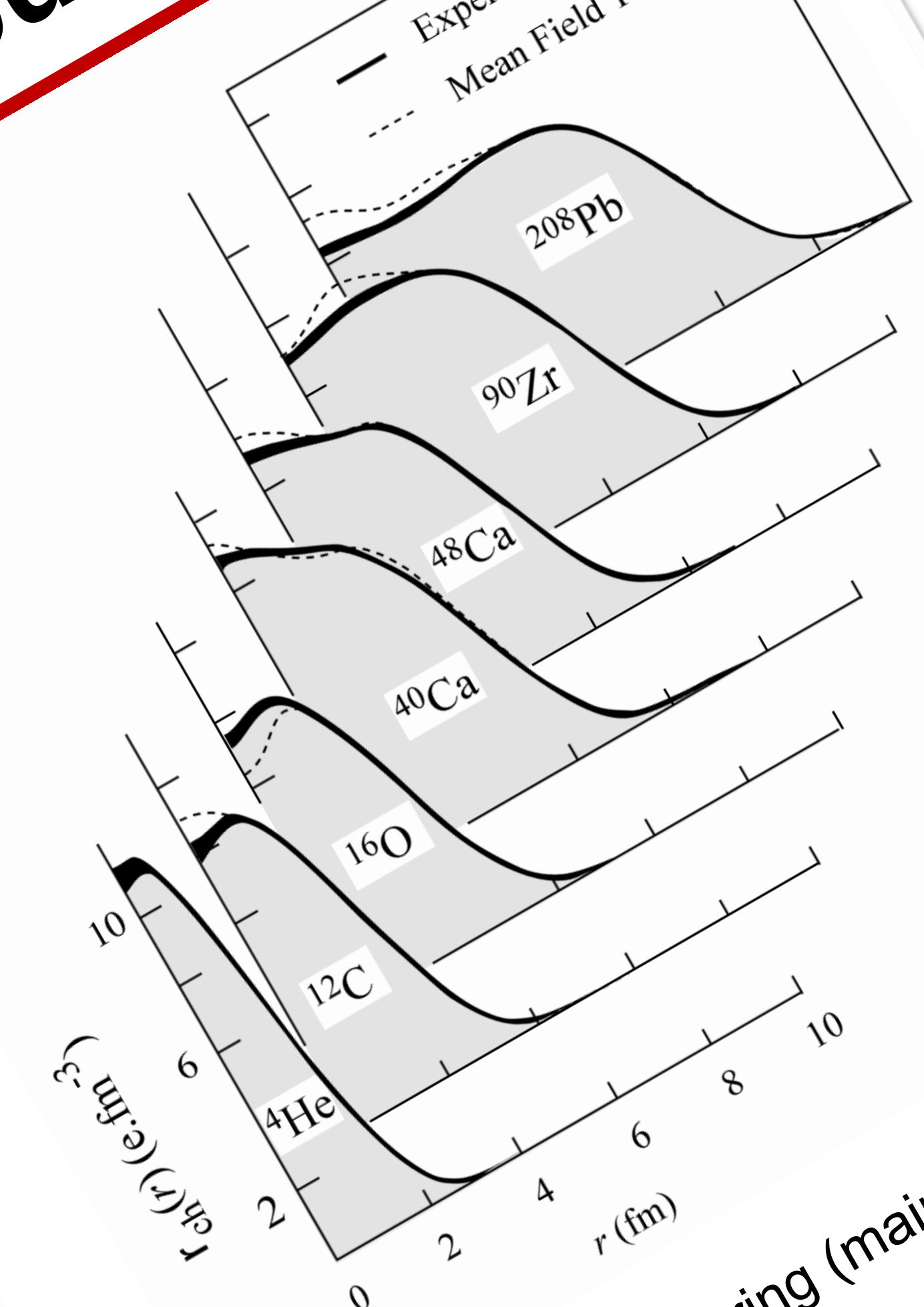
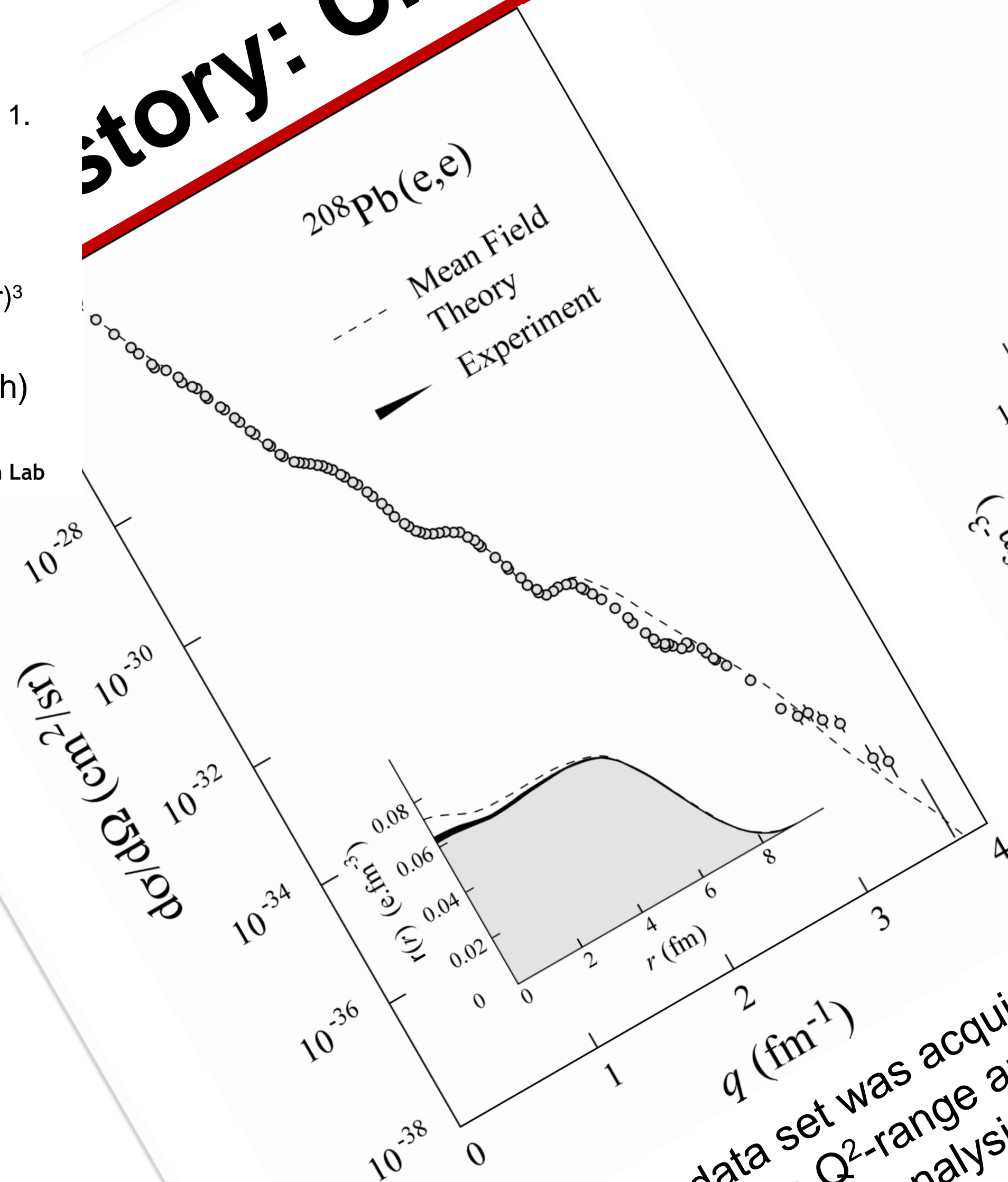
2nd term gives direct R_{RMS} measurement (for q low enough)

At higher q pattern looks like slit scattering with radius R

34

Jefferson Lab

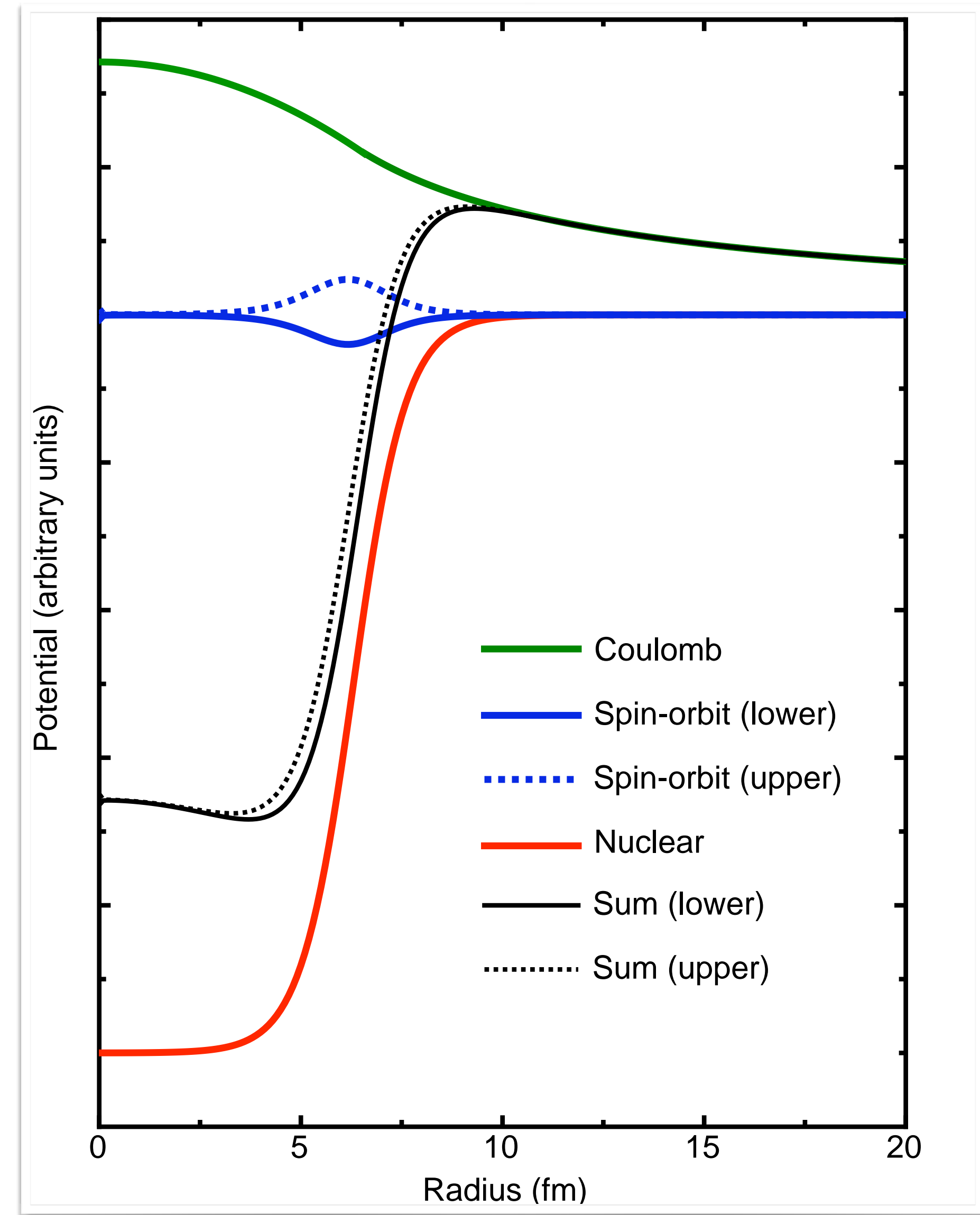
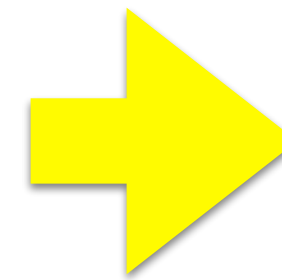
Story: Charge Distribution



data set was acquired on elastic electron scattering (mainly from the Q²-range and for variety of nuclei analysis provided accurate results on charge distribution Density-Dependent Hartree-Fock calculation

Charge radii

Target nucleus	E_x [MeV]	J^π	S (e,e'p)	r_0 [fm]	S (d, 3 He) literature	S (d, 3 He) reanalysis
^{12}C	0.000	$3/2^-$	1.72(11) [43]	1.35(2)	2.98 [44]	1.72
	2.125	$1/2^-$	0.26(2)	1.65(2)	0.69	0.27
	5.020	$3/2^-$	0.20(2)	1.51(2)	0.31	0.11
^{16}O	0.000	$1/2^-$	1.27(13) [45]	1.37(3)	2.30 [46]	1.02
	6.320	$3/2^-$	2.25(22)	1.28(2)	3.64	1.94
^{31}P	0.000	0^+	0.40(3) [47]	1.27(2)	0.62 [48]	0.36
	2.239	2^+	0.60(5)	1.18(3)	0.72	0.49
	3.498	2^+	0.28(2)	1.12(3)	0.30	0.19
^{40}Ca	0.000	$3/2^+$	2.58(19) [49,50]	1.30(5)	3.70 [51]	2.30
	2.522	$1/2^+$	1.03(7)	1.28(6)	1.65	1.03
^{51}V	0.000	$7/2^-$	0.37(3) [3]	1.30(3)	0.73 [52]	0.30 [5]
	1.554	$7/2^-$	0.16(2)	1.31(4)	0.39	0.15
	2.675	$7/2^-$	0.33(3)	1.32(3)	0.64	0.26
	3.199	$7/2^-$	0.49(4)	1.34(3)	1.05	0.39
	4.410	$1/2^+$	0.28(3)	1.22(3)	0.63	0.22
^{90}Zr	6.045	$1/2^+$	0.35(3)	1.27(4)	1.10	0.30
	0.000	$1/2^-$	0.72(7) [3]	1.32(3)	1.80 [53]	0.60 [54]
	0.909	$9/2^+$	0.54(5)	1.31(2)	1.25	0.30
^{142}Nd	1.507	$3/2^-$	1.86(14)	1.27(2)	3.90	1.20
	1.745	$5/2^-$	2.77(19)	1.30(2)	8.90	2.40
	0.000	$5/2^+$	1.39(23) [55]	1.29(9)	2.53 [56]	1.25
	0.145	$7/2^+$	3.14(43)	1.26(8)	6.28	3.79
^{206}Pb	1.118	$11/2^-$	0.56(7)	1.28(8)	0.74	0.36
	1.300	$1/2^+$	0.05(1)	1.26(9)	0.11	0.07
	0.000	$1/2^+$	0.68(6) [57]	1.23(9)	1.15 [58]	1.03
^{208}Pb	0.203	$3/2^+$	1.10(9)	1.27(9)	1.77	0.99
	0.616	$5/2^+$	0.32(3)	1.23(8)	0.52	0.44
	1.151	$3/2^+$	0.52(5)	1.28(9)	0.66	0.37
	1.479	$11/2^-$	3.58(32)	1.25(9)	6.94	5.21
	0.000	$1/2^+$	0.98(9) [57]	1.25(8)	1.8 [59]	1.5
^{208}Pb	0.350	$3/2^+$	2.31(22)	1.23(8)	3.8	2.2
	1.350	$11/2^-$	6.85(68)	1.16(9)	7.7	5.4
	1.670	$5/2^+$	2.93(28)	1.19(8)	3.5	3.1
	3.470	$7/2^+$	2.06(20)	1.15(9)	3.5	2.9



He isotopes

Which is larger charge radius? ${}^4\text{He}$, ${}^6\text{He}$, ${}^8\text{He}$?

Which is larger matter radius? ${}^4\text{He}$, ${}^6\text{He}$, ${}^8\text{He}$?

... and how were they determined?

He isoto

Which is larg

Which is larg

... and how v

REVIEWS OF MODERN PHYSICS, VOLUME 85, OCTOBER–DECEMBER 2013

Colloquium: Laser probing of neutron-rich nuclei in light atoms

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Z.-C. Yan

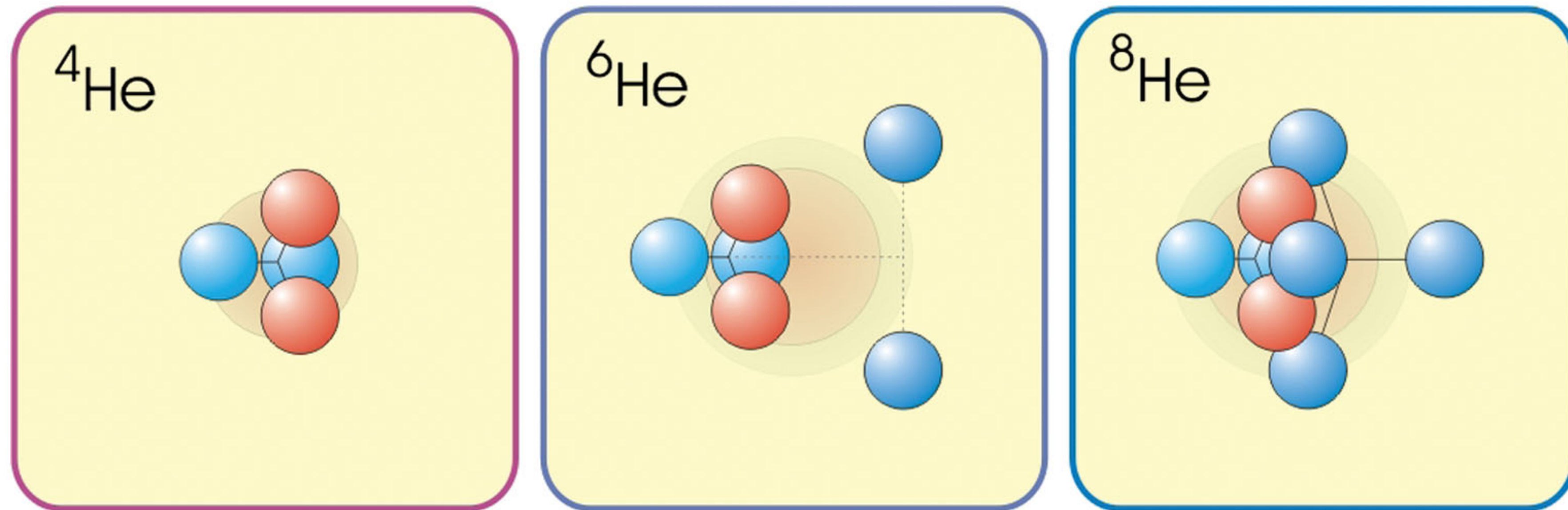
*State Key Laboratory of Magnetic Resonance and Atomic and Molecular Physics,
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Chinese Academy of Sciences, Wuhan 430071, China, and Department of Physics,
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(published 2 October 2013)

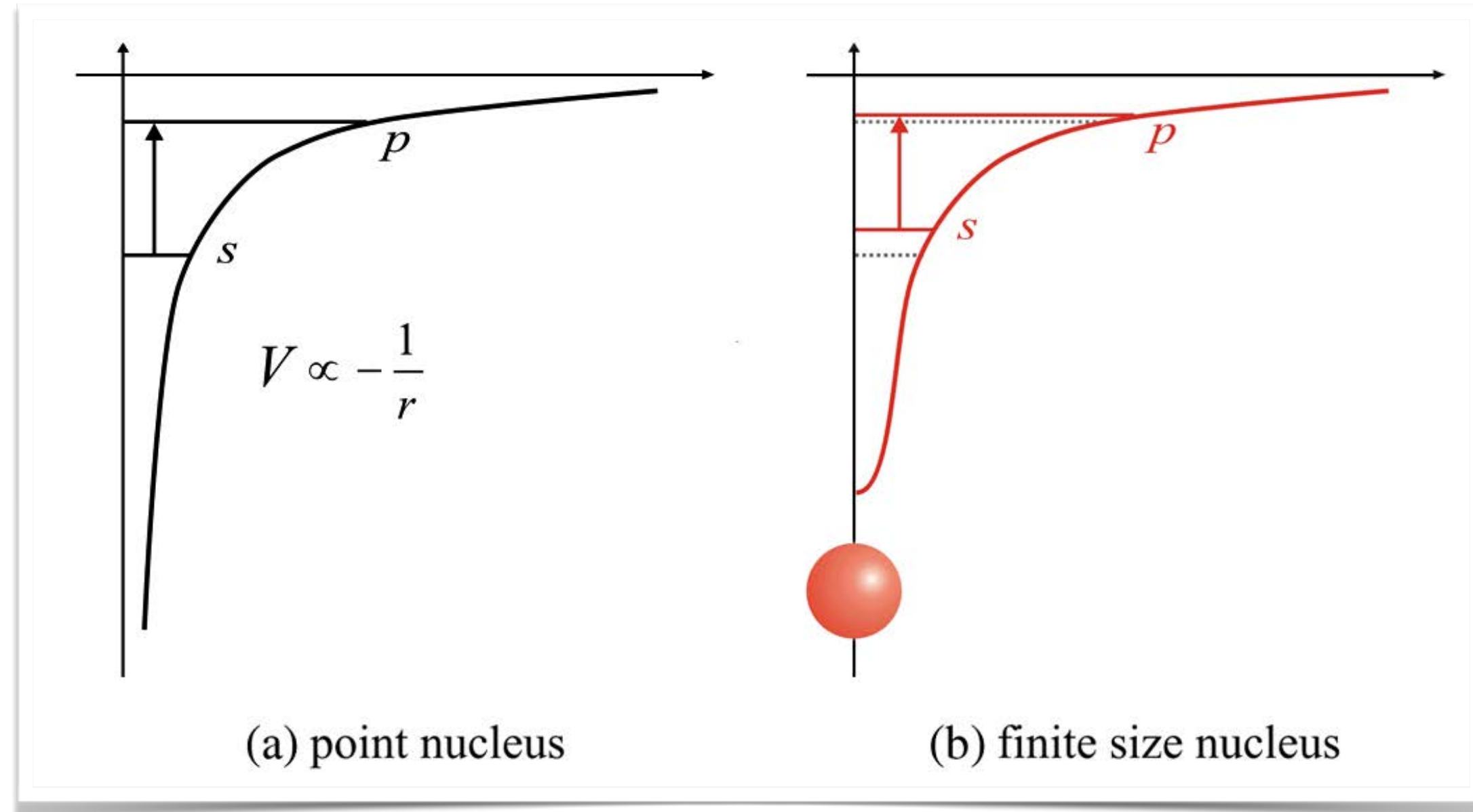
He isotopes

Which is larger charge radius? ${}^4\text{He}$, ${}^6\text{He}$, ${}^8\text{He}$?

Which is larger matter radius? ${}^4\text{He}$, ${}^6\text{He}$, ${}^8\text{He}$?

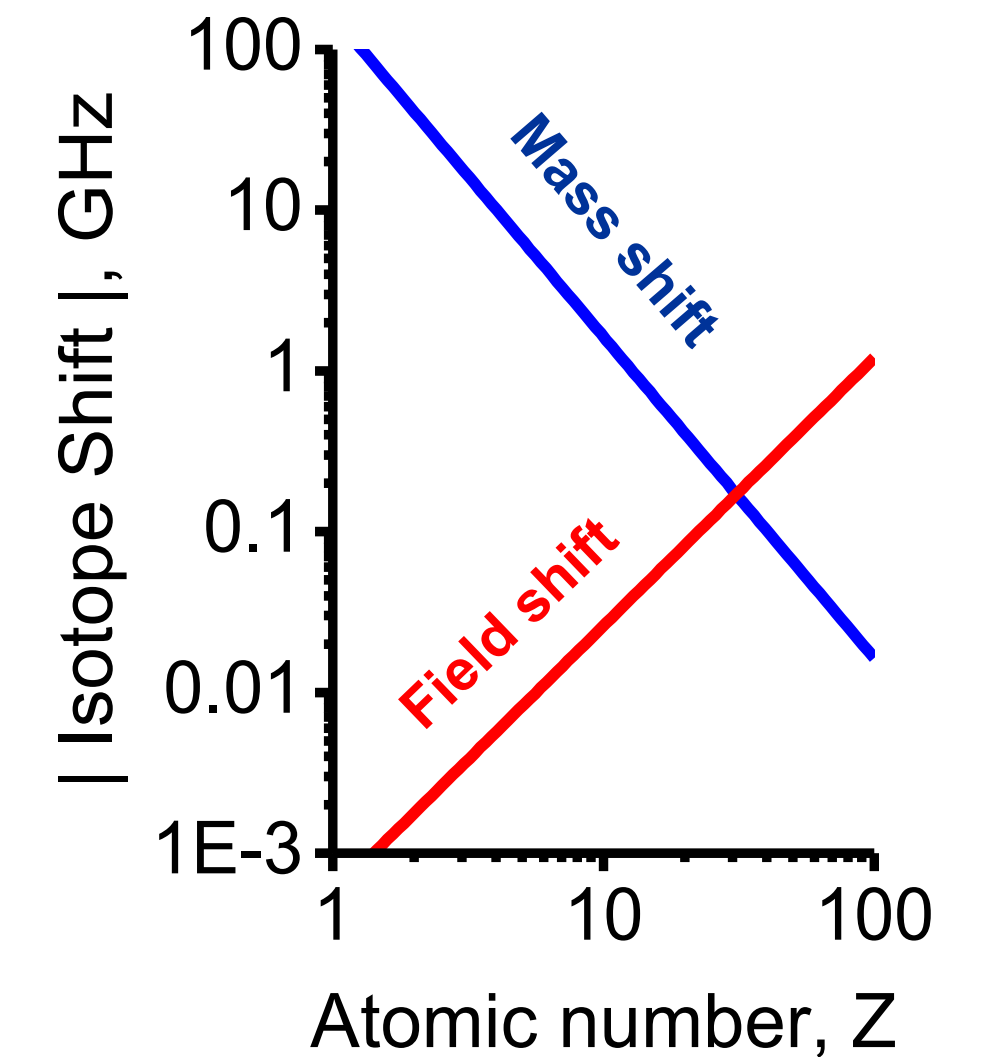
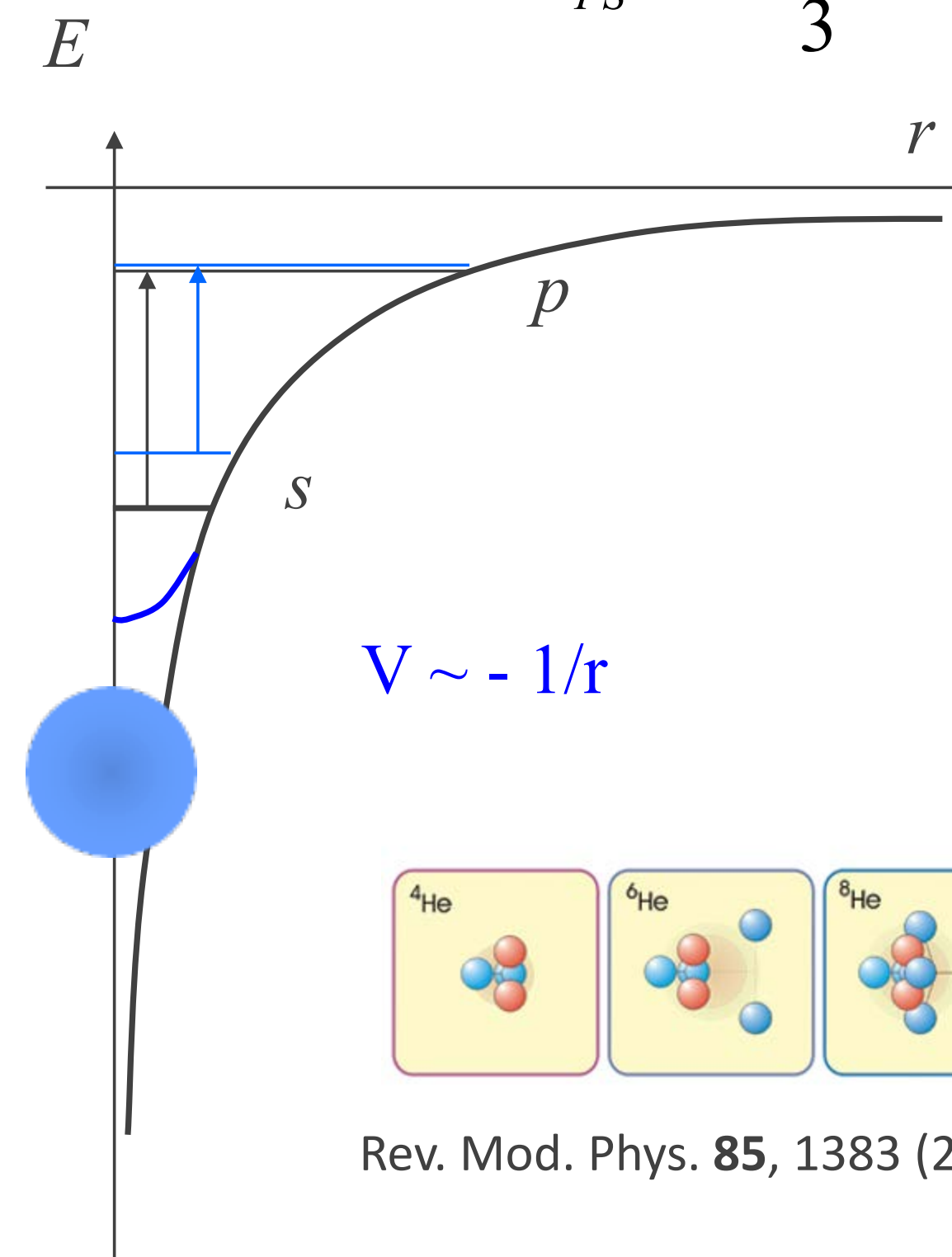


Charge radii, isotope shift



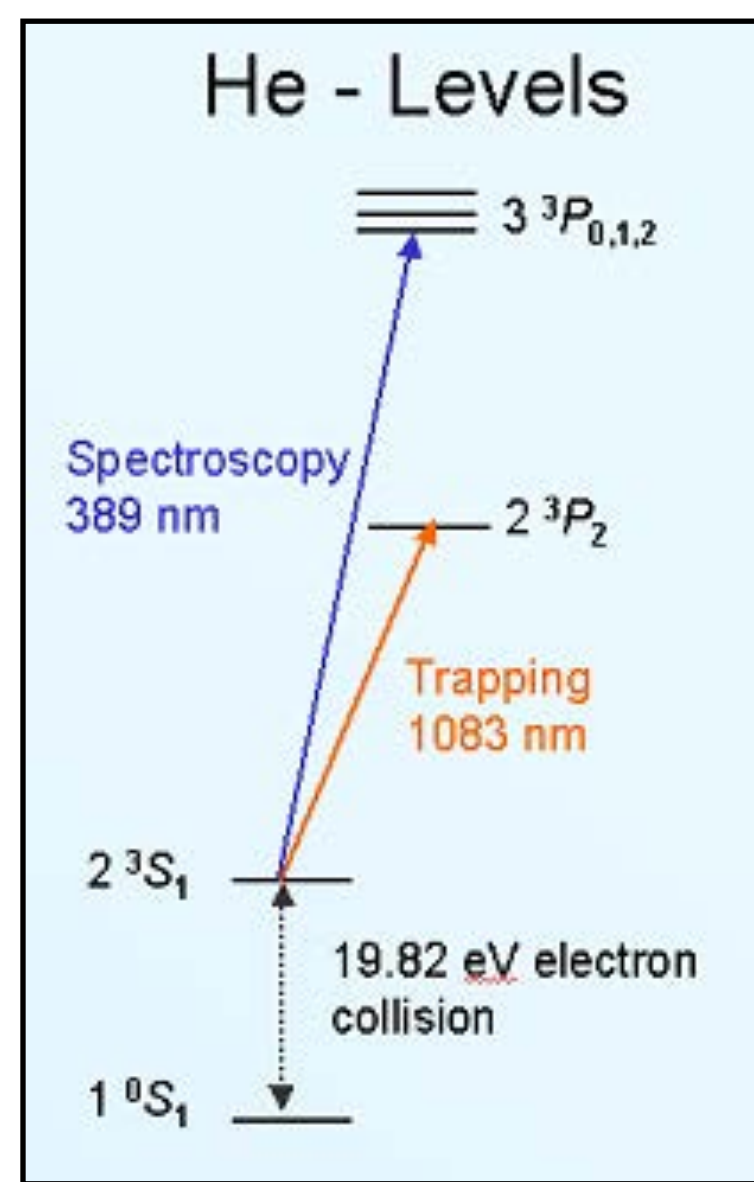
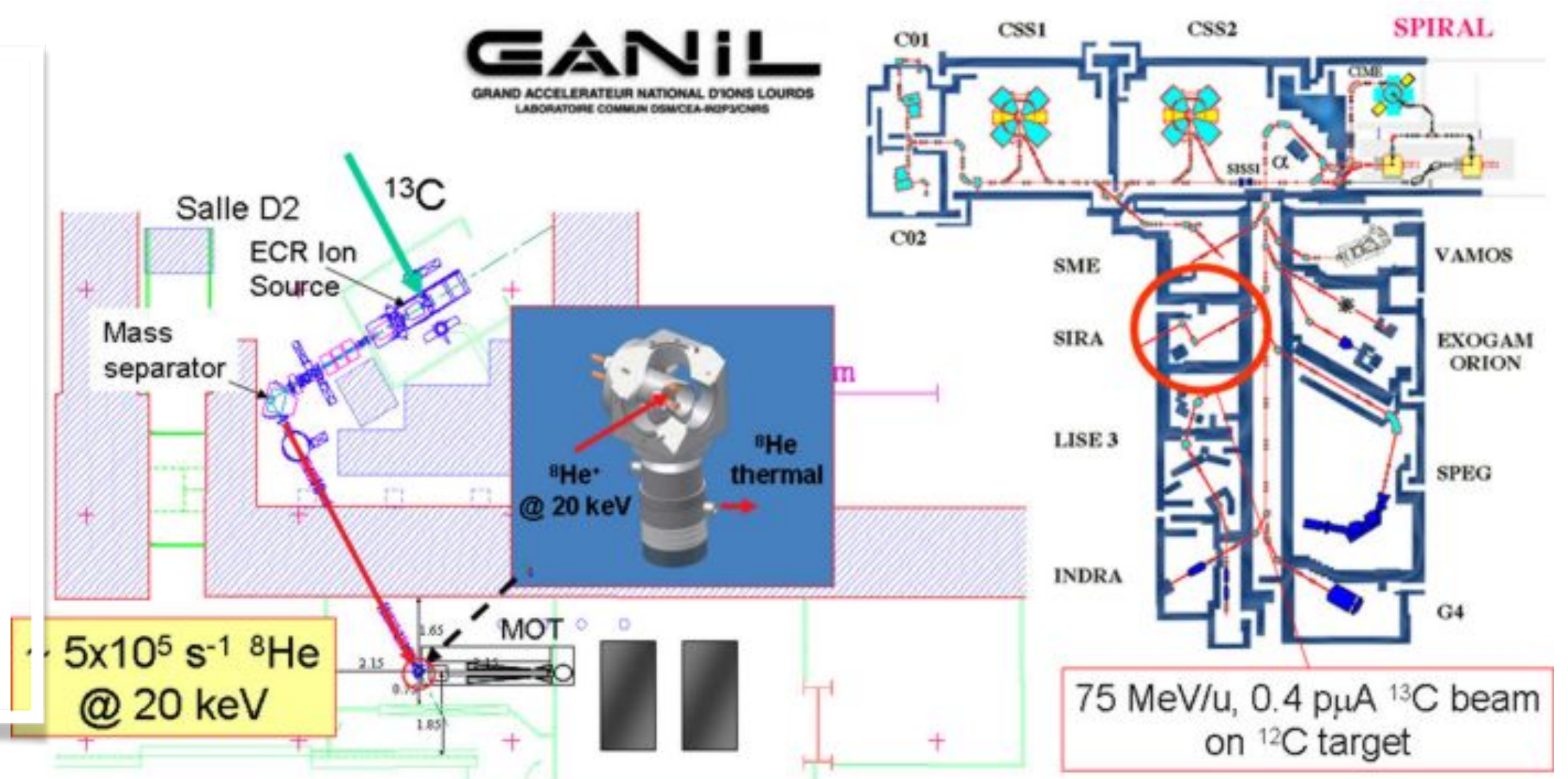
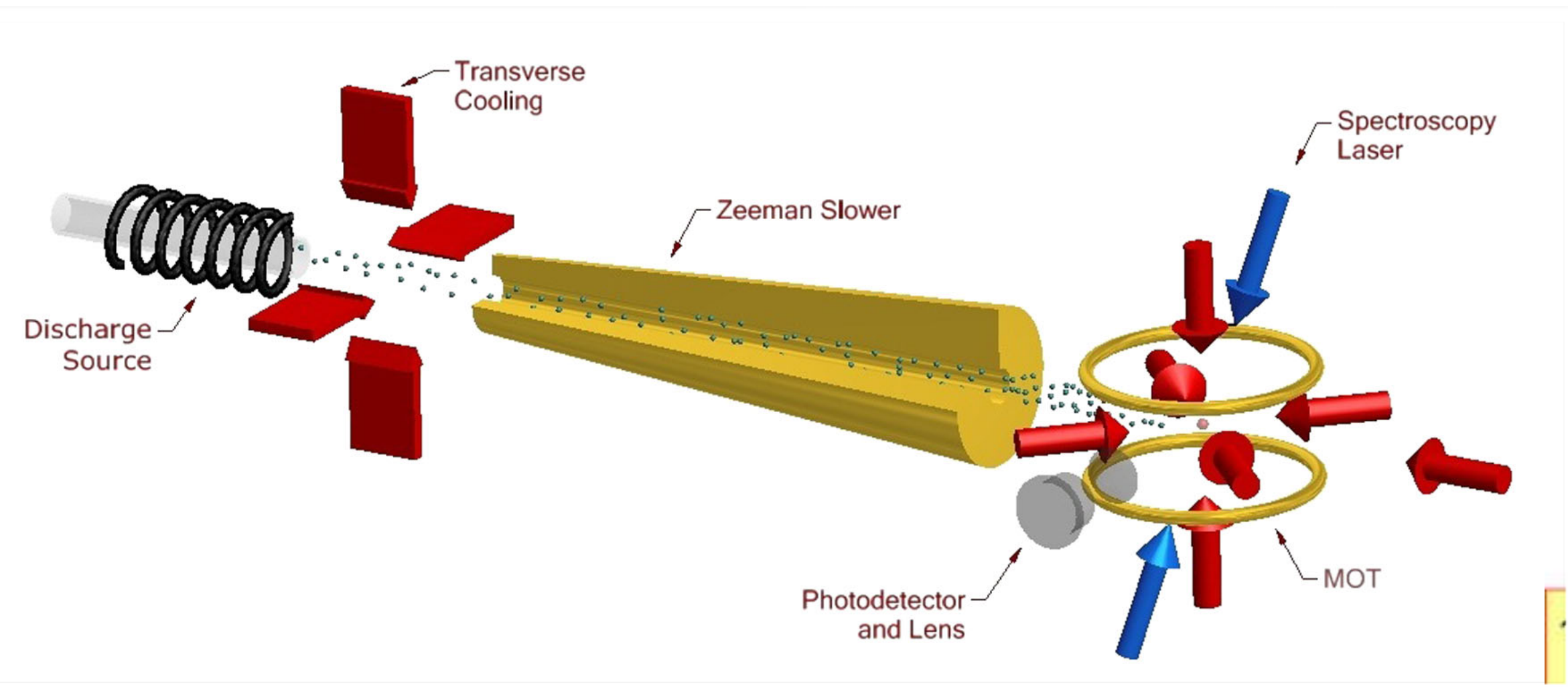
Field (Volume) Shift

$$\delta v_{FS} = -\frac{2\pi}{3} Ze^2 \cdot \Delta|\Psi(0)|^2 \cdot \delta\langle r^2 \rangle^{AA'}$$

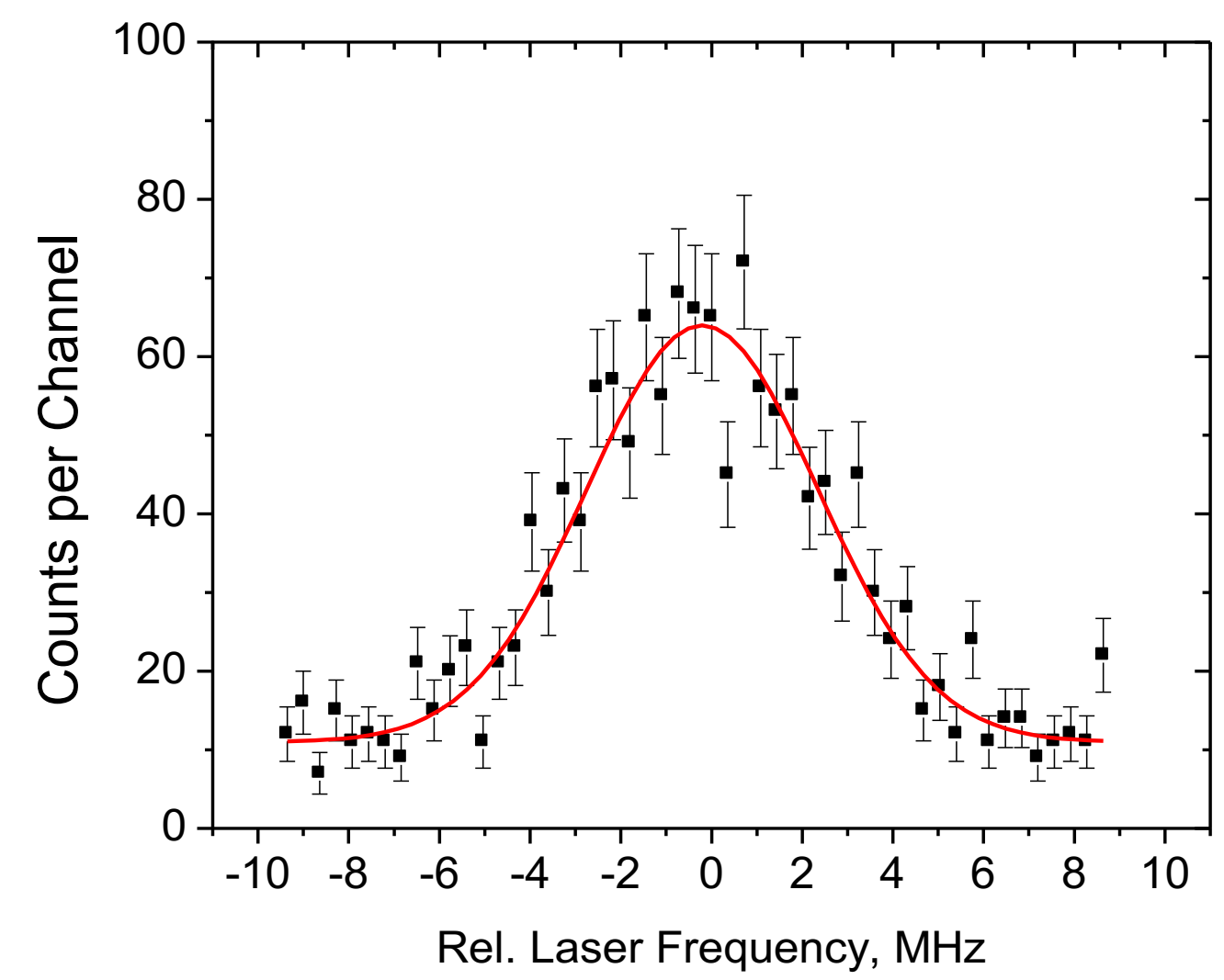
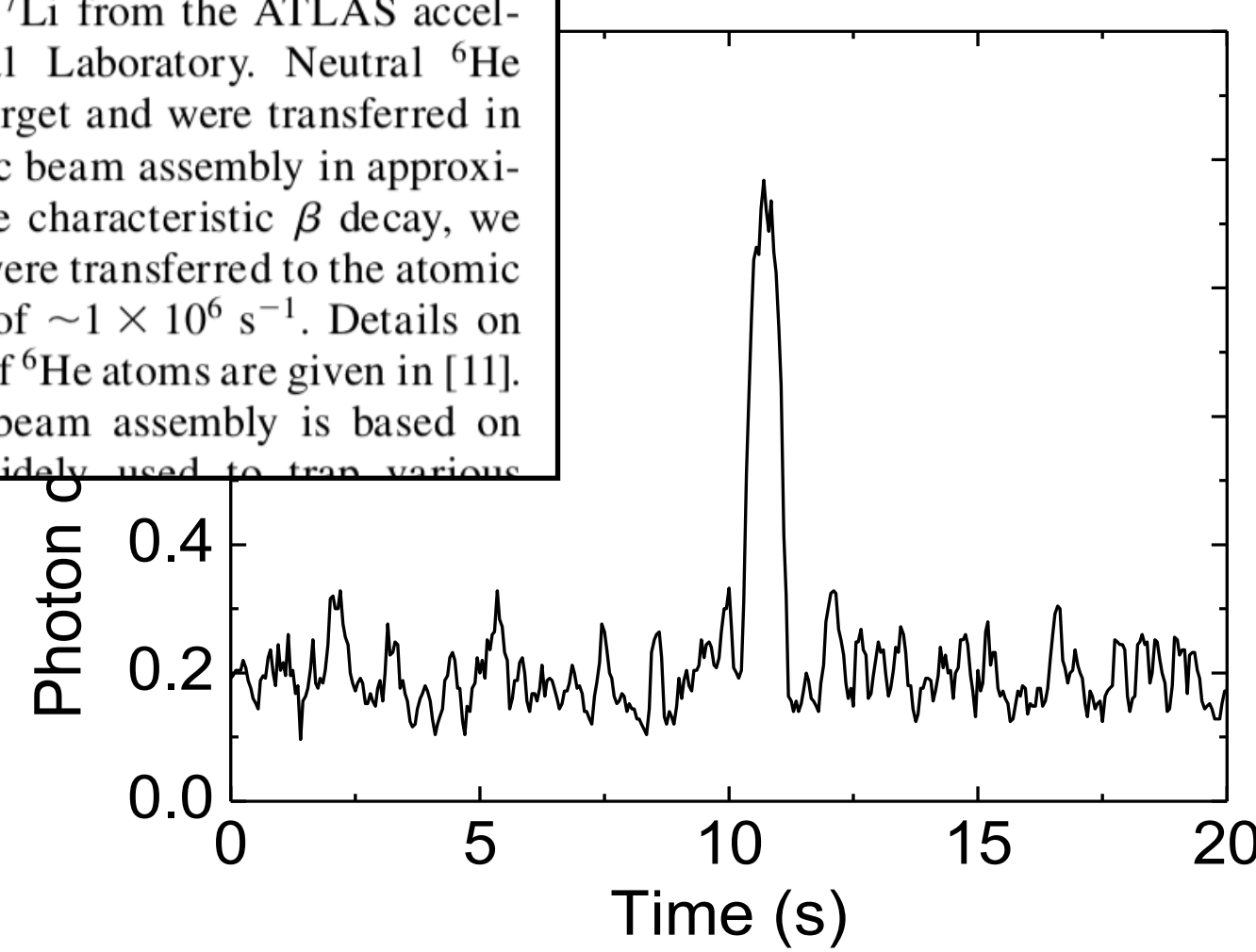


- Field shift, mass shift, and radius

E.g. ^8He



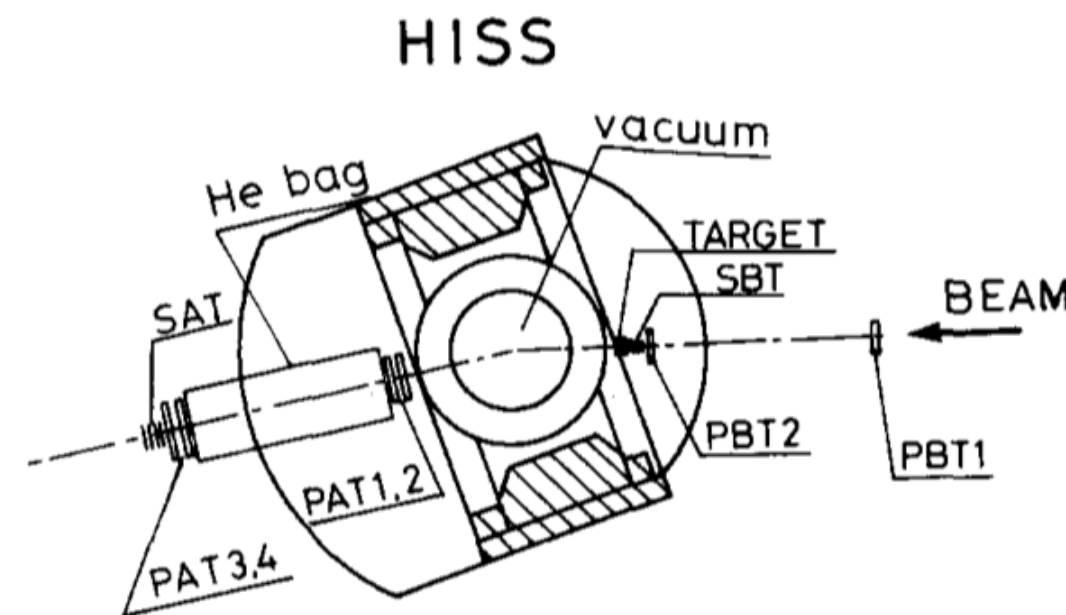
work, ^6He nuclei were produced in a hot (750 °C) graphite target via the $^{12}\text{C}(^7\text{Li}, ^6\text{He})^{13}\text{N}$ reaction with a 100 pA, 60 MeV beam of ^7Li from the ATLAS accelerator at Argonne National Laboratory. Neutral ^6He atoms diffused out of the target and were transferred in vacuum to the nearby atomic beam assembly in approximately 1 s. By detecting the characteristic β decay, we established that ^6He atoms were transferred to the atomic beam assembly at the rate of $\sim 1 \times 10^6 \text{ s}^{-1}$. Details on the production and transfer of ^6He atoms are given in [11]. Our design of the atomic beam assembly is based on a type of MOT system widely used to trap various



Matter radii

- Matter radii can be determined from interaction cross sections

$$\sigma_I(p, t) = \pi [R_I(p) + R_I(t)]^2$$



Volume 160B, number 6

PHYSICS LETTERS

17 October 1985

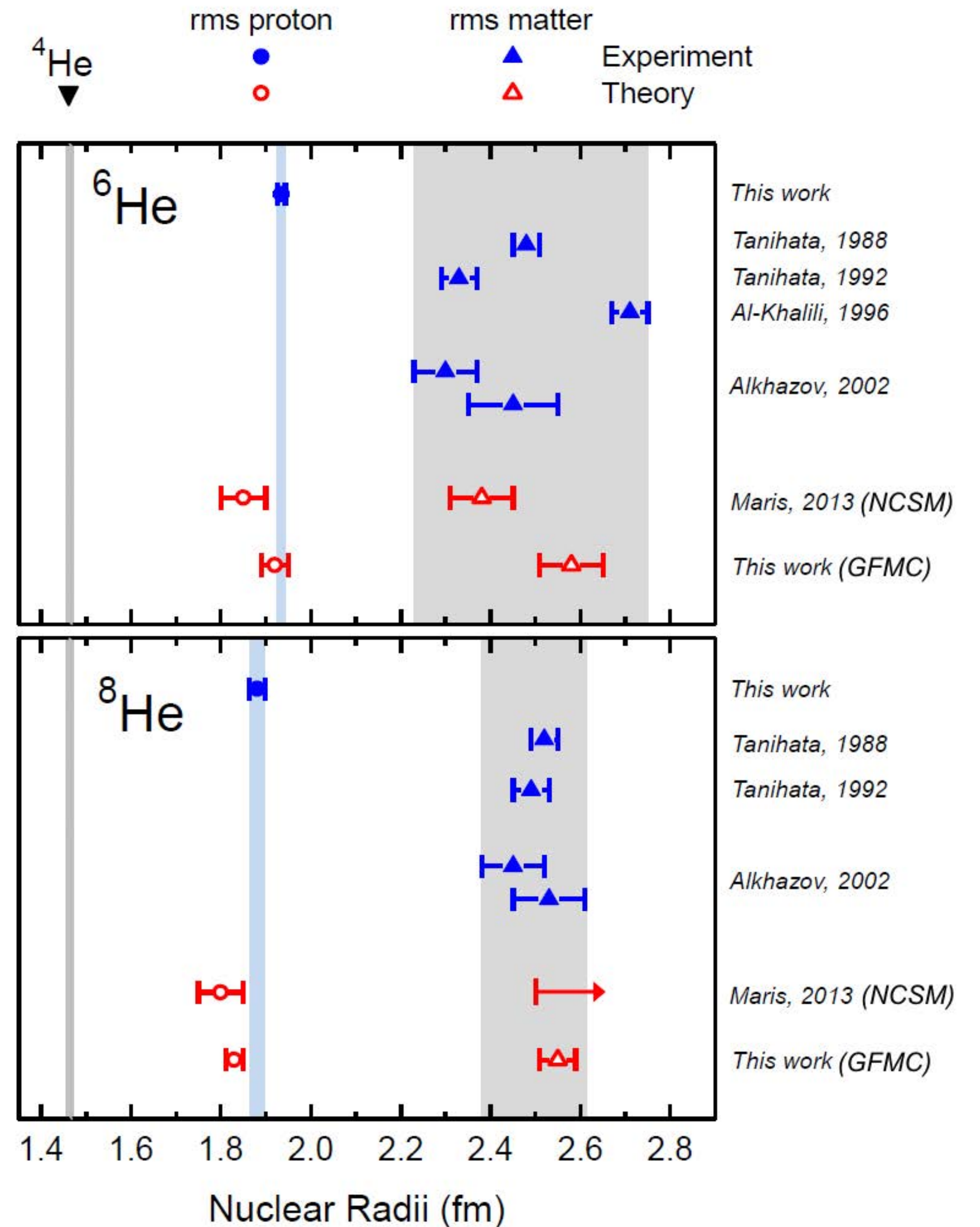
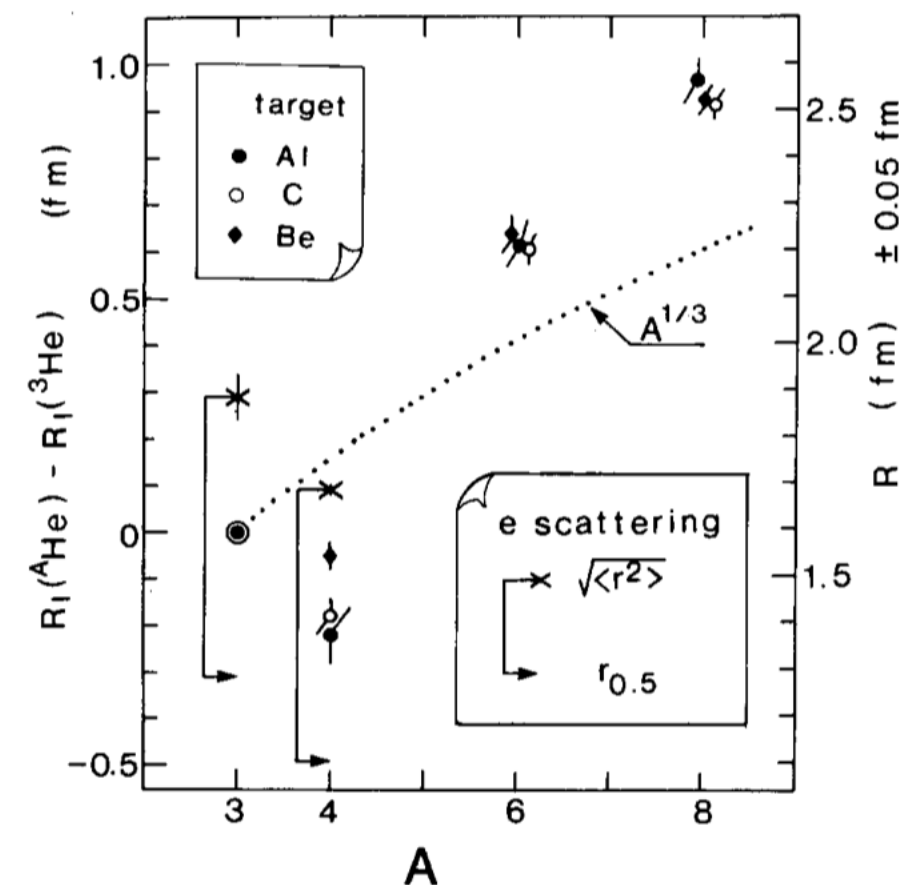
Table 1

Interaction cross sections (σ_I) of He isotopes (σ_I in mb). The listed errors include statistical and systematic errors. The largest systematic errors were due to uncertainties in the estimation of scattering-out probabilities of non-interacting nuclei.

Target	Beam			
	^3He	^4He	^6He	^8He
Be	498 ± 4	485 ± 4	672 ± 7	757 ± 4
C	550 ± 5	503 ± 5	722 ± 6	817 ± 6
Al	850 ± 9	527 ± 26^a	1063 ± 8	1197 ± 9
^4He		262 ± 19^a		

^a) Data from ref. [4]. The present value of $^4\text{He} + \text{C}$ cross section agrees with the known value within quoted errors.

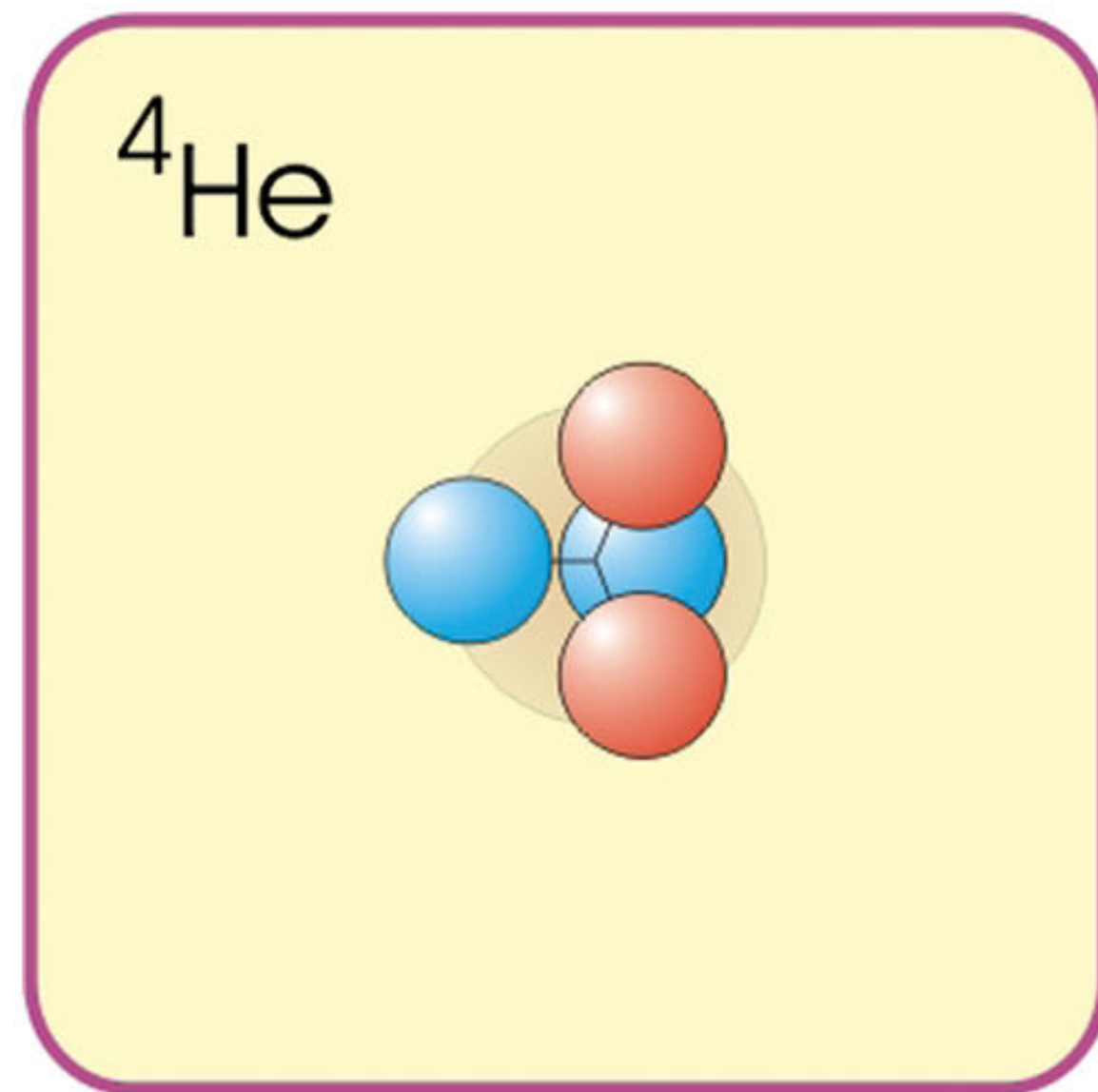
sections on the target thickness was observed. The interaction cross sections σ_I thus determined are



He isotopes

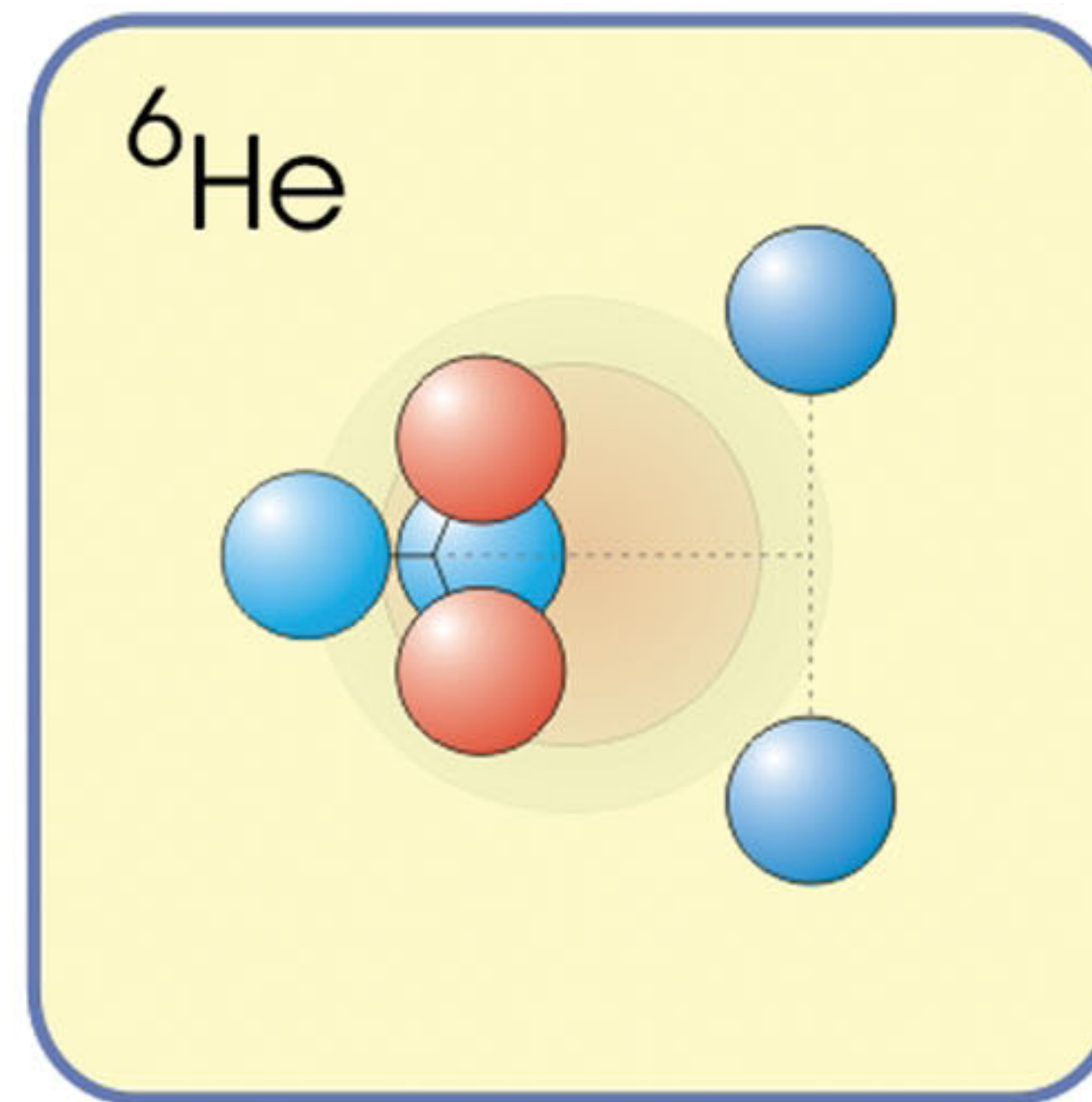
Which is larger charge radius? ${}^4\text{He}$, ${}^6\text{He}$, ${}^8\text{He}$?

Which is larger matter radius? ${}^4\text{He}$, ${}^6\text{He}$, ${}^8\text{He}$?



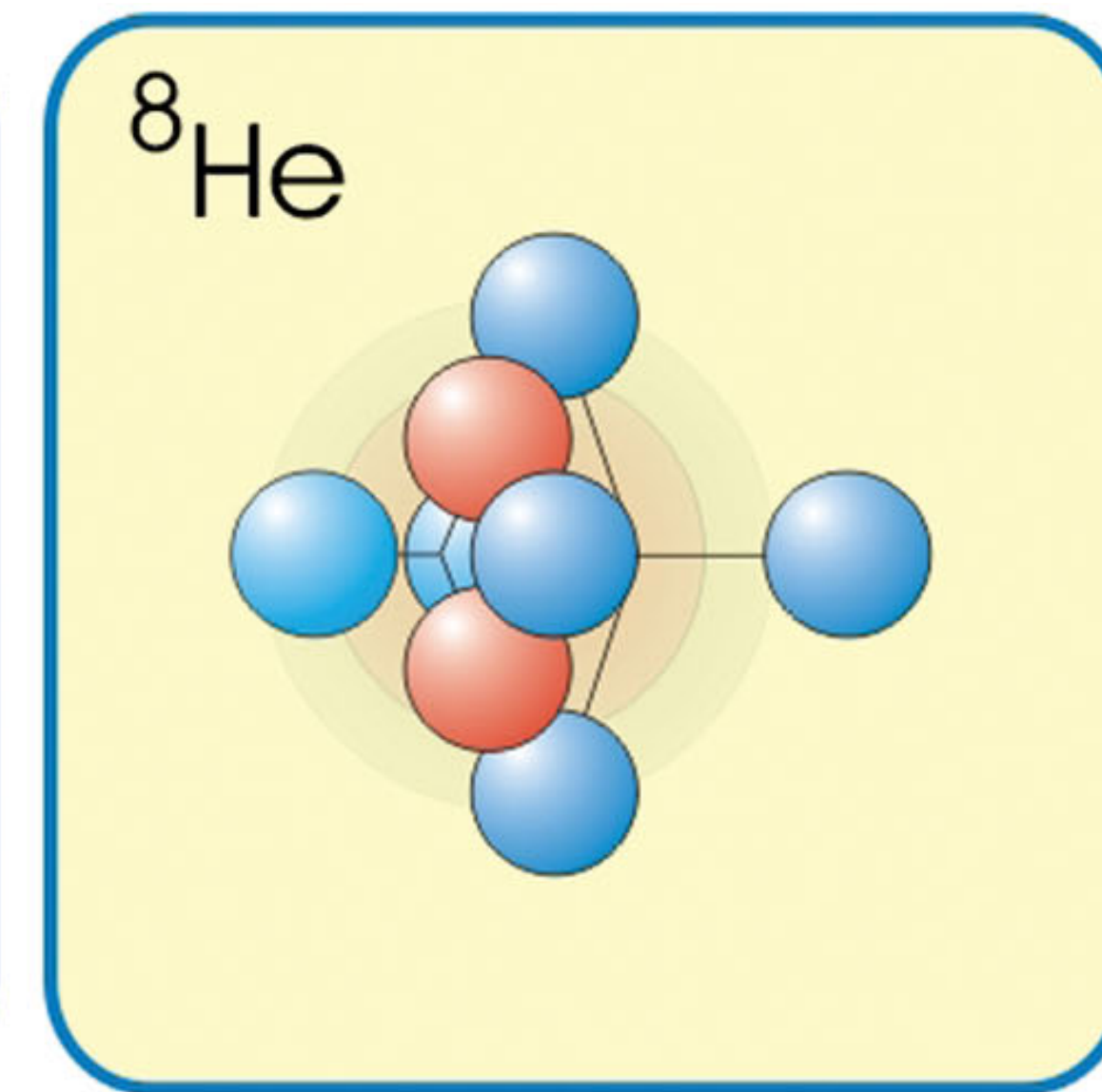
1.681(4) fm

1.63(3) fm



2.060(8) fm

2.33(4) fm

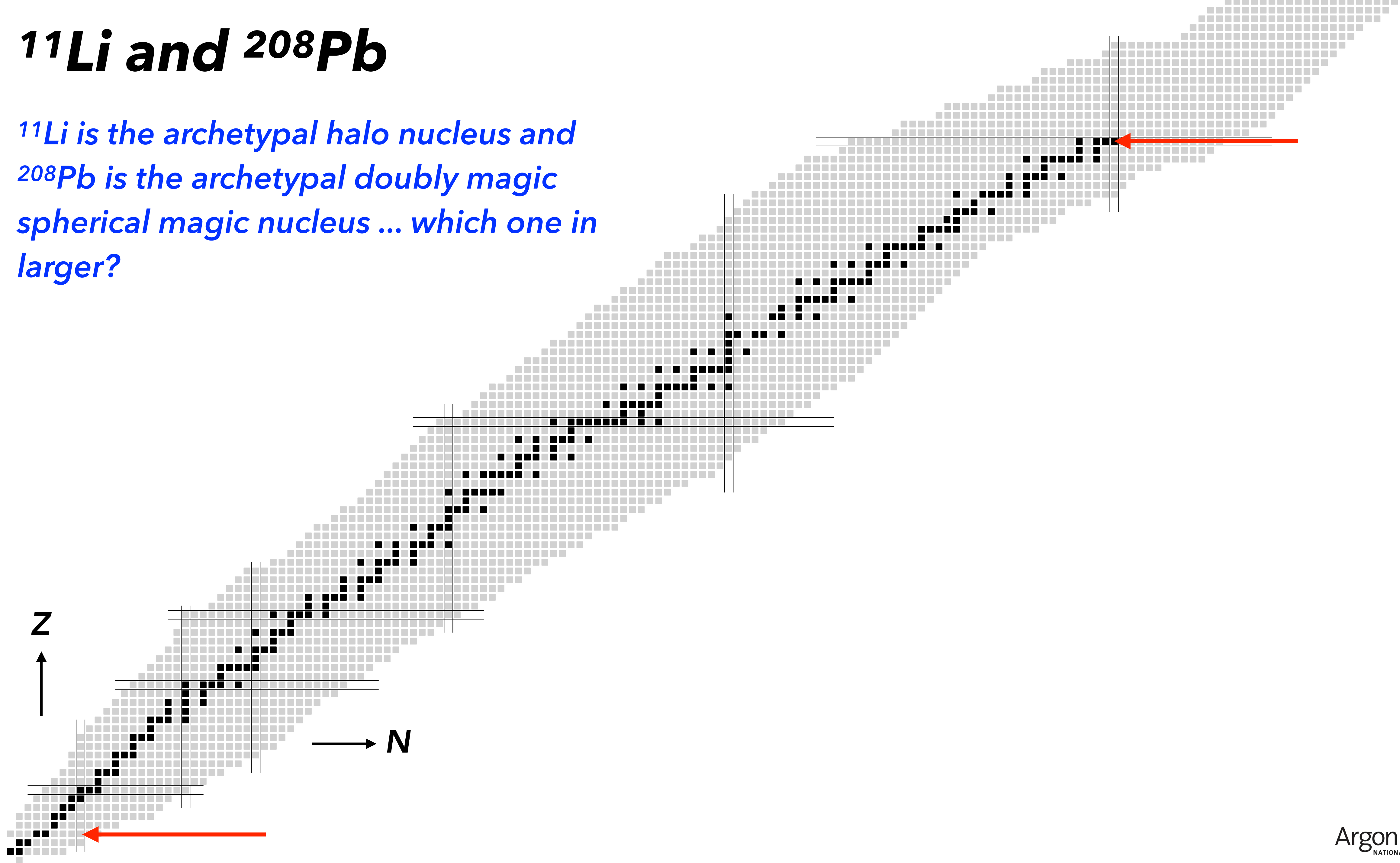


1.959(16) fm

2.49(4) fm

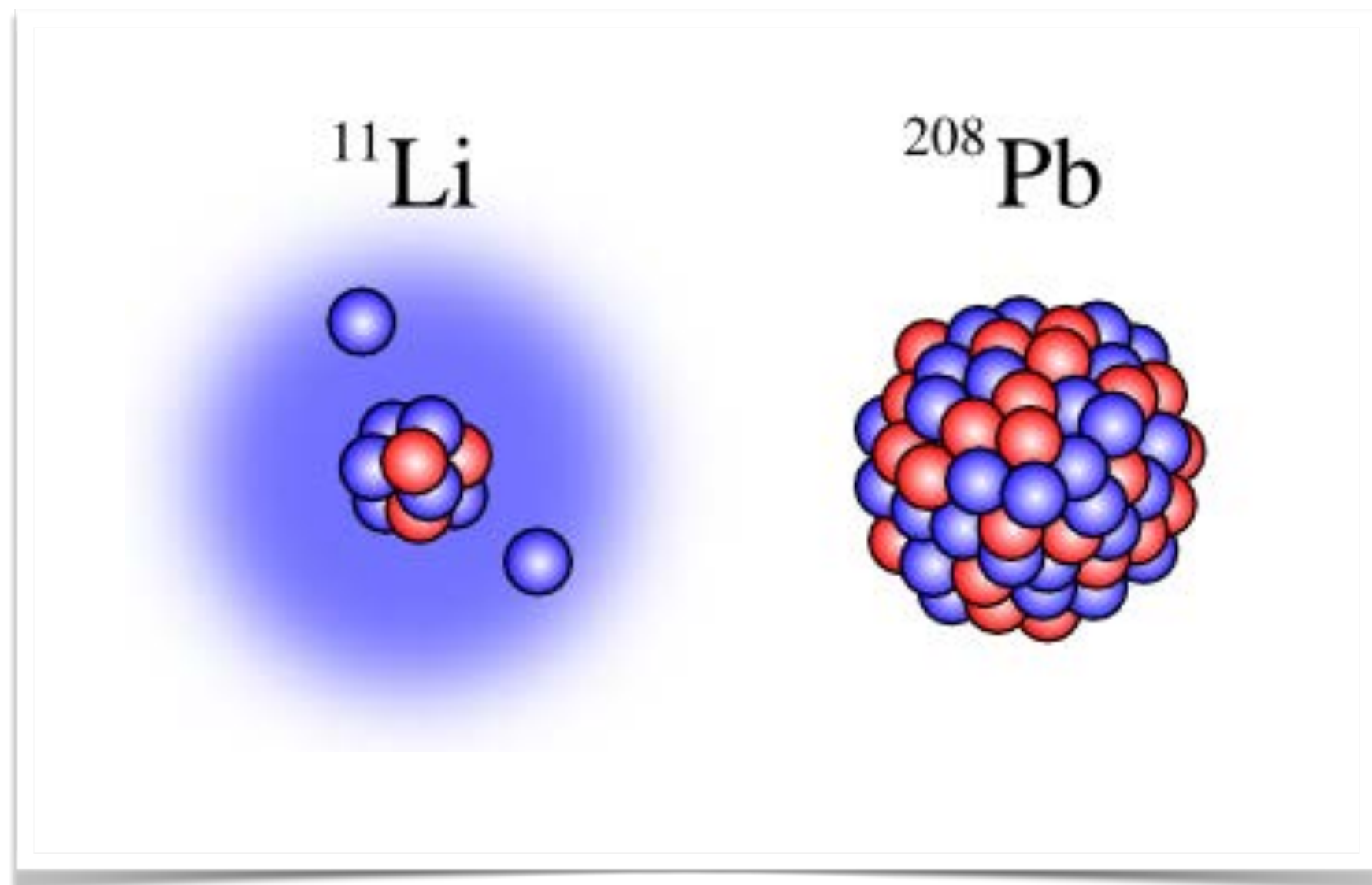
^{11}Li and ^{208}Pb

^{11}Li is the archetypal halo nucleus and ^{208}Pb is the archetypal doubly magic spherical magic nucleus ... which one is larger?



^{11}Li and ^{208}Pb

^{11}Li is the archetypal halo nucleus and ^{208}Pb is the archetypal doubly magic spherical magic nucleus ... which one is larger? About the same size!



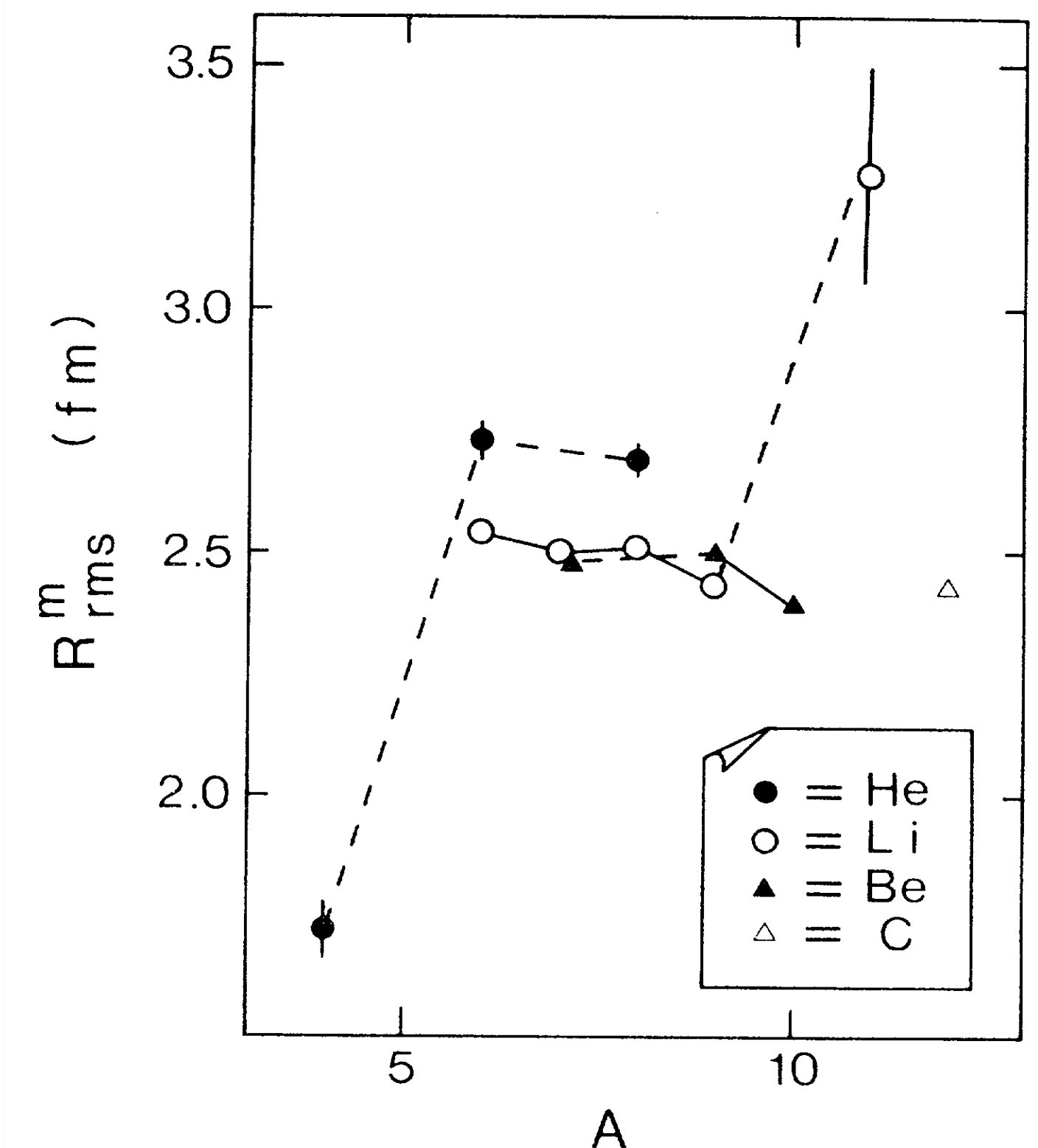
Recently, exotic-isotope beams, produced through the projectile-fragmentation process in high-energy heavy-ion reactions, were used to measure the interaction cross sections (σ_I) for all the known He isotopes.¹ This novel technique of using exotic nuclear beams makes it possible to study systematically properties of unstable nuclei. In the present paper, we report the σ_I for all the known Li isotopes (^6Li , ^7Li , ^8Li , ^9Li , and ^{11}Li) and ^7Be , ^9Be , and ^{10}Be on the target nuclei Be, C, and Al at 790 MeV/nucleon. A firm basis has been empirically established by use of a Glauber-type calculation to extract root mean square (rms) nuclear radii from the σ_I .

The Li isotopes, except ^{11}Li , and the Be isotopes were produced as secondary beams through projectile fragmentation of the 800-MeV/nucleon ^{11}B accelerated by the Bevalac at the Lawrence Berkeley Laboratory. The beam of ^{11}Li was produced from a ^{20}Ne primary beam. The isotopes produced in a production target of Be were separated by rigidity with the beam-line magnet system as described in previous papers.^{1,2} The rigidity-separated isotopes were further identified before incidence on a reaction target by velocity [time-of-flight (TOF)] and by charge (pulse height in scintillation counters). No contamination more than 10^{-3} was observed in any selected isotope beam.

The interaction cross section (σ_I) was measured by a transmission experiment using the large-acceptance spectrometer as in the measurement of the He isotopes.¹ Here σ_I is defined as the total reaction cross section for the change of proton and/or neutron number in the incident nucleus. The obtained σ_I are listed in Table I. The largest systematic error on σ_I , up to about 0.3%, came from uncertainties in the estimation of the scattering-out probability of the nonin-

TABLE I. Interaction cross sections (σ_I) in millibarns.

Beam	Be	Target C	Al
^6Li	651 ± 6	688 ± 10	1010 ± 11
^7Li	686 ± 4	736 ± 6	1071 ± 7
^8Li	727 ± 6	768 ± 9	1147 ± 14
^9Li	739 ± 5	796 ± 6	1135 ± 7
^{11}Li		1040 ± 60	
^7Be	682 ± 6	738 ± 9	1050 ± 17
^9Be	755 ± 6	806 ± 9	1174 ± 11
^{10}Be	755 ± 7	813 ± 10	1153 ± 16



I. Tanihata et al., Phys. Rev. Lett. **54**, 2676 (1985)

(Not sure who to credit for figure ... found as image online ...)

Nuclei ... neutron stars

Neutron rich matter in heaven and on Earth

J. Piekarewicz^{1,*} and F. J. Fattoyev^{2,†}

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(Dated: July 8, 2019)

Despite a length-scale difference of 18 orders of magnitude, the internal structure of neutron stars and the spatial distribution of neutrons in atomic nuclei are profoundly connected.

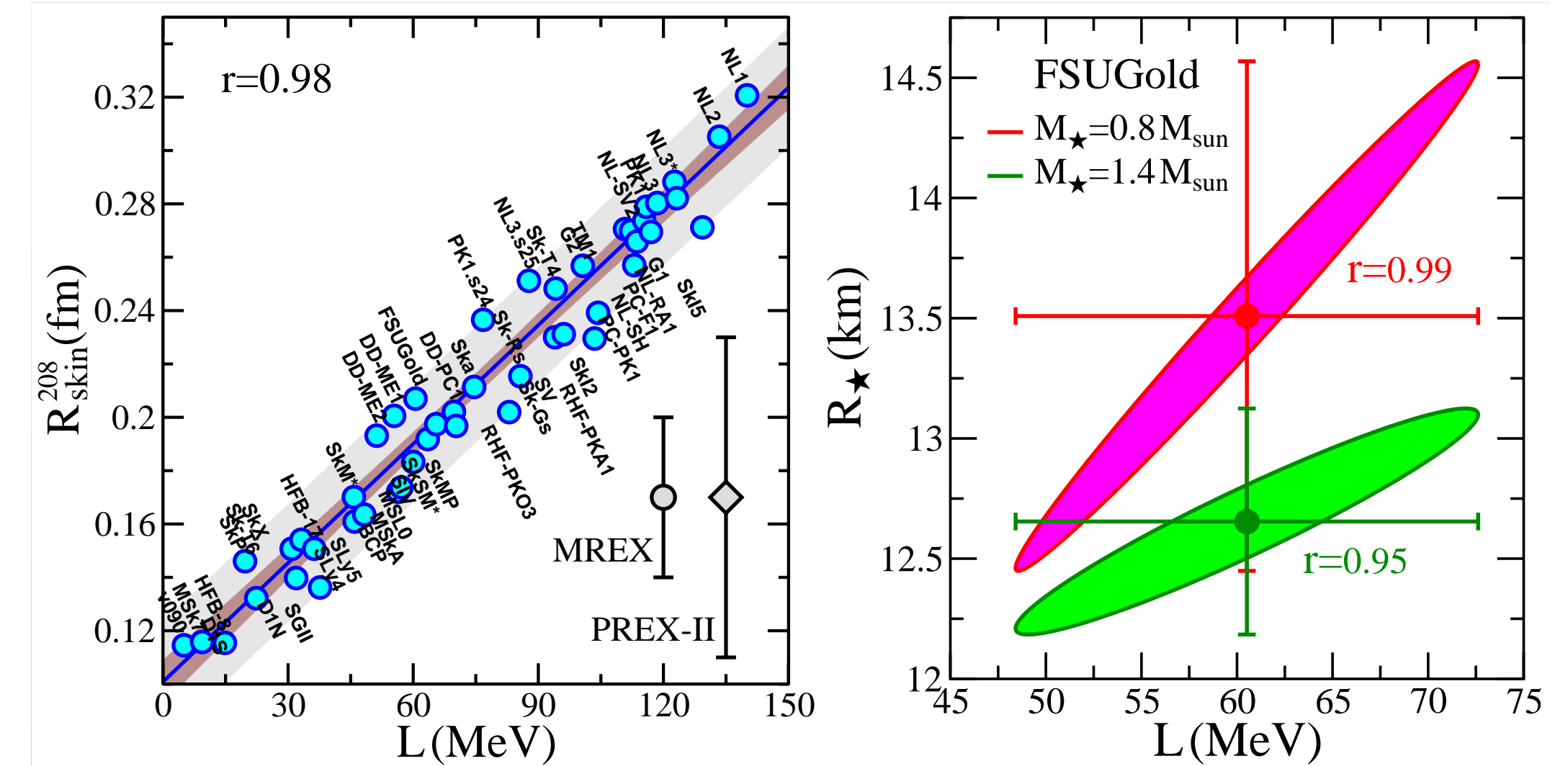
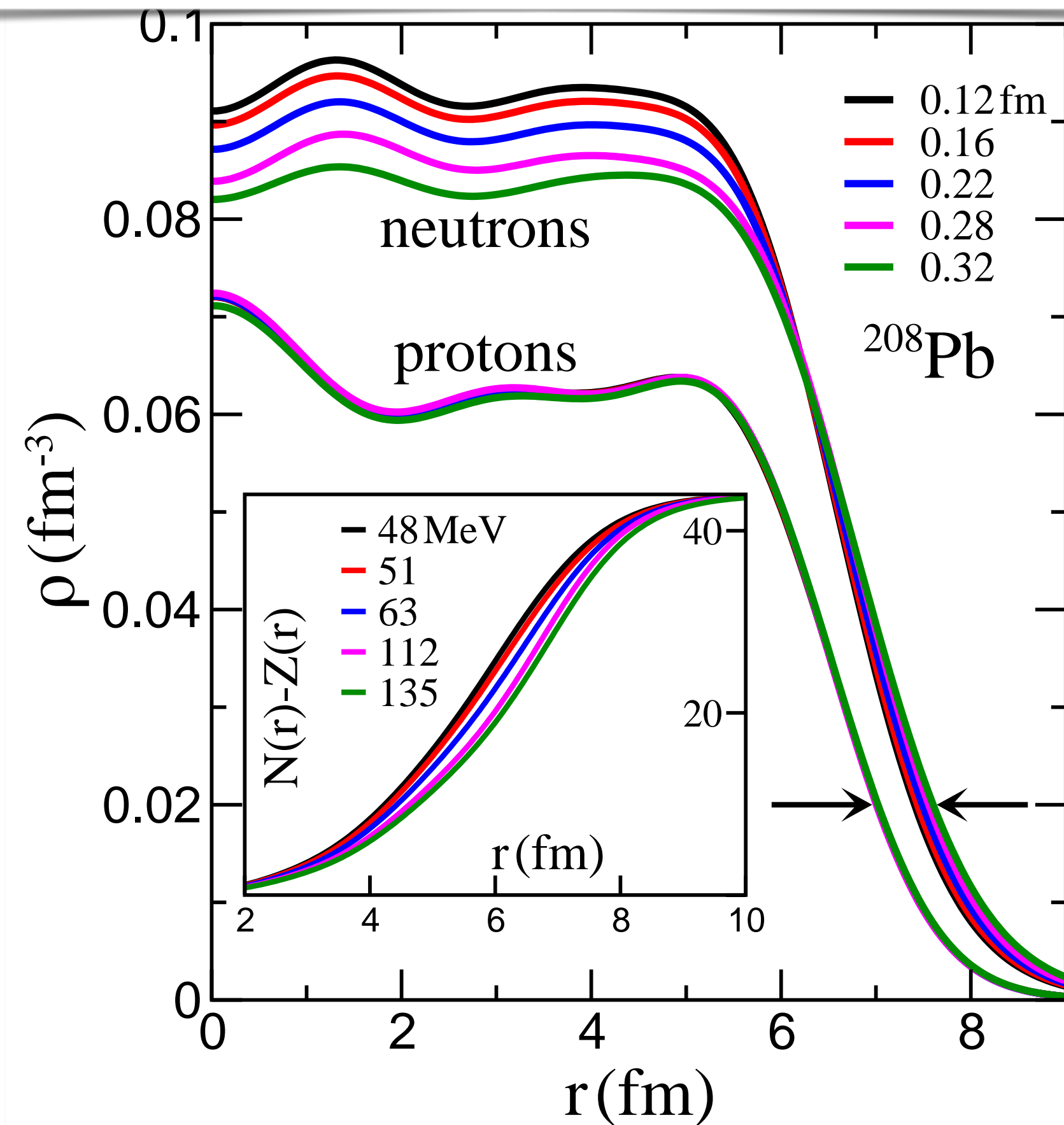


FIG. 3: **Connecting the very small to the very big.** Despite a difference in size of 18 orders of magnitude, the symmetry pressure L controls both the neutron skin thickness of ^{208}Pb as well as the radius of a neutron star. On the left hand panel a large set of highly successful models are used to illustrate the correlation between L and R_{skin}^{208} ; figure adapted from Ref. [5]. The right hand panel displays the correlation between L and neutron star radii for one of these models: “FSUGold” [6].

Nuclei ... neutron stars

Neutron rich matter in heaven and on Earth

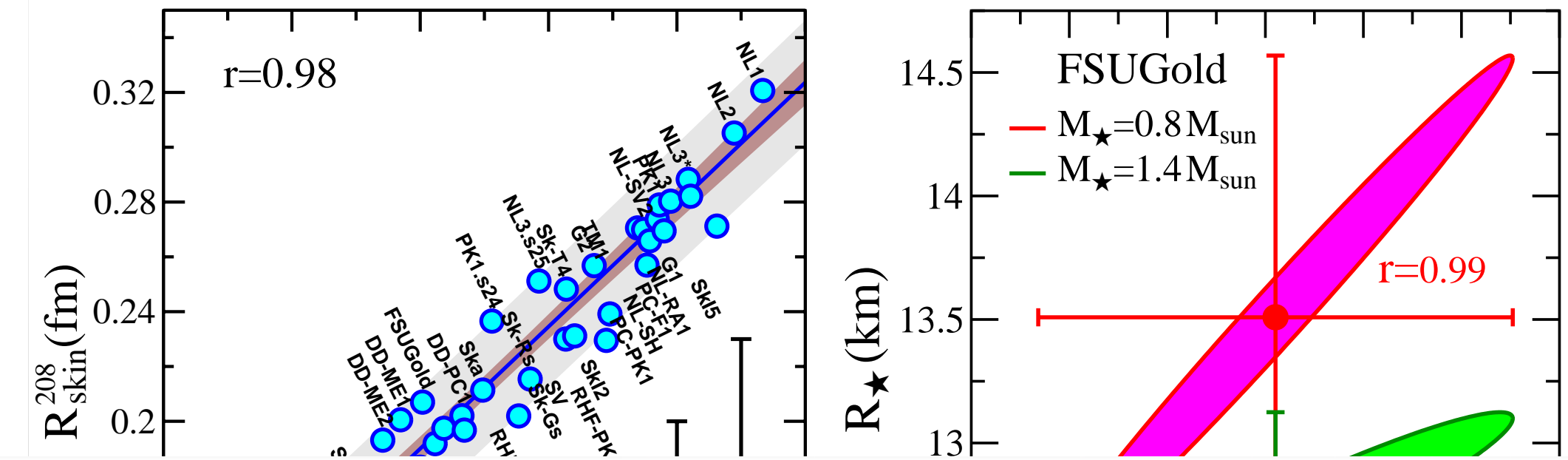
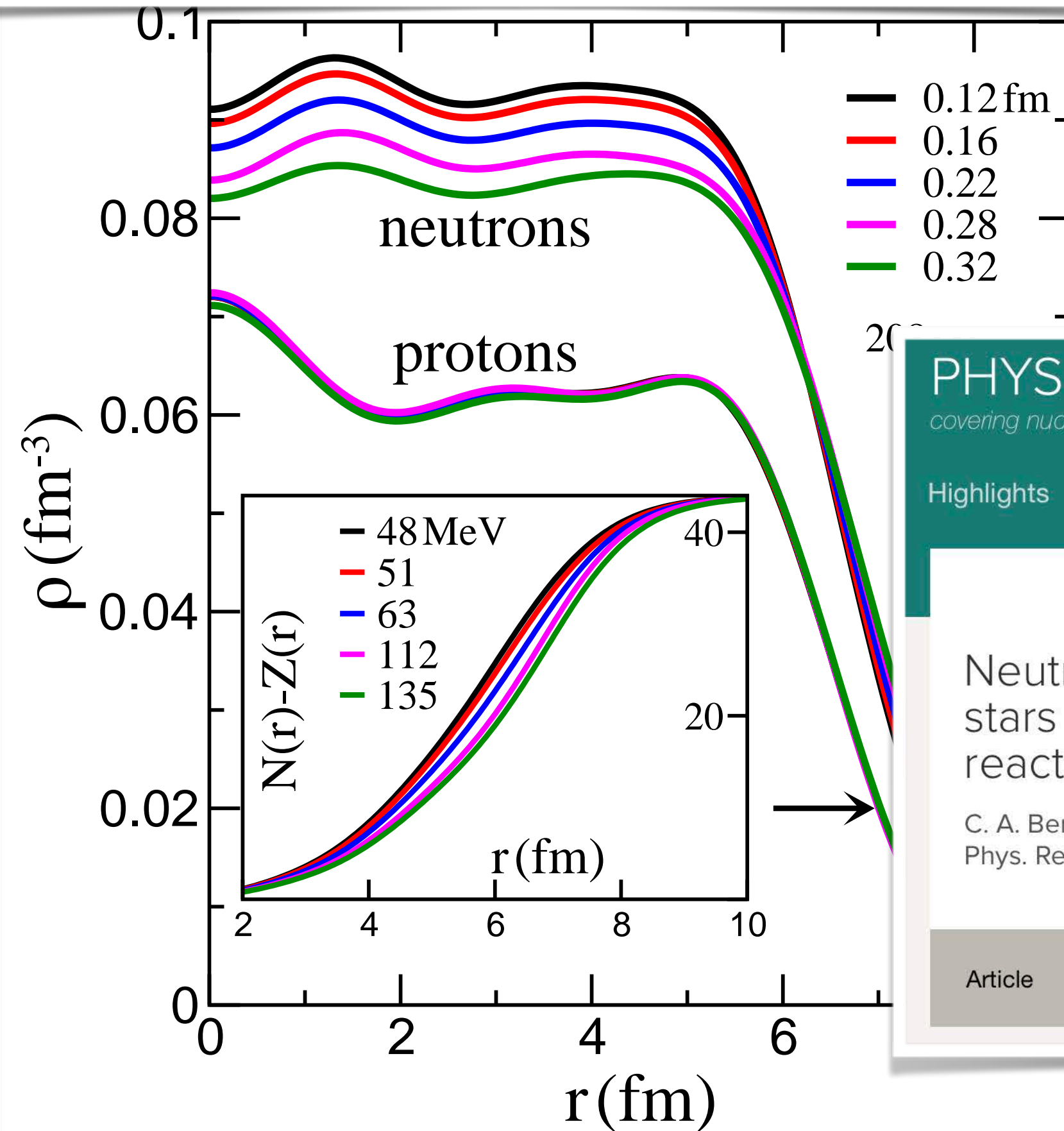
J. Piekarewicz^{1,*} and F. J. Fattoyev^{2,†}

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Despite a length-scale difference of 18 orders of magnitude, the internal structure of neutron stars and the spatial distribution of neutrons in atomic nuclei are profoundly connected.



PHYSICAL REVIEW C

covering nuclear physics

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Neutron skins as laboratory constraints on properties of neutron stars and on what we can learn from heavy ion fragmentation reactions

C. A. Bertulani and J. Valencia
Phys. Rev. C **100**, 015802 – Published 15 July 2019



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symmetry and panel m Ref. [5]. [6].

Shapes (and sizes)

ISOLDE and Hg beams

nature
physics

LETTERS

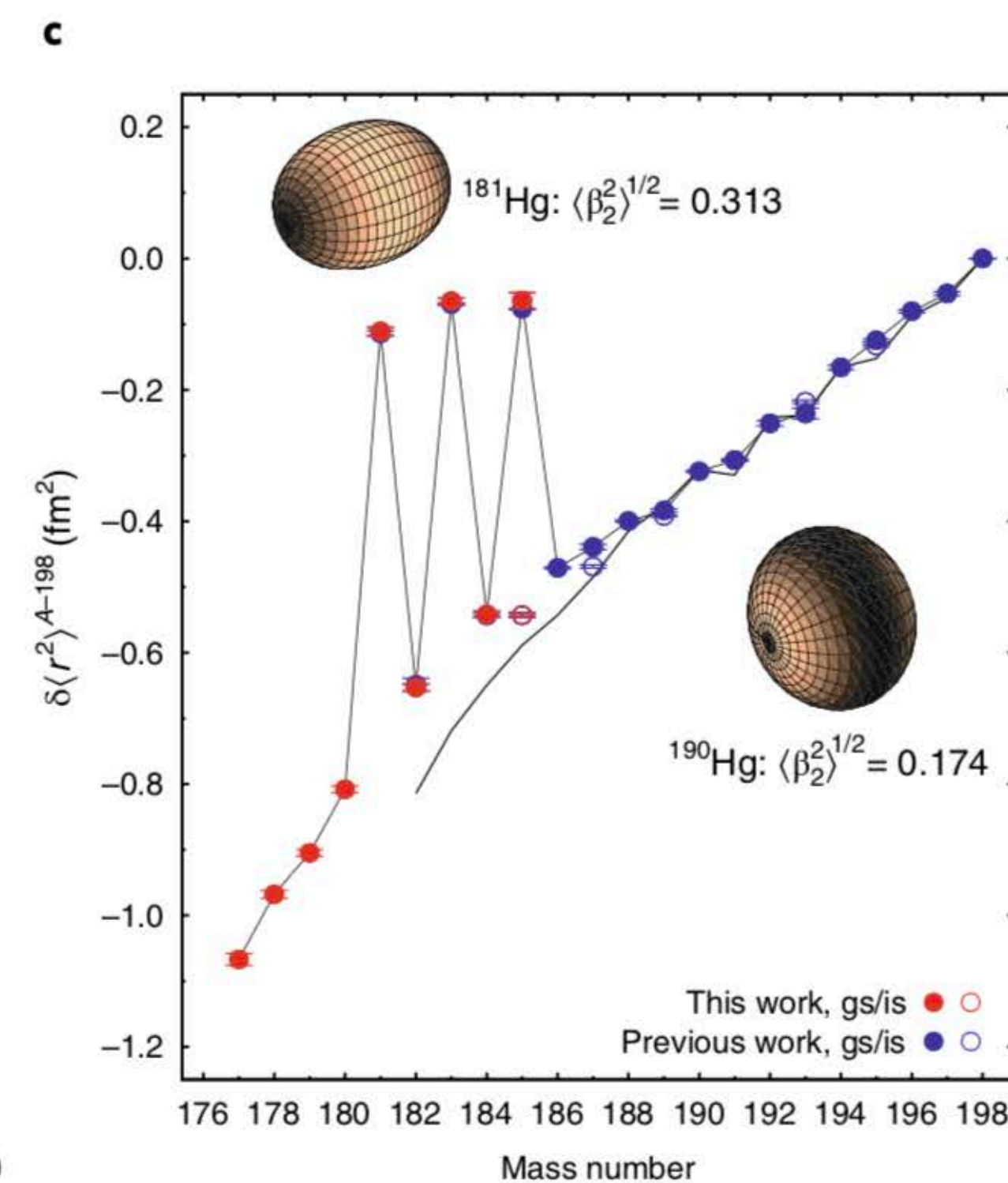
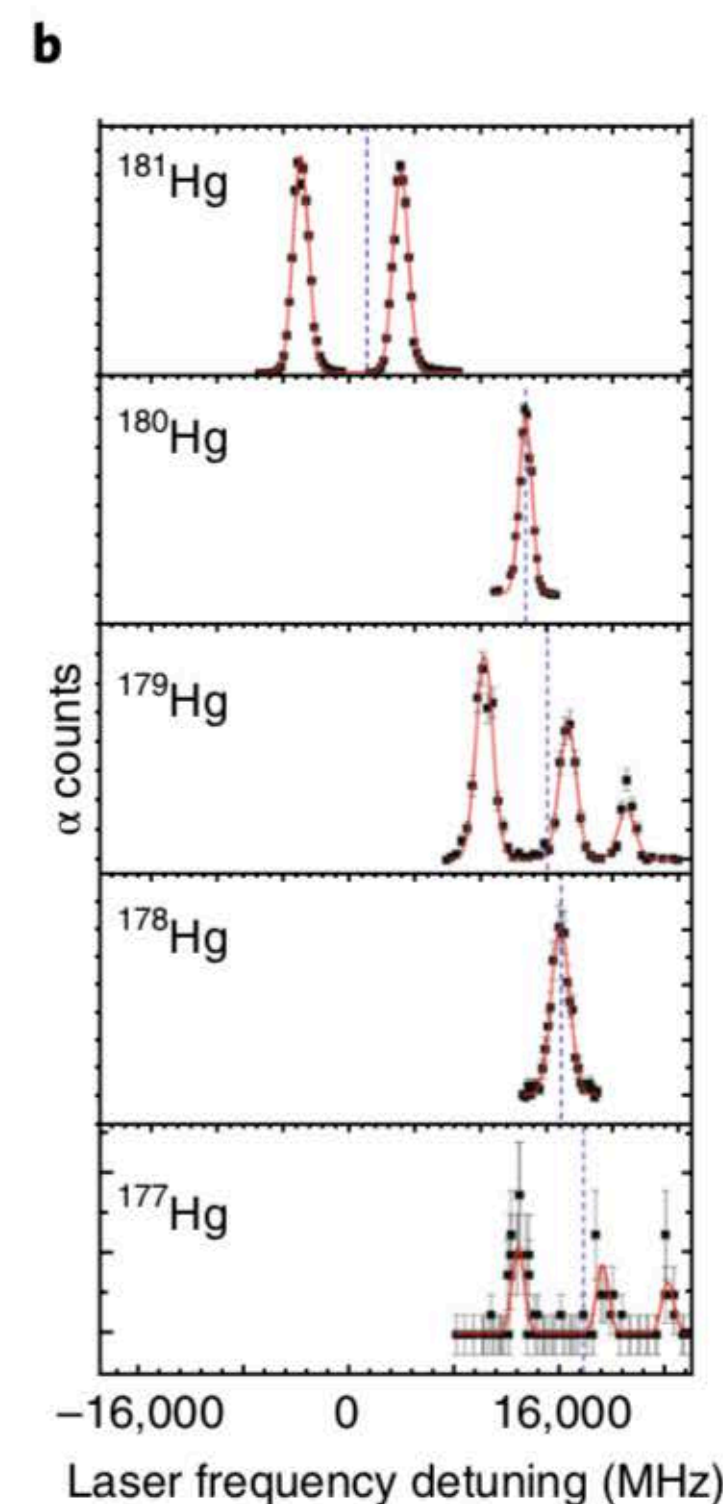
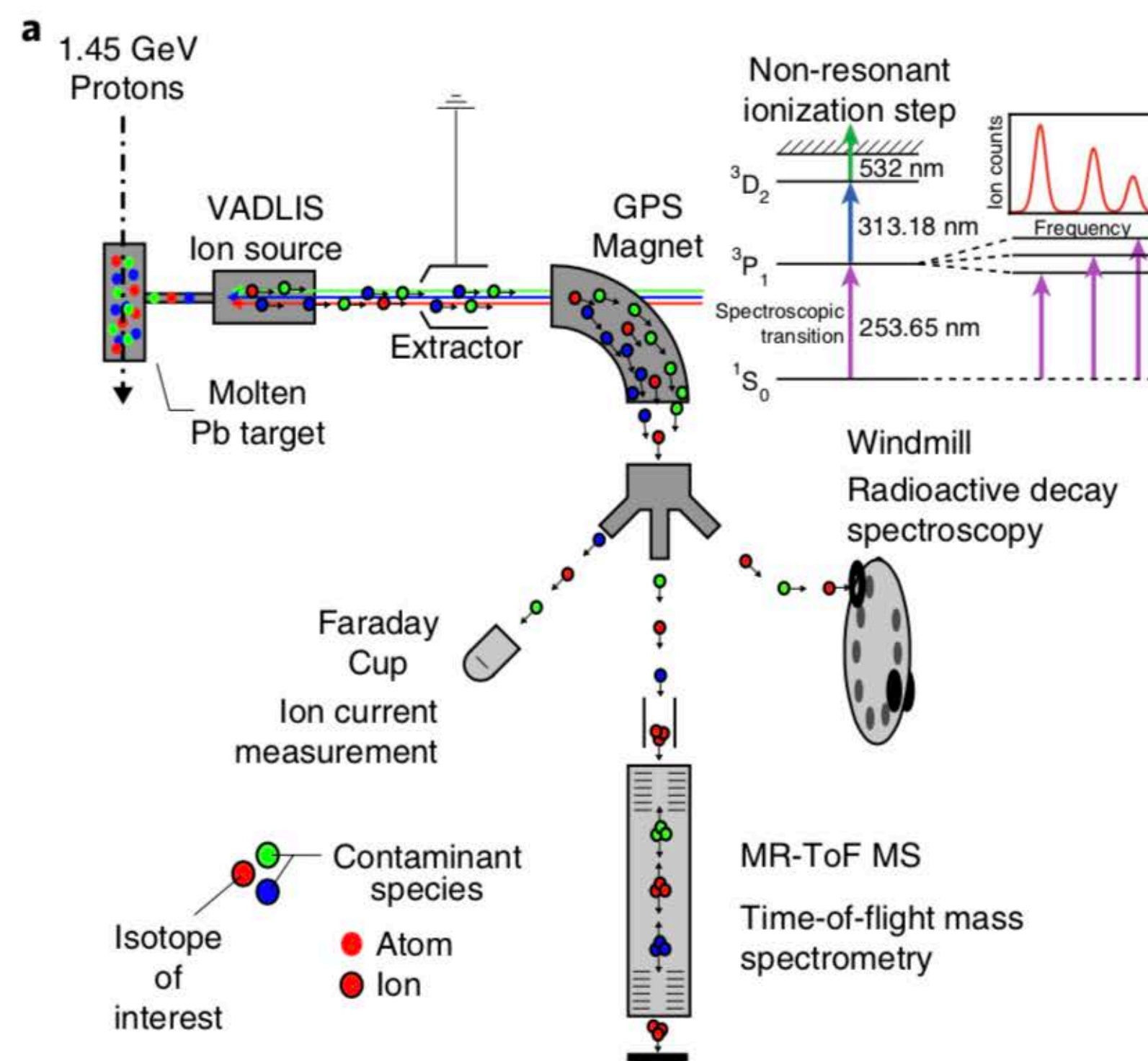
<https://doi.org/10.1038/s41567-018-0292-8>

Characterization of the shape-staggering effect in mercury nuclei

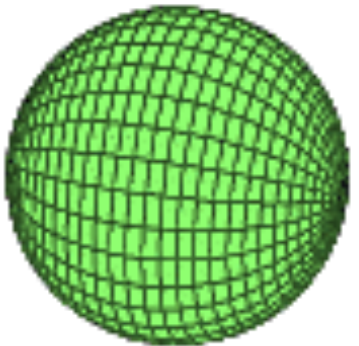
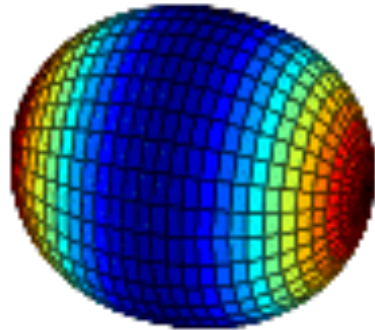
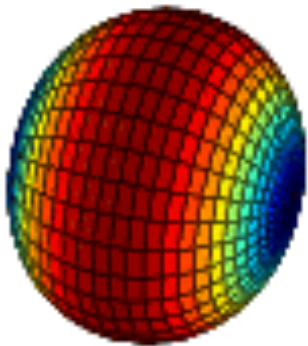
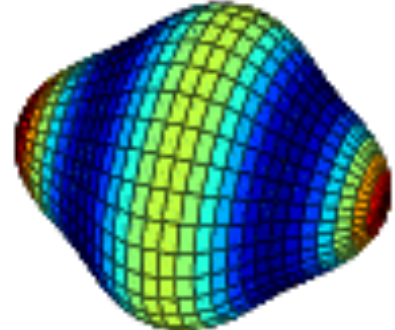
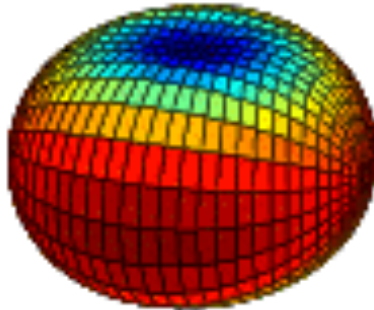
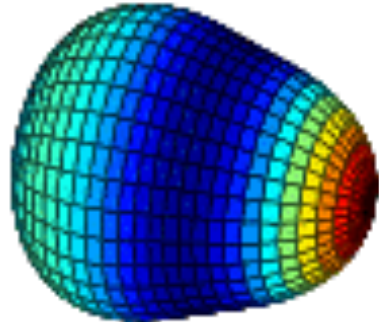
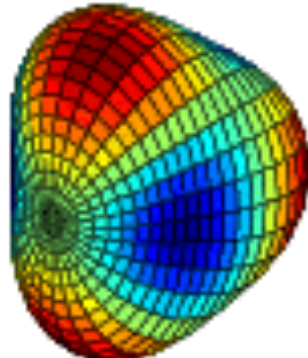
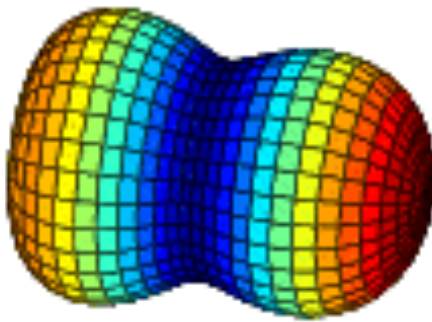
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J. Dobaczewski⁶, G. J. Farooq-Smi³,
L. Ghys³, M. Huyse³, S. Kreim⁸, D.
T. Otsuka^{3,4,12,13,14}, A. Pastore⁶, M.
P. Spagnoletti¹⁰, C. Van Beveren³,
F. Wienholtz¹⁵, R. N. Wolf⁸, A. Zaccaro¹

LETTERS

NATURE PHYSICS



Shapes

$\beta_{\lambda\mu} = 0$	$\beta_{20} > 0$	$\beta_{20} < 0$	$\beta_{40} > 0$
			
$\beta_{22} \neq 0$	$\beta_{30} \neq 0$	$\beta_{32} \neq 0$	$\beta_{20} \gg 0$
			

$$R(\theta, \phi) = c(\alpha_{\lambda\mu})R_0 \left[1 + \sum_{\lambda=0}^{\infty} \sum_{\mu=-\lambda}^{\lambda} \alpha_{\lambda\mu} Y_{\lambda\mu}(\theta, \phi) \right]$$

Multipole order: 2^λ

2^0 = monopole - breathing mode

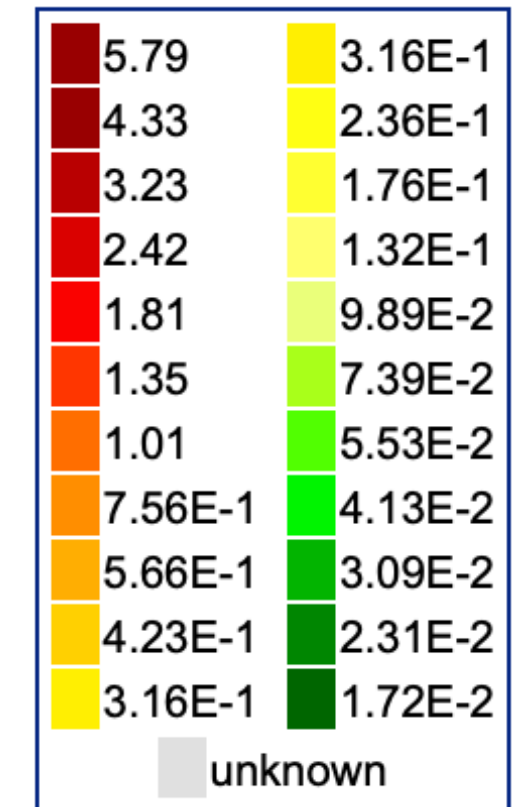
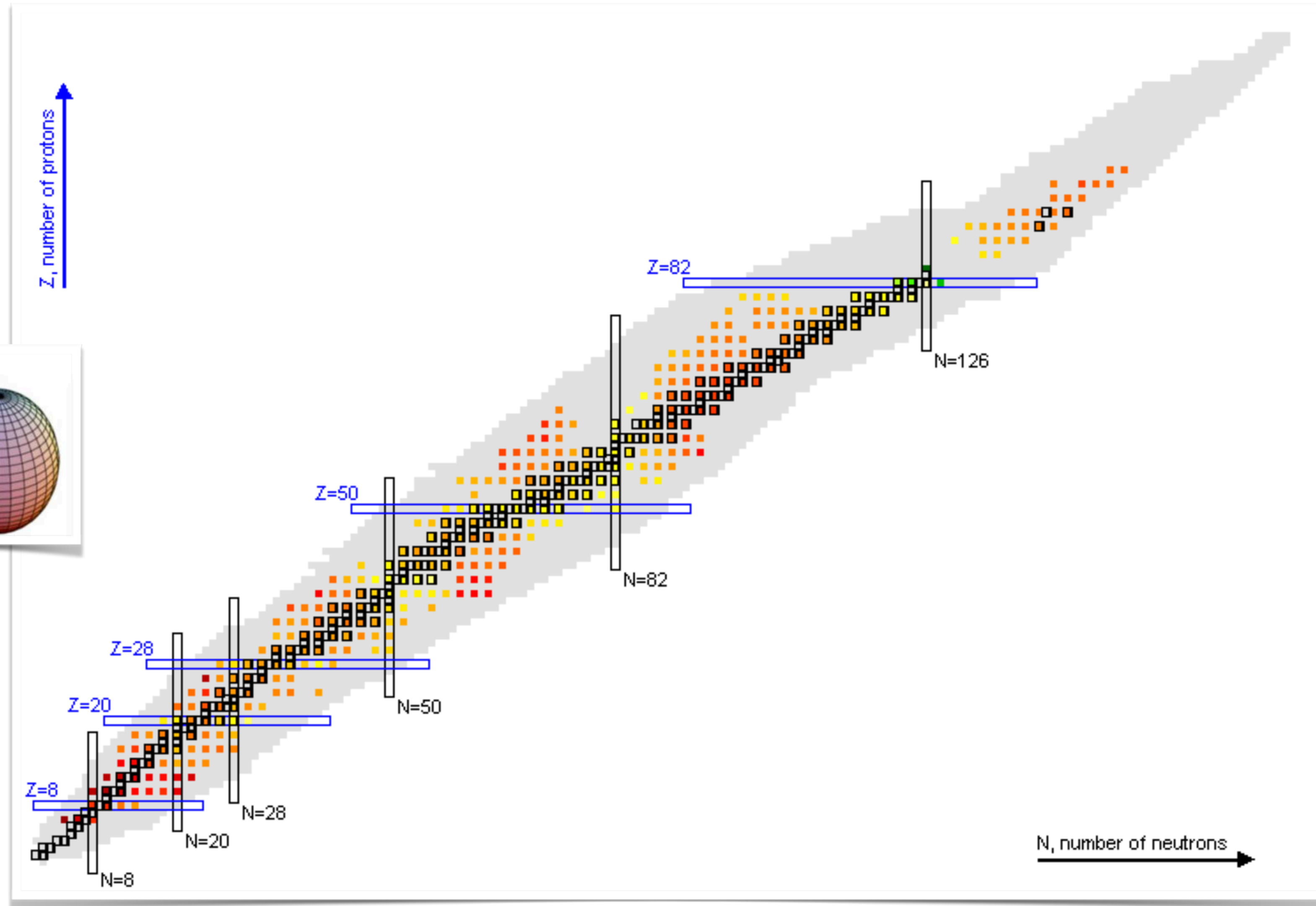
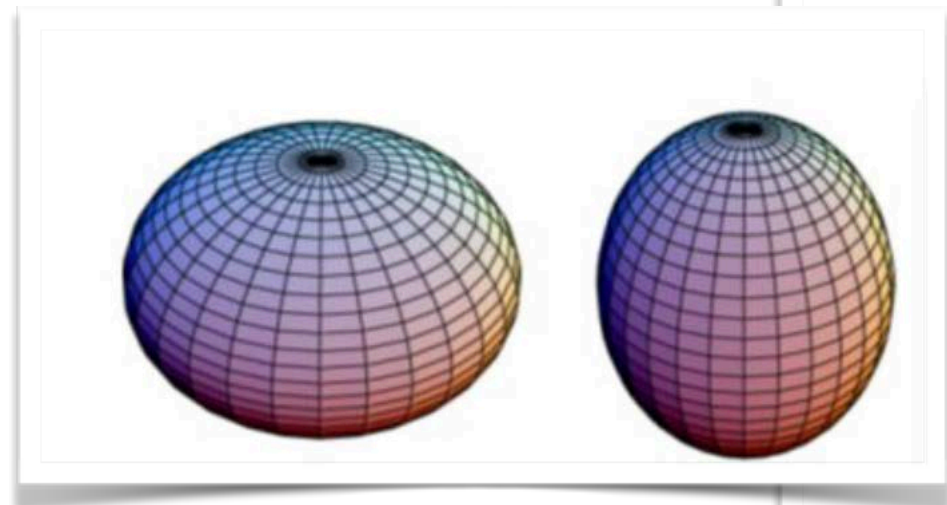
2^1 = dipole - centre of mass shift

2^2 = quadrupole - axial deformation

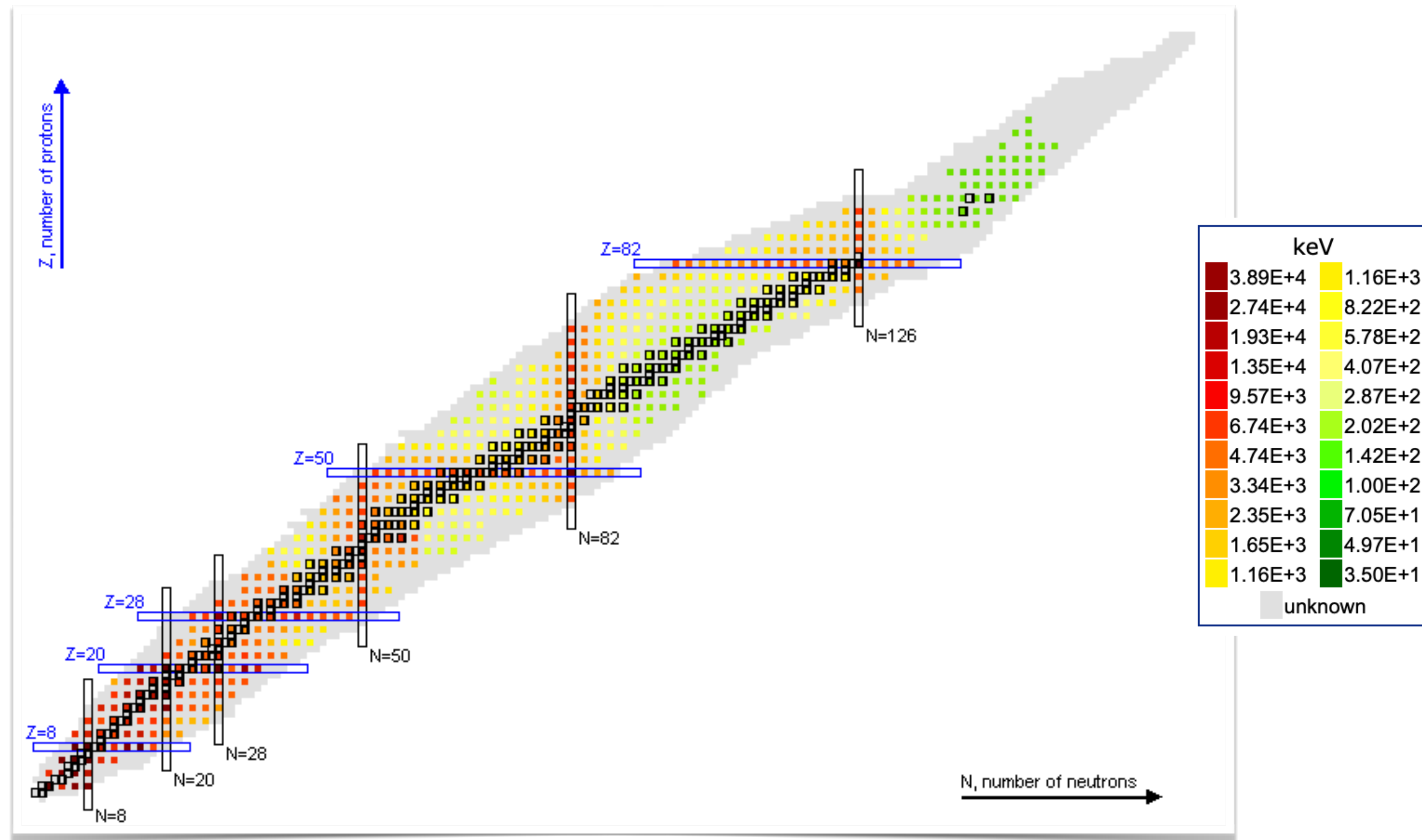
2^3 = octupole - asymmetric deformation

2^4 = hexadecapole - pinching

Shapes, spectroscopy



Shapes, spectroscopy



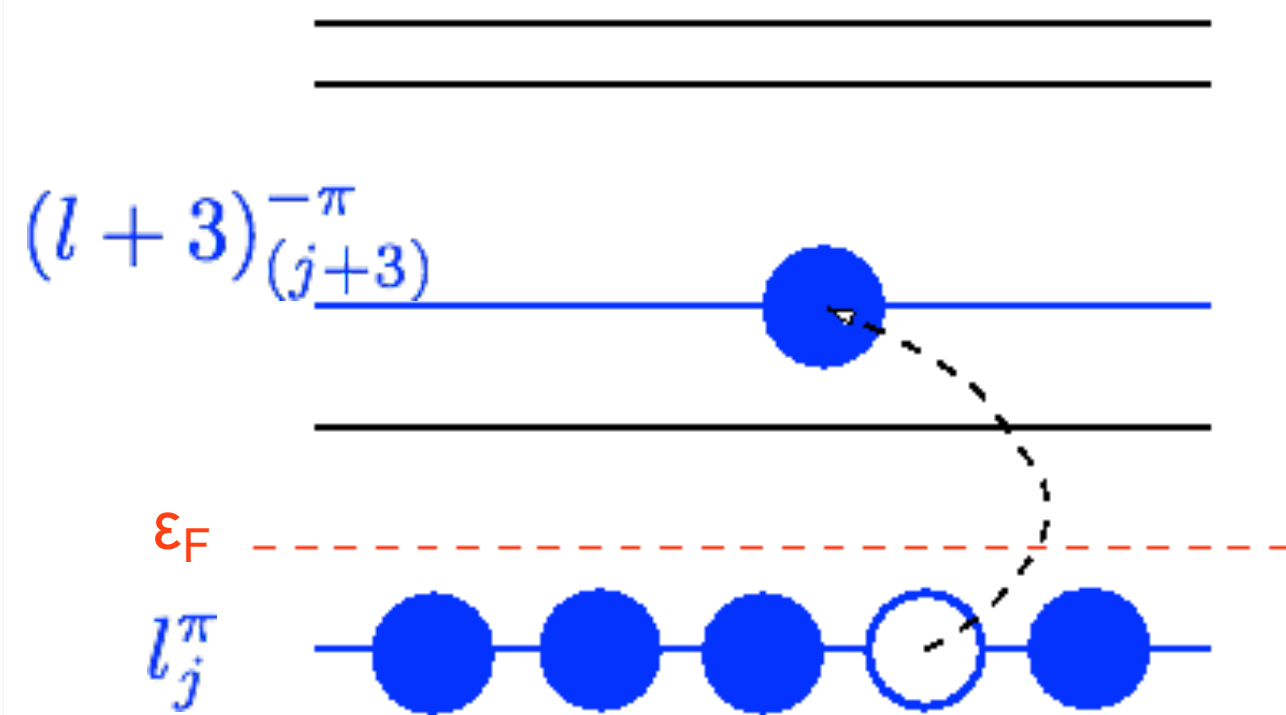
<https://www.nndc.bnl.gov/nudat2/>, energy of the first excited $2+$ state

Strange shapes ... "pears"

Reflection Asymmetry

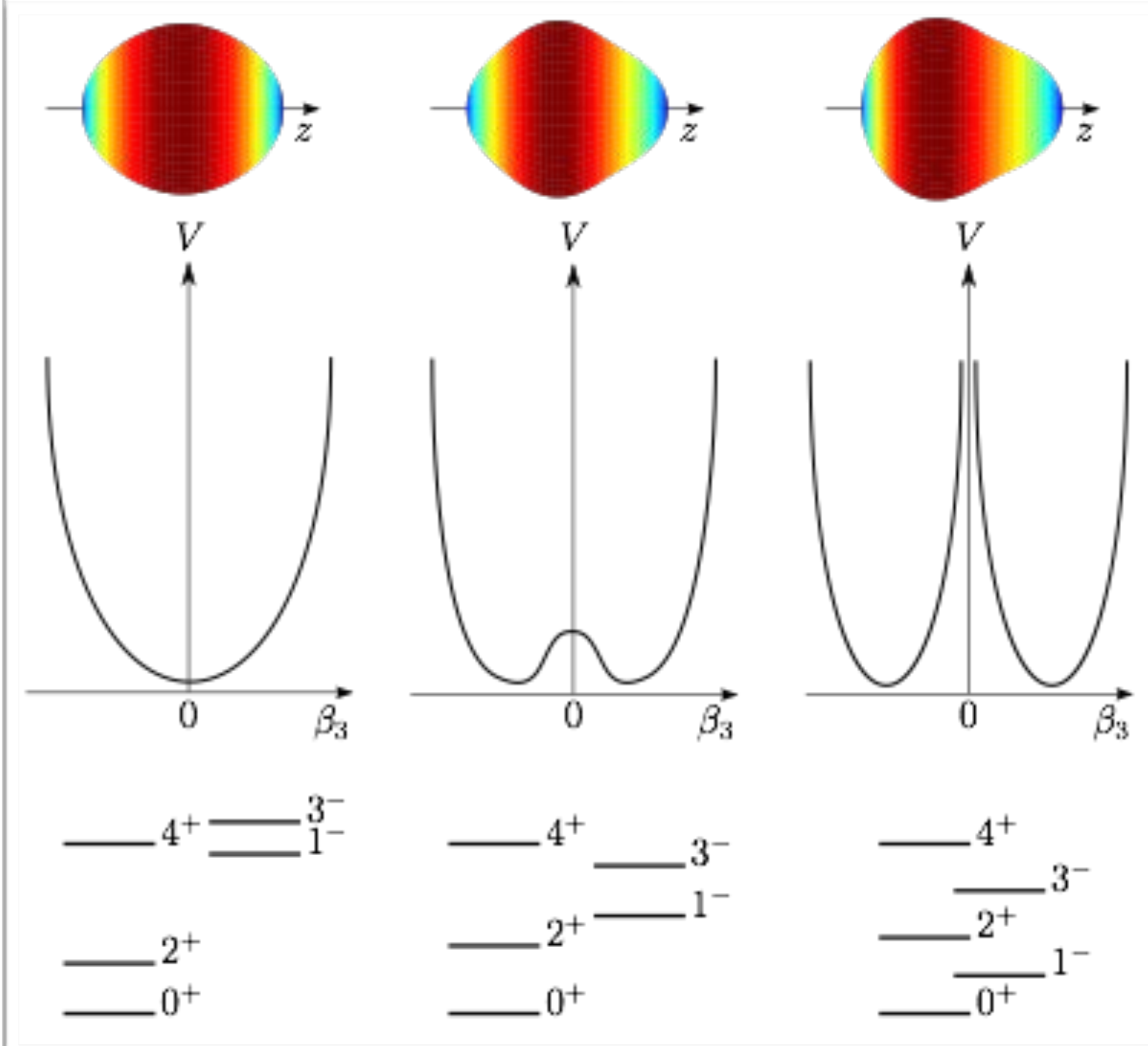
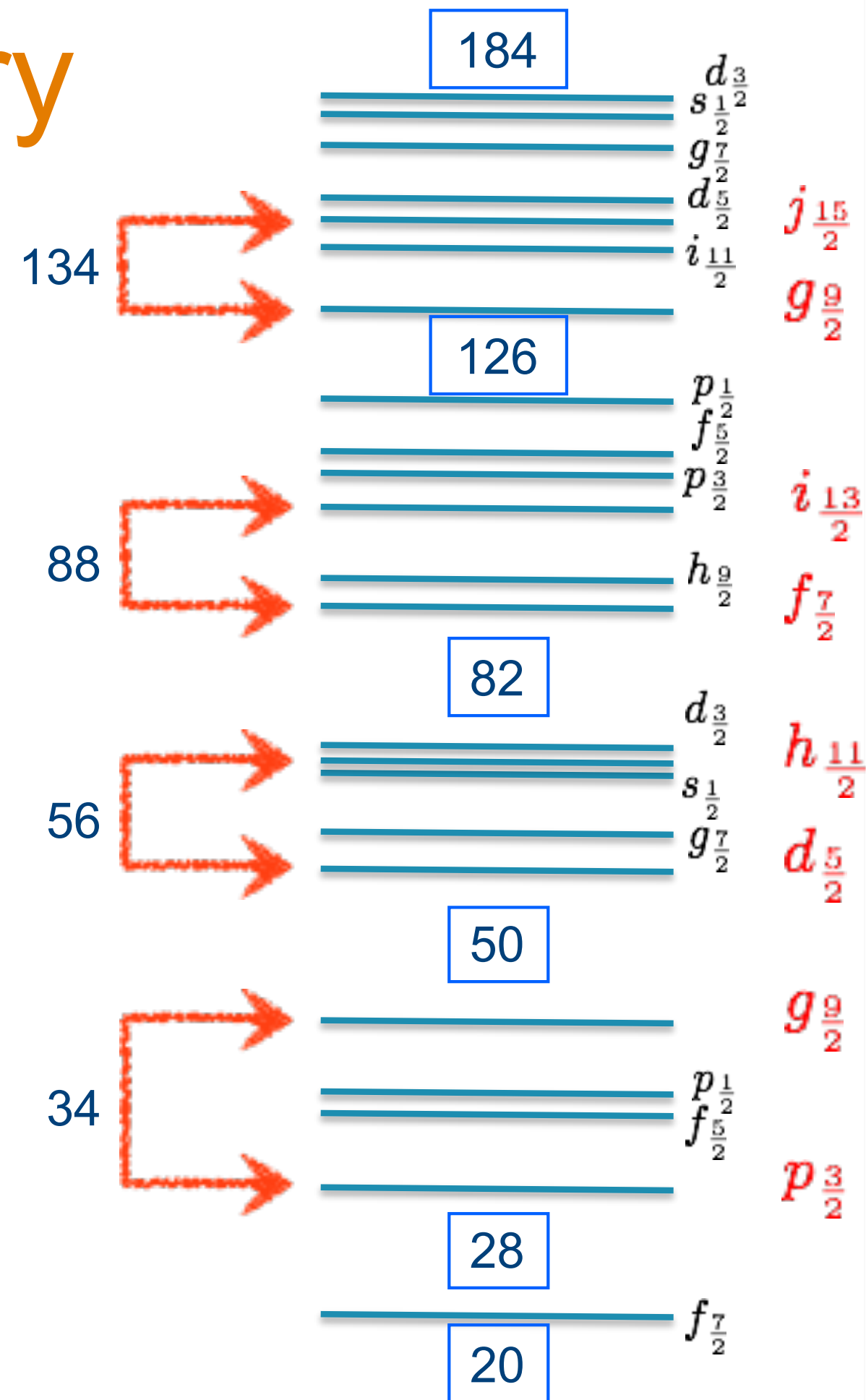
Microscopically driven...

Intruder orbitals of opposite parity and $\Delta J, \Delta L = 3$ close to the Fermi level

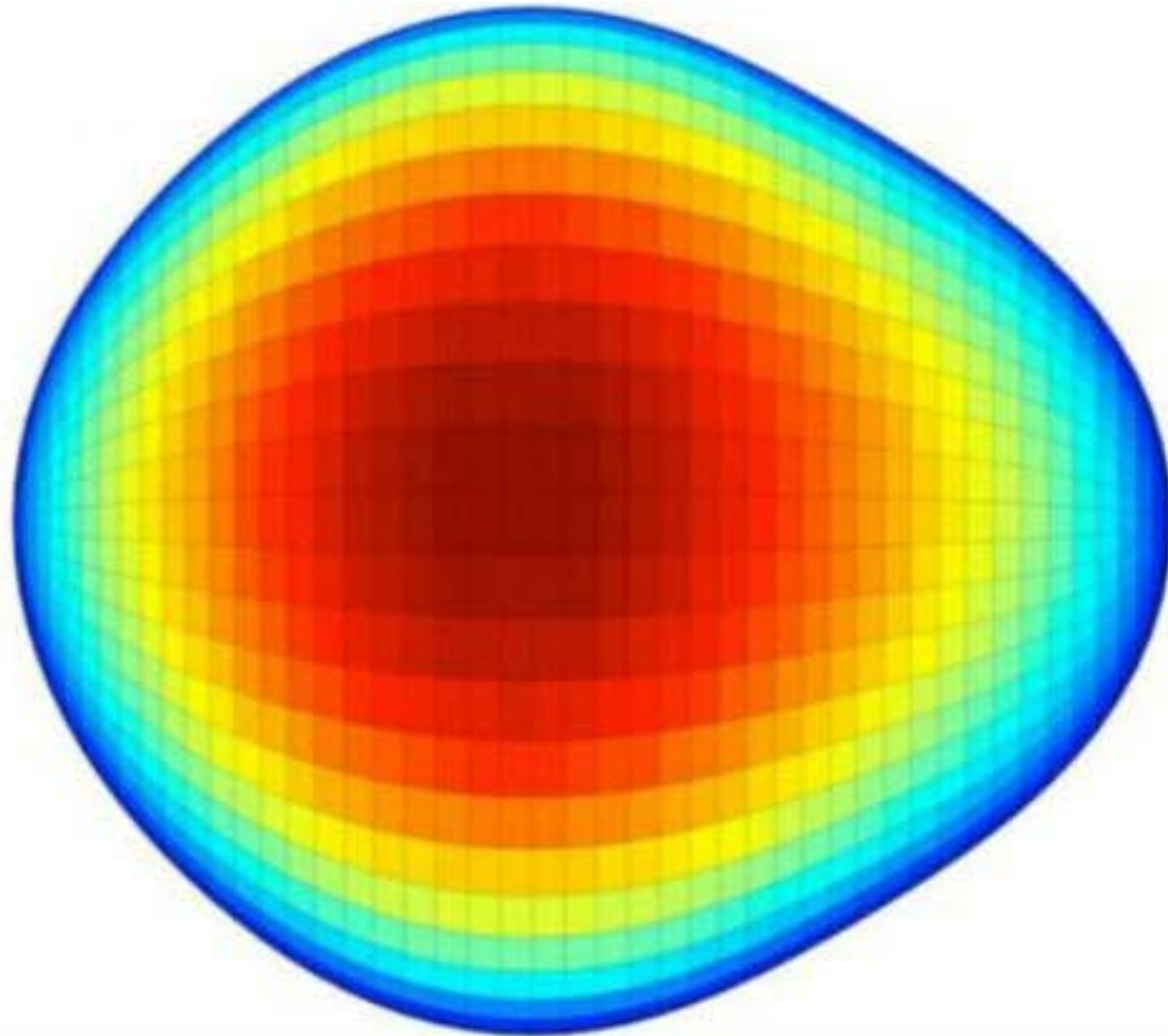


Enhancement of octupole part of nucleon-nucleon force...

- Small energy gap
 - High sub-state density
- } Heaviest nuclei



It's all gone ... pear shaped



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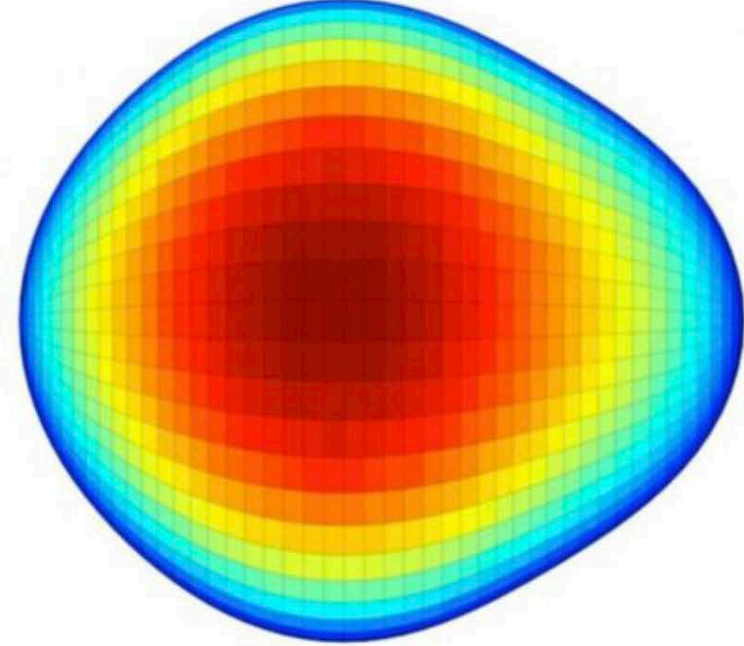
Scotland Scotland Politics Scotland Business Edinburgh, Fife & East Glasgow & West

Pear-shaped nuclei discovery challenges time travel hopes

By Kenneth Macdonald
BBC Scotland Science Correspondent

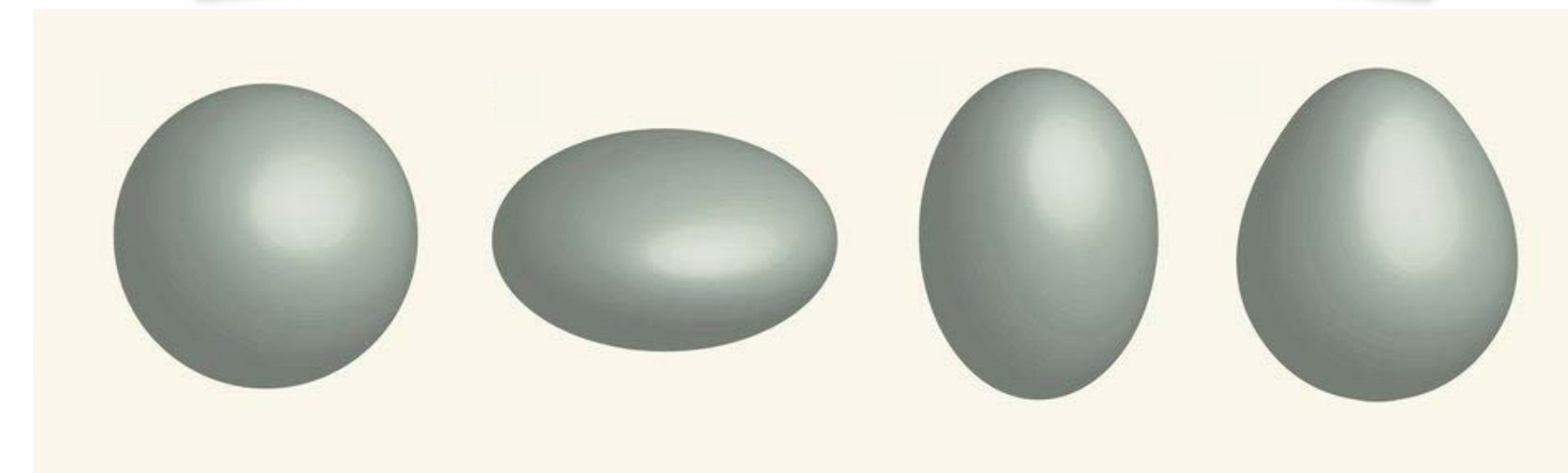
23 June 2016

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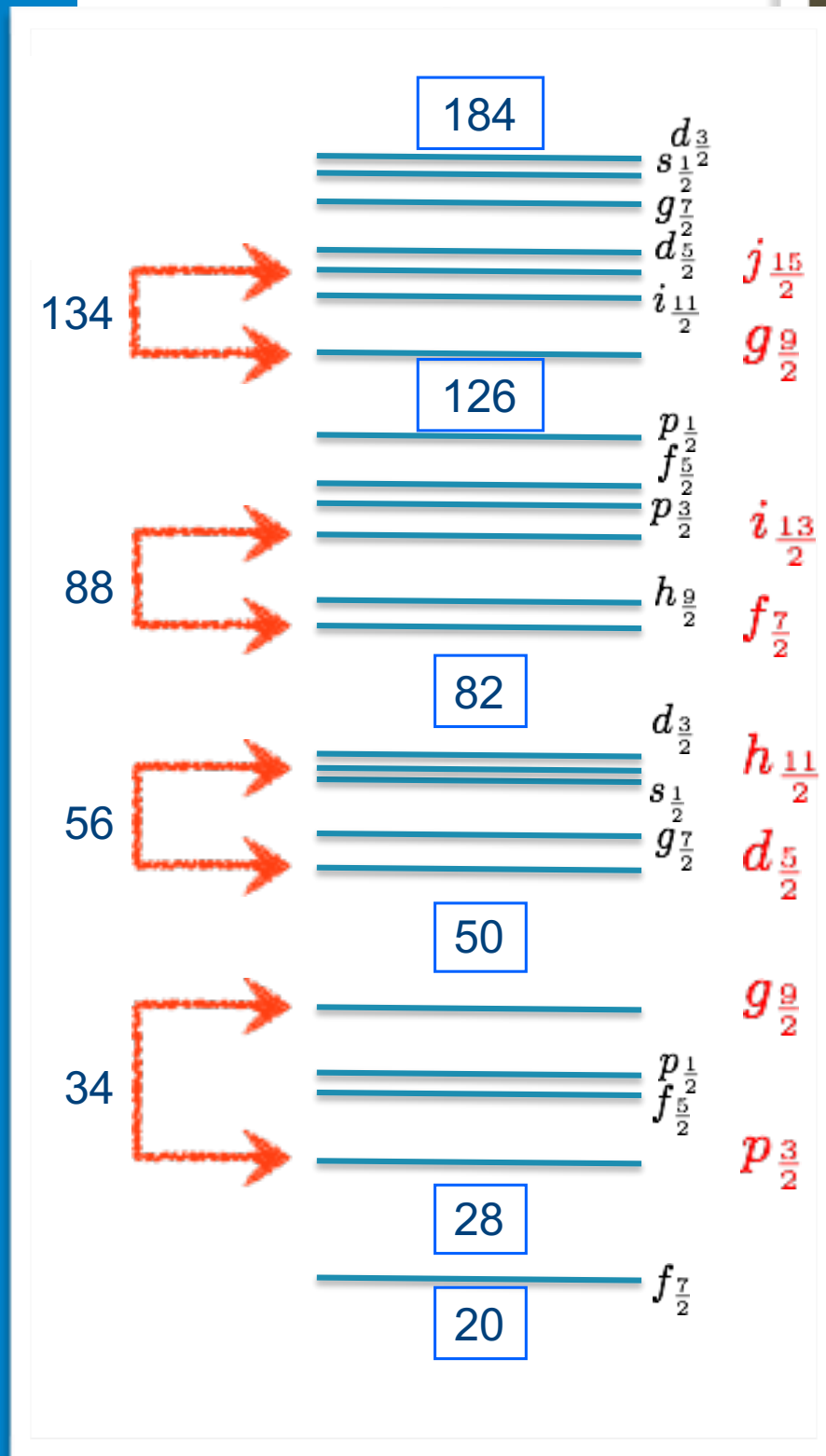
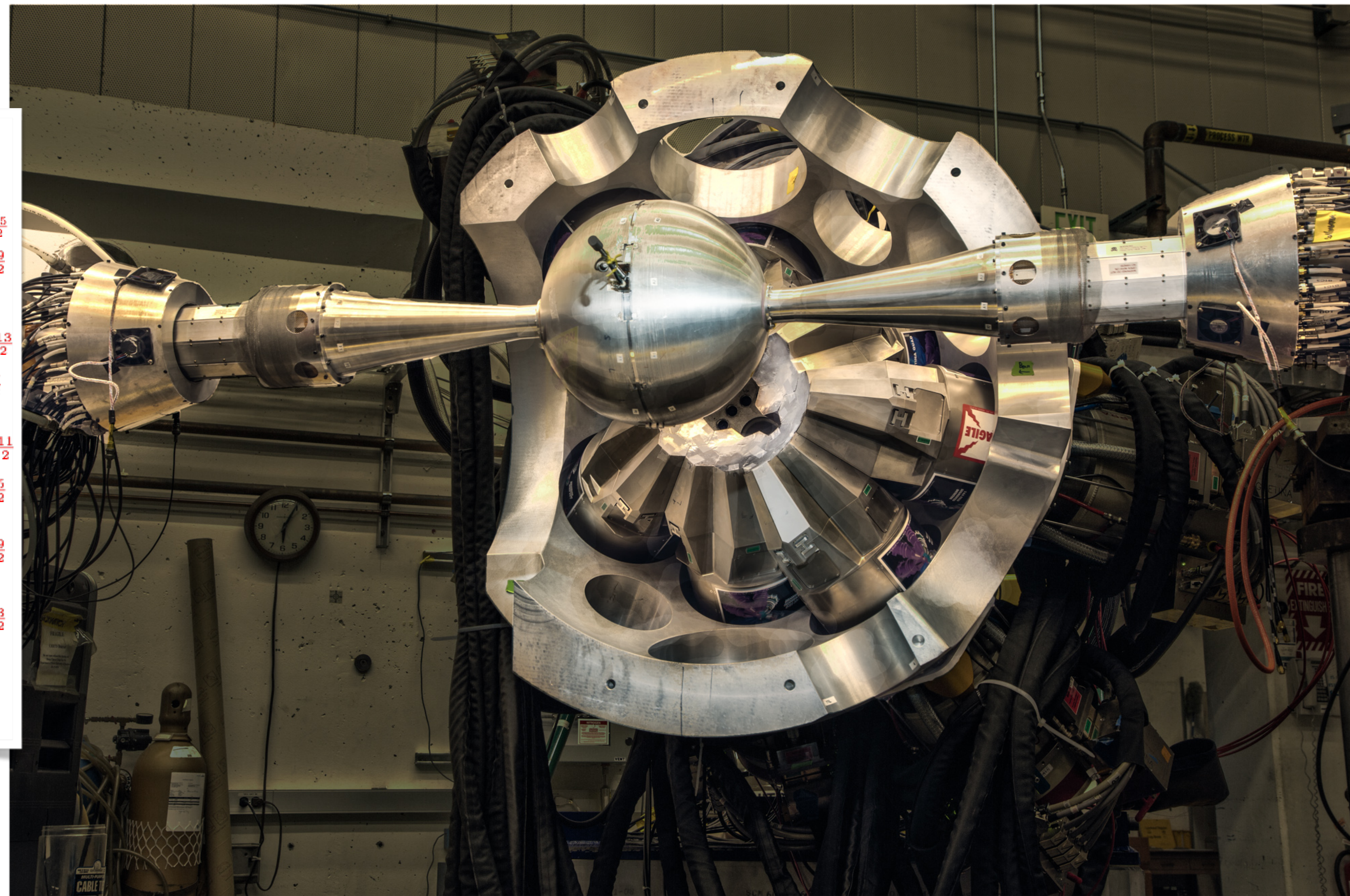
CERN

The pear shape was first discovered in the nucleus of the isotope Radium-224

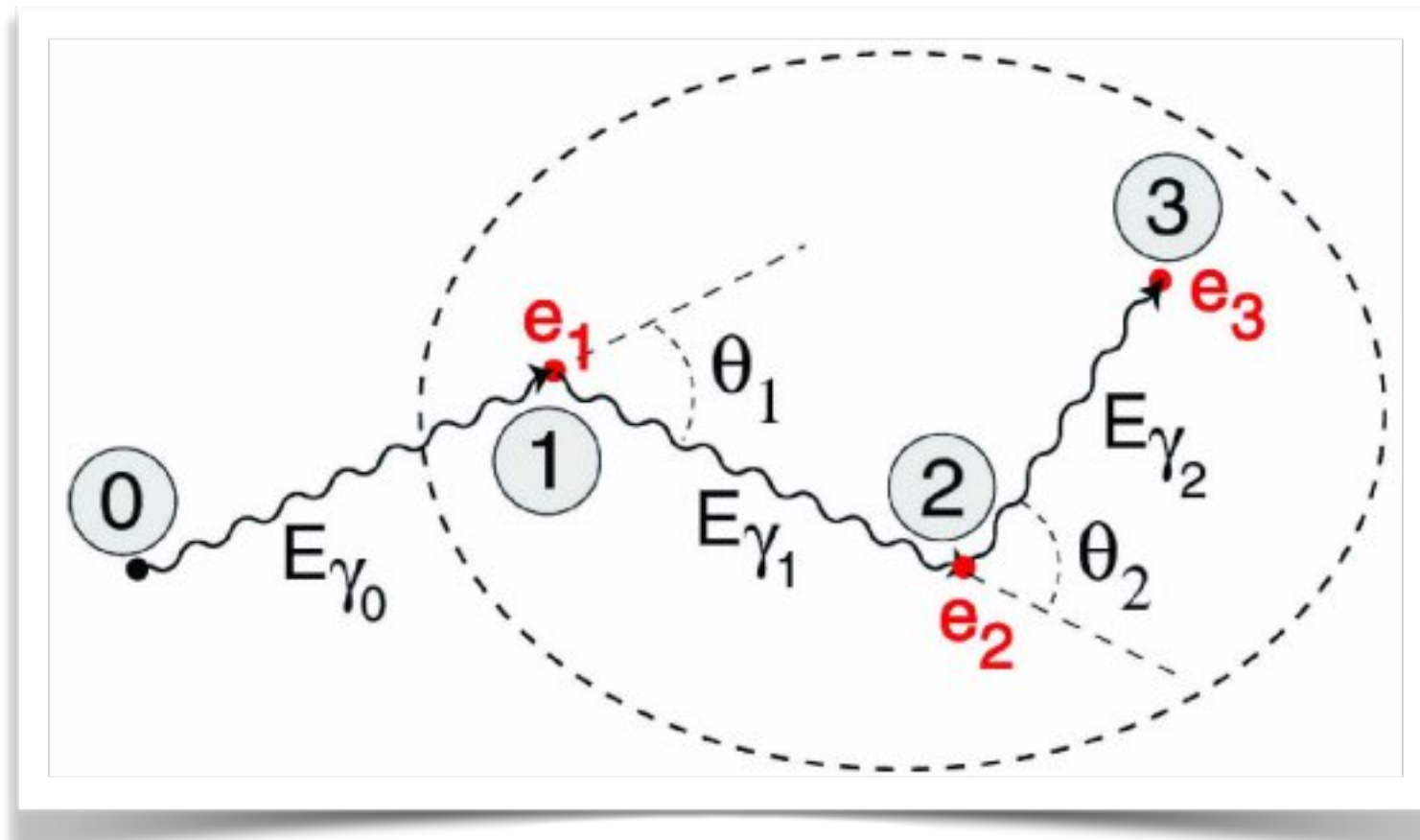


See write up in *Nature* **497**, 190 (2013) [by C. J. (Kim) Lister ...] ... and the BBC

Take a look at the Ba isotopes



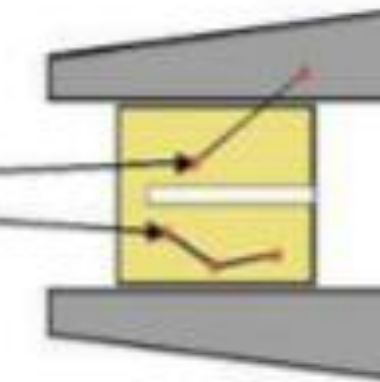
Gamma ray tracking



Motivation of γ -ray tracking

Compton Suppressed

$\epsilon_{ph} \sim 10\%$ $\Omega \sim 40\%$
 $N_{det} \sim 100$ $\theta \sim 8^\circ$



- 50% of solid angle taken by the AC shields
- large opening angle \rightarrow poor energy resolution at high recoil velocity

Ge Sphere

$\epsilon_{ph} \sim 50\%$ $\Omega \sim 80\%$
 $N_{det} \sim 1000$ $\theta \sim 3^\circ$



- too many detectors needed to avoid summing effects
- opening angle still too big for very high recoil velocity

Tracking Array

$\epsilon_{ph} \sim 50\%$ $\Omega \sim 80\%$
 $N_{det} \sim 100$ $\theta \sim 10^\circ$



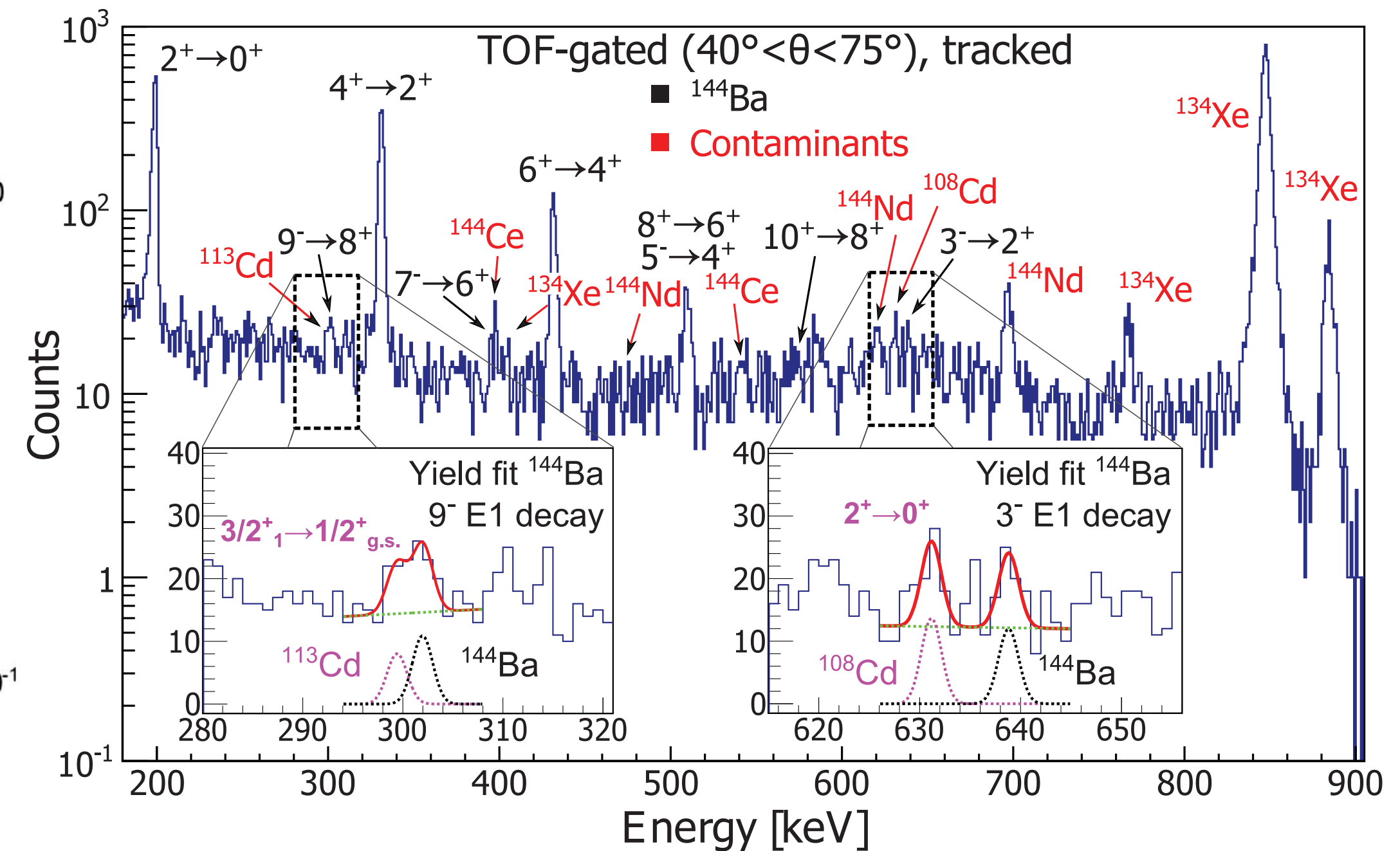
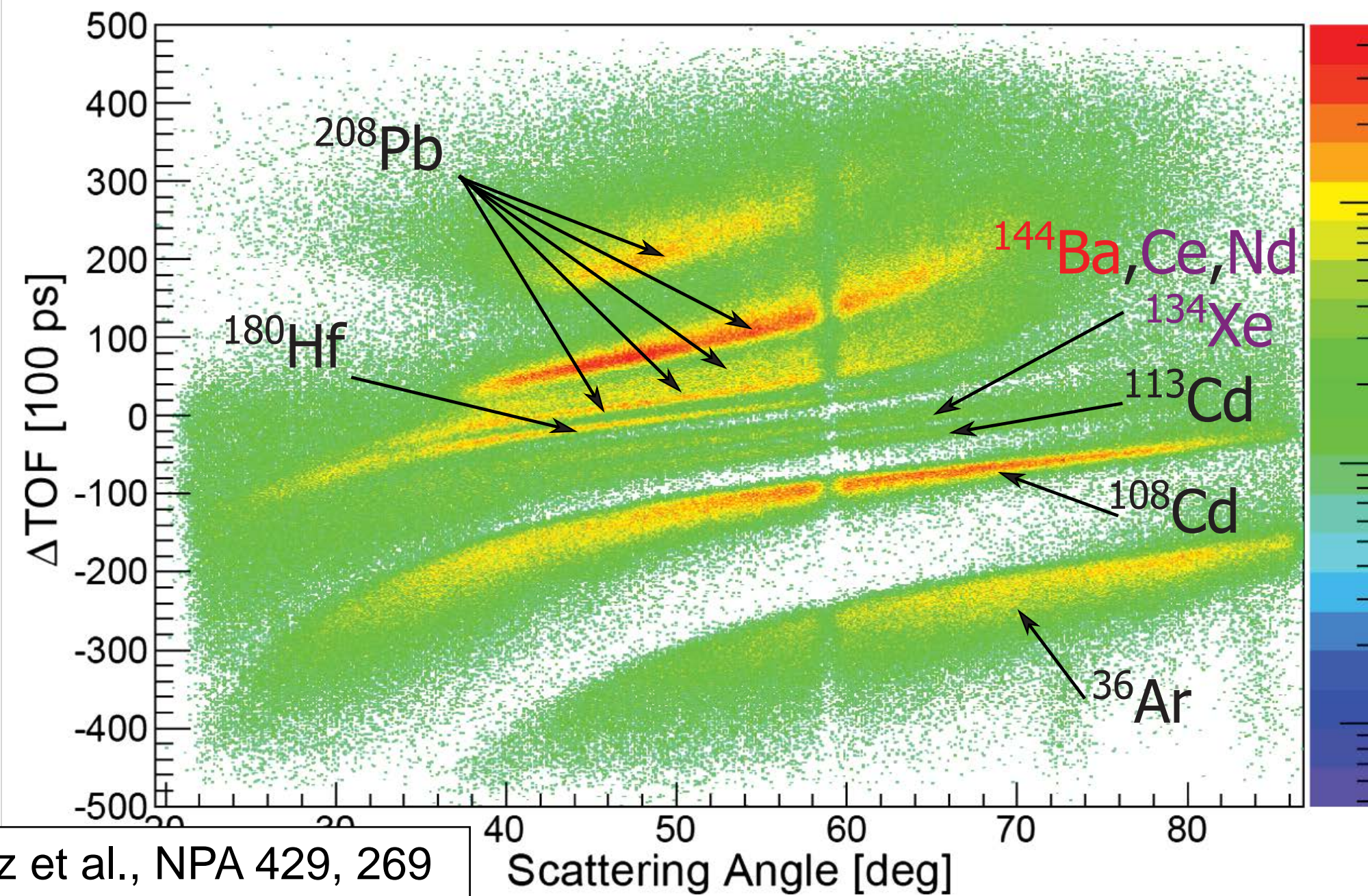
Smarter use of Ge detectors

- segmented detectors
- digital electronics
- timestamping of events
- analysis of pulse shapes
- tracking of γ -rays

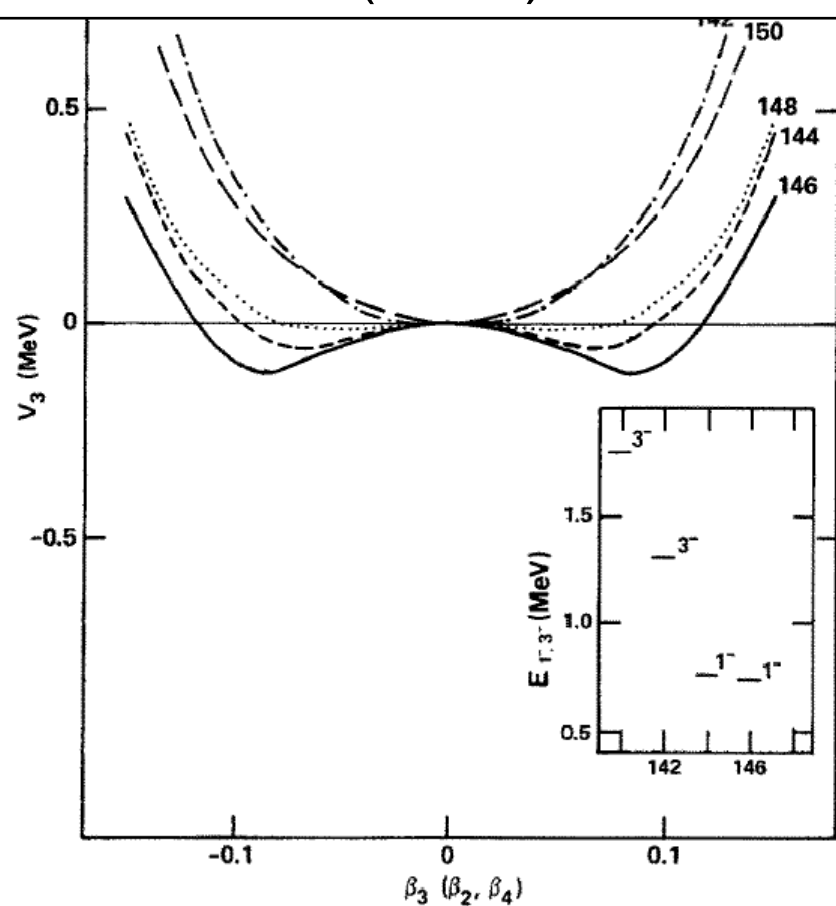
Pulse Shape Analysis \rightarrow $\theta_{eff} \sim 1^\circ$
 Gamma-ray Tracking \rightarrow $N_{eff} \sim 10000$

from Calorimetric to Position Sensitive operation mode

Coulomb excitation of ^{144}Ba



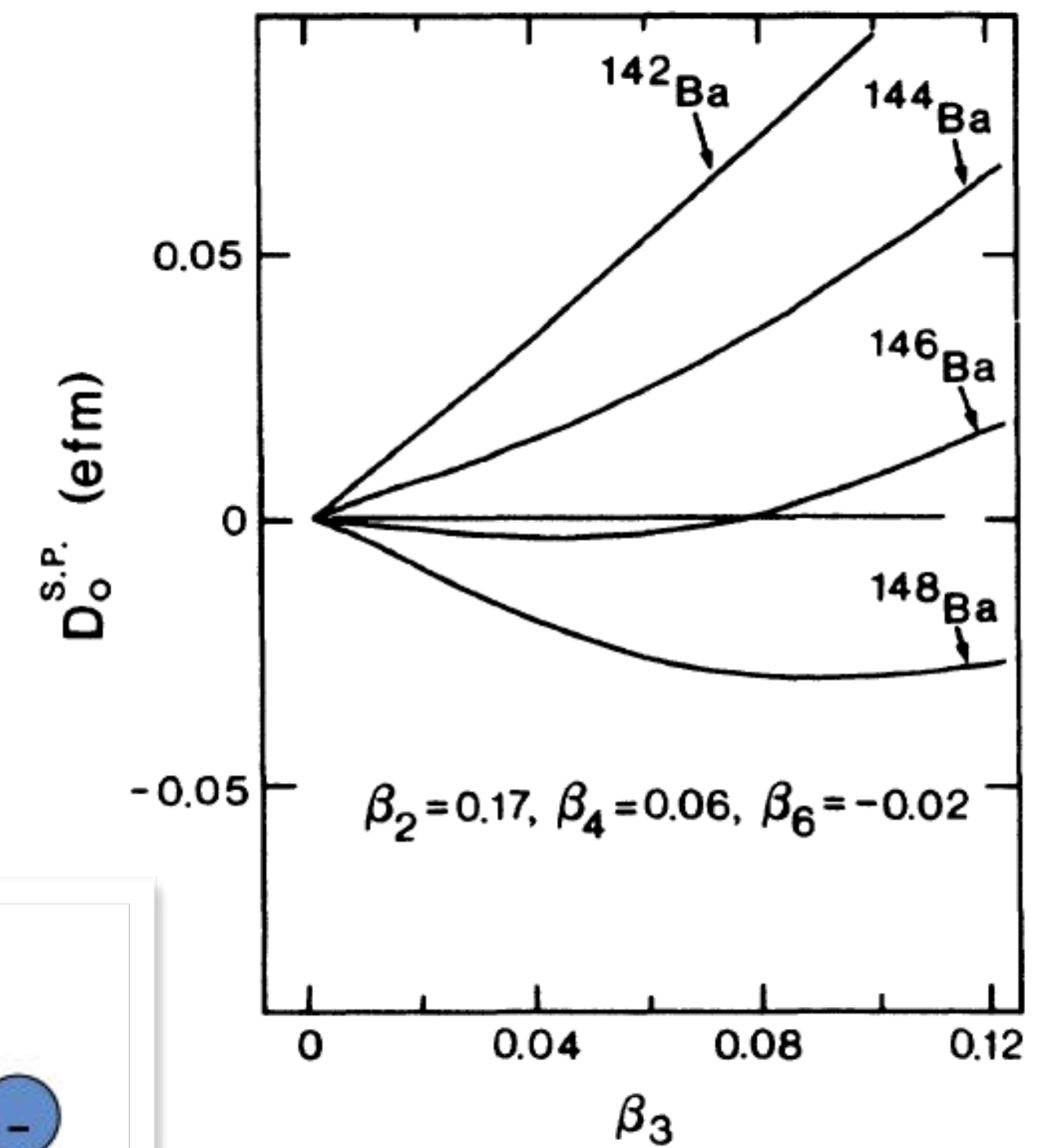
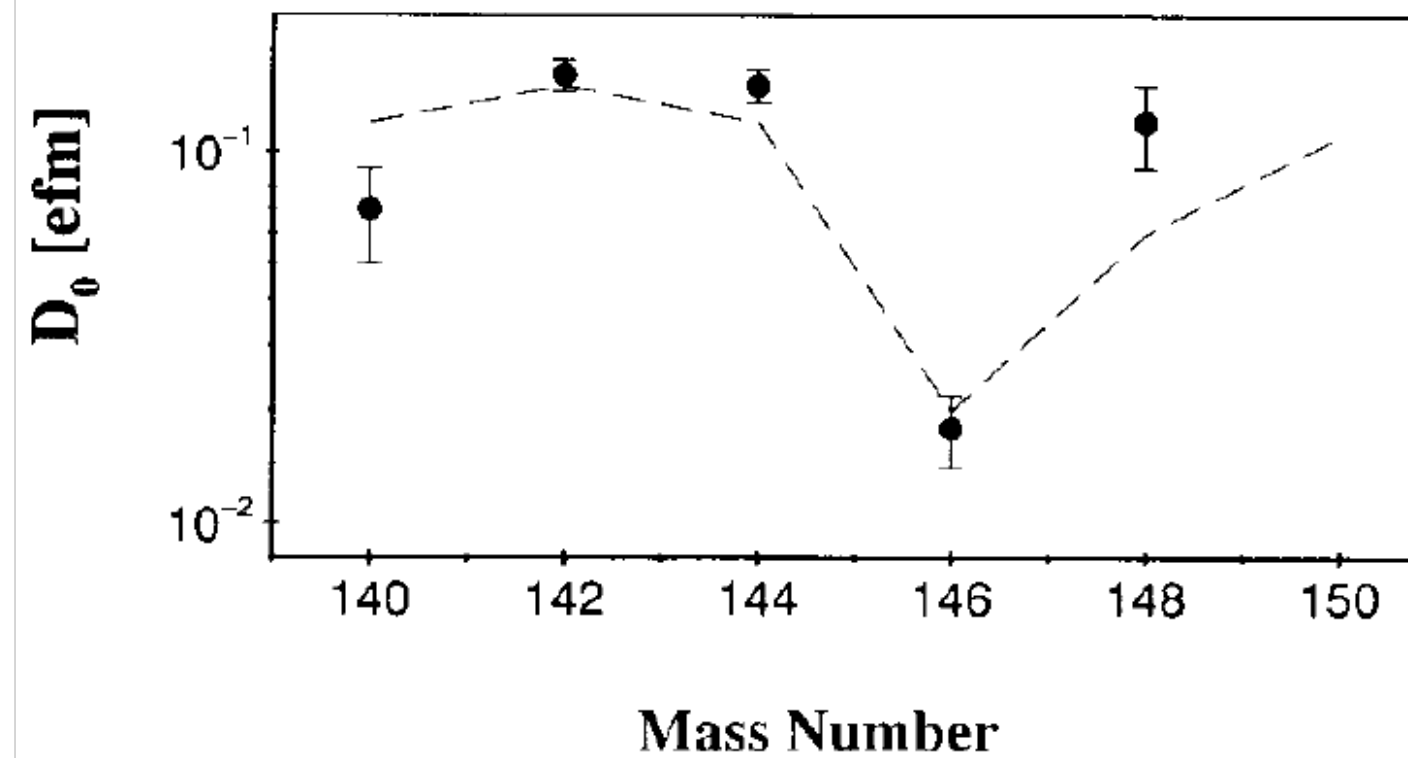
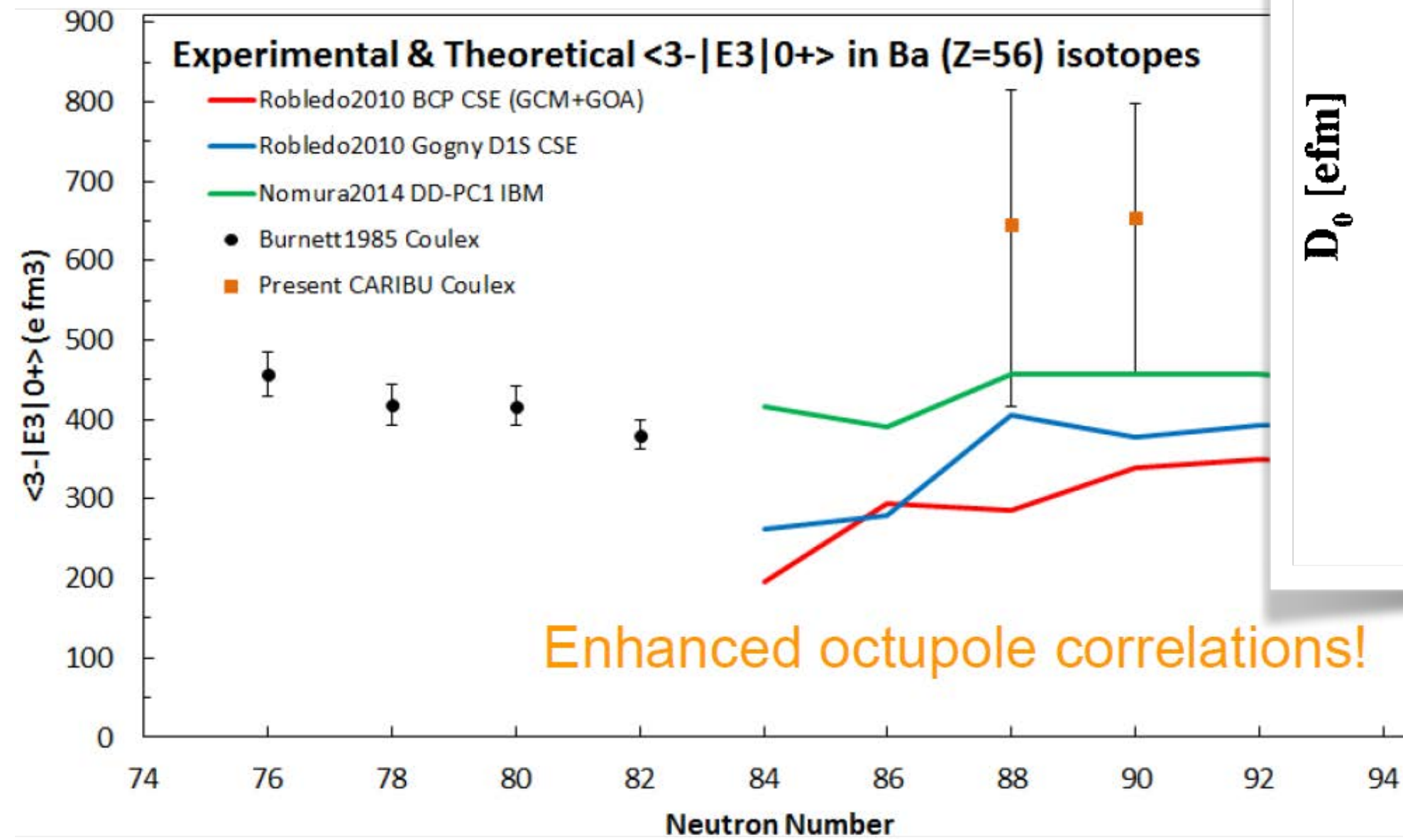
W. Nazarewicz et al., NPA 429, 269 (1984)



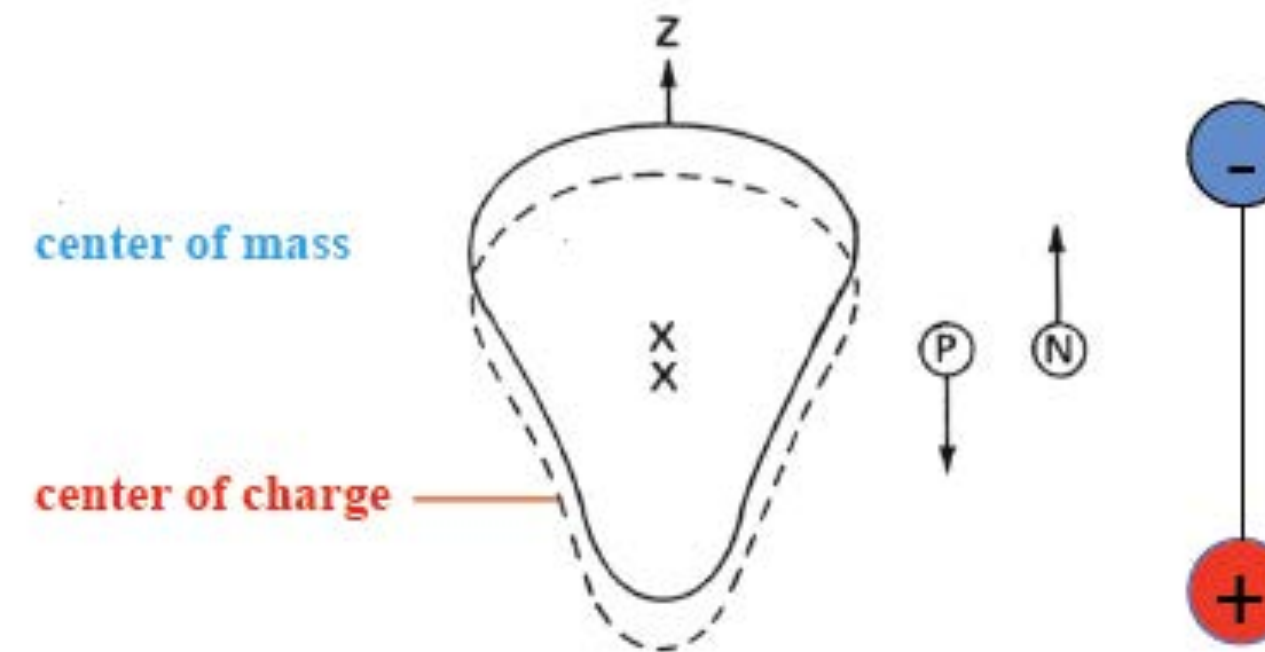
B. Bucher, S. Zhu et al., Phys. Rev. Lett. 116, 112503 (2016)

- ^{144}Ba lies in a region where suspected **enhanced octupole correlations** occur
 - where long-range interactions between **$\Delta j = \Delta l = 3$** configurations, namely the $\pi h_{11/2} \otimes \pi d_{5/2}$ and $\nu i_{13/2} \otimes \nu f_{7/2}$, occur
- Coulomb excitation is a reliable probe to extract $B(E3)$ values
 - **$B(E3) = 48(+25_{-34})$ W.u., consistent with octupole collectivity**

Coulomb excitation of ^{144}Ba



Induced Electric Dipole Moment D_0



- Enhanced octupole correlations firmly established in both ^{144}Ba and ^{146}Ba
- Suggests a 'region' not just isolated cases (how far does it extend?)
- Behavior of the dipole strength shows interesting behavior

$$D_0 = c_{1d} Z A \beta_2 \beta_3$$

$$D_0^{\text{s.p.}} = \frac{e}{3.6} \left[\frac{N}{A} \langle z \rangle_p - \frac{Z}{A} \langle z \rangle_n \right]$$

Summary

- *The many varied techniques associated with determining even simple properties of nuclei can give tremendous insights ... and we have not even delved (much) into the microscopic structure of these systems yet*
- *... next two lectures focus on single-particle structure as probed through direct reactions*